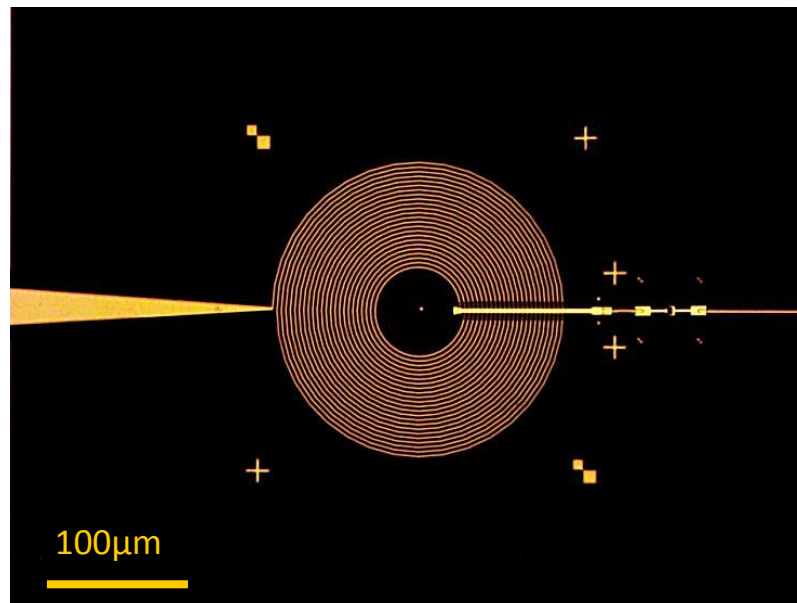


Anti-bunched photons emitted by a dc biased Josephson junction

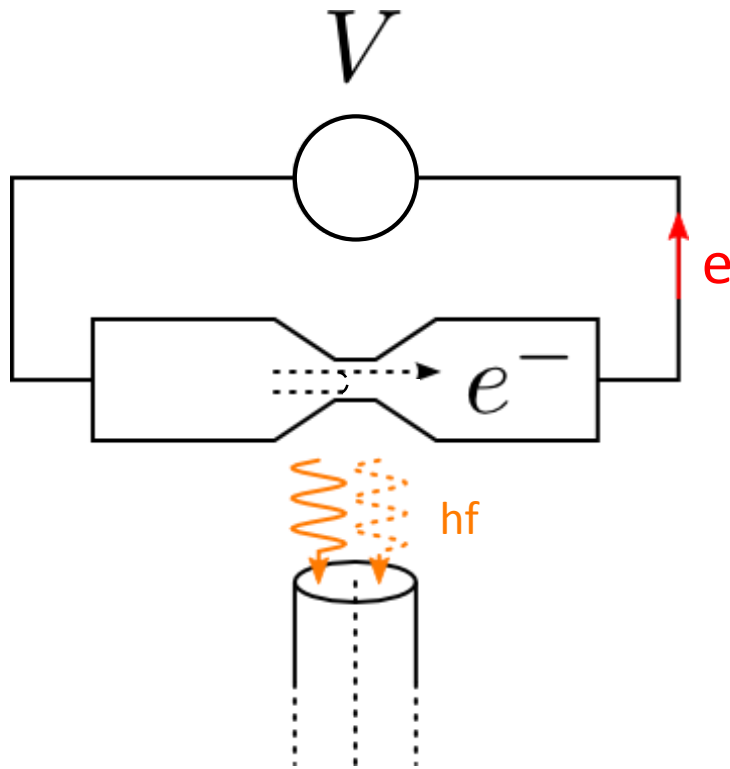
Chloé Rolland, M. Westig, I. Moukharski, D. Vion, H. Le sueur, P. Joyez, C. Altimiras, P. Roche, D. Estève and Fabien Portier

collab.: J. Ankerhold, B. Kubala

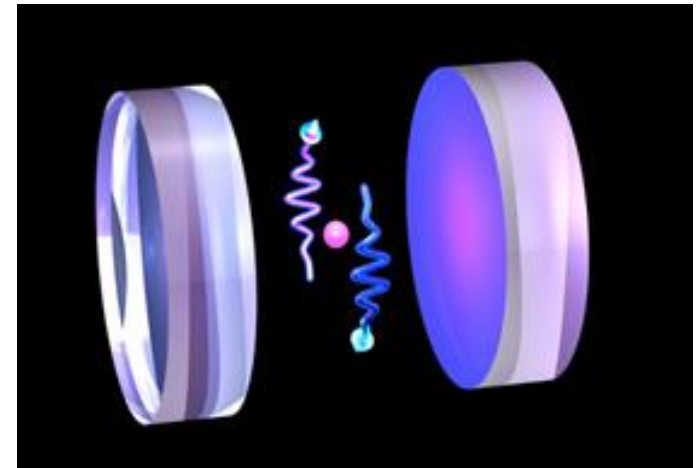


Context: quantum optics of quantum conductors

Quantum conductors



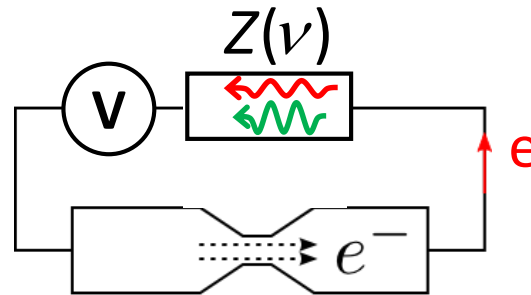
Atomic Physics: cavity QED



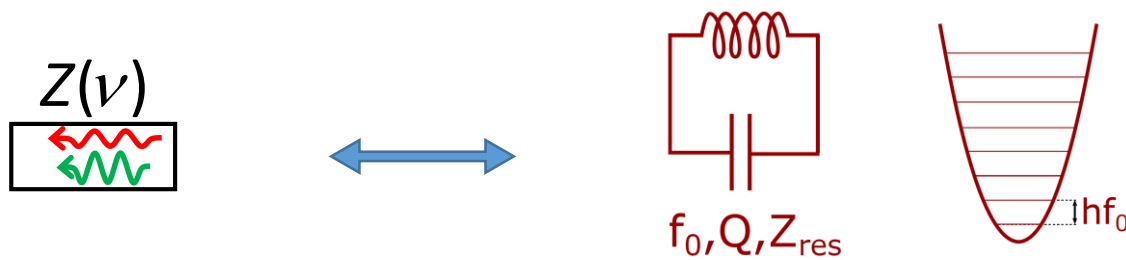
⇒ Many open questions (properties of the emitted radiation, **strong feedback** of coupling to electromagnetic modes)

⇒ Excellent understanding (dressed atom formalism...)

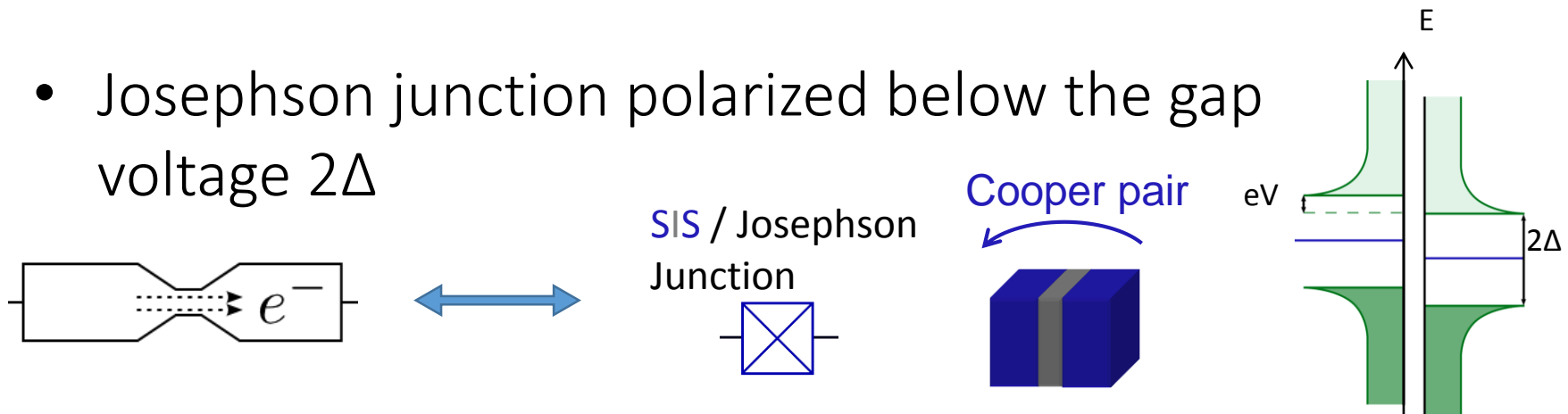
Simplifying the system



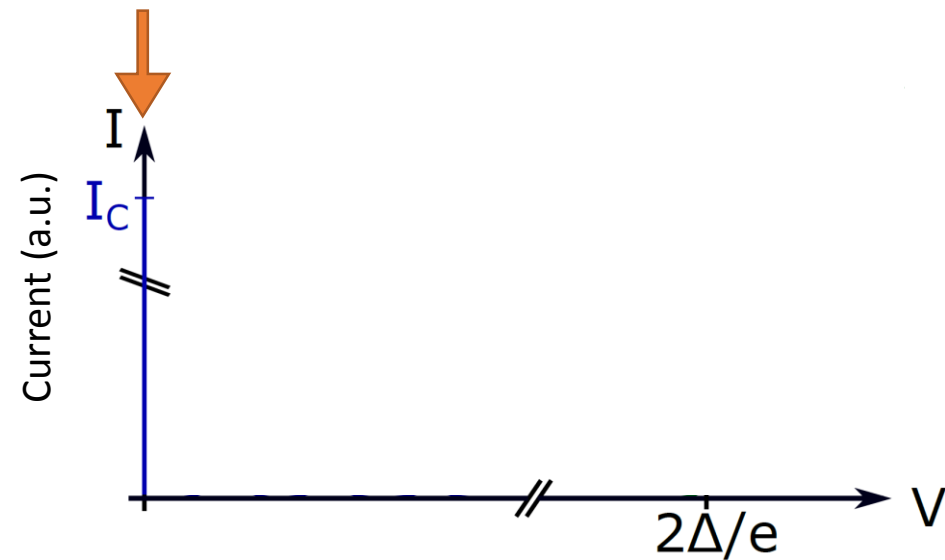
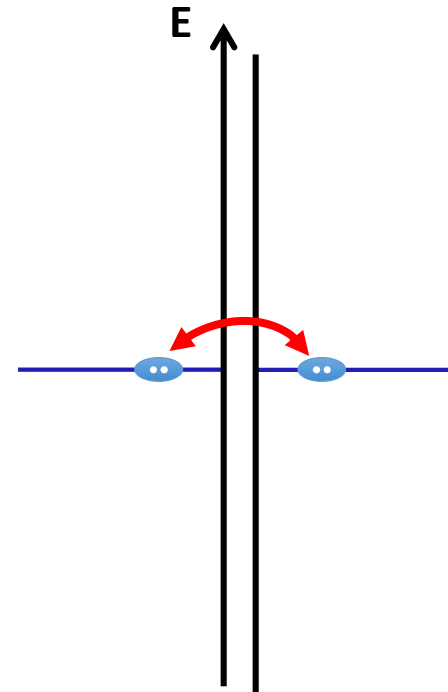
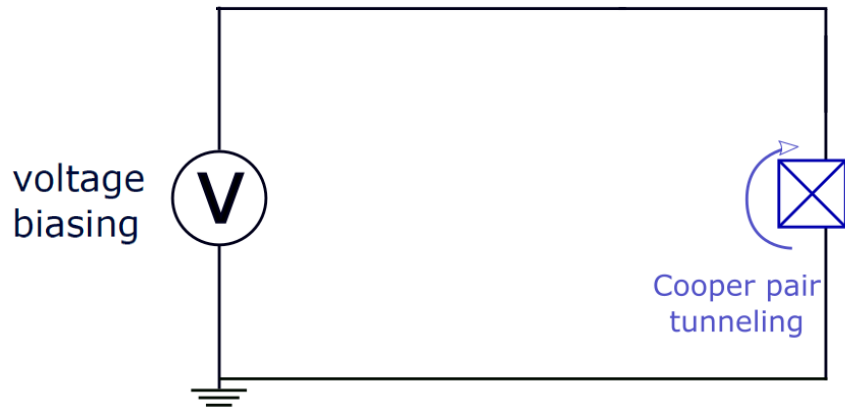
- Environment: single mode



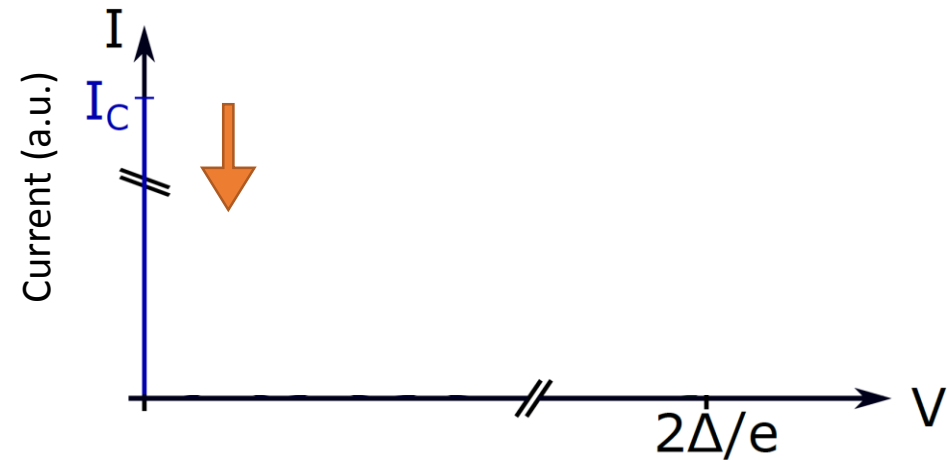
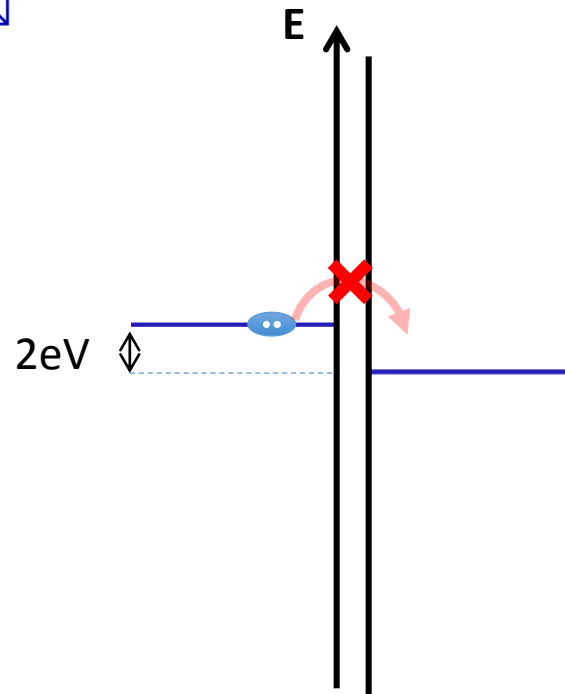
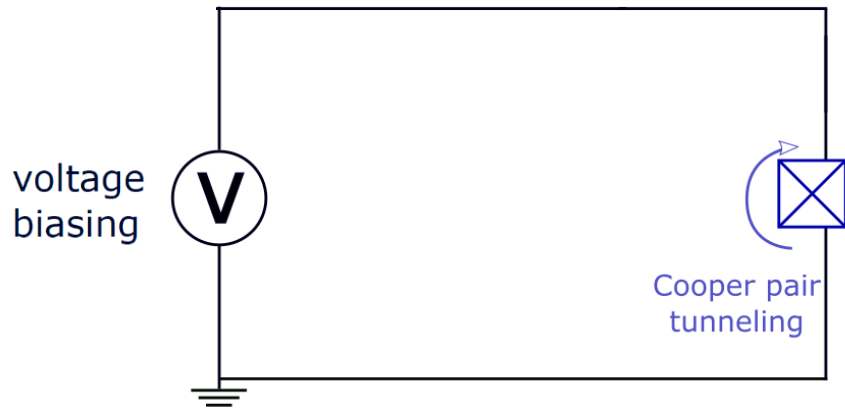
- Josephson junction polarized below the gap voltage 2Δ



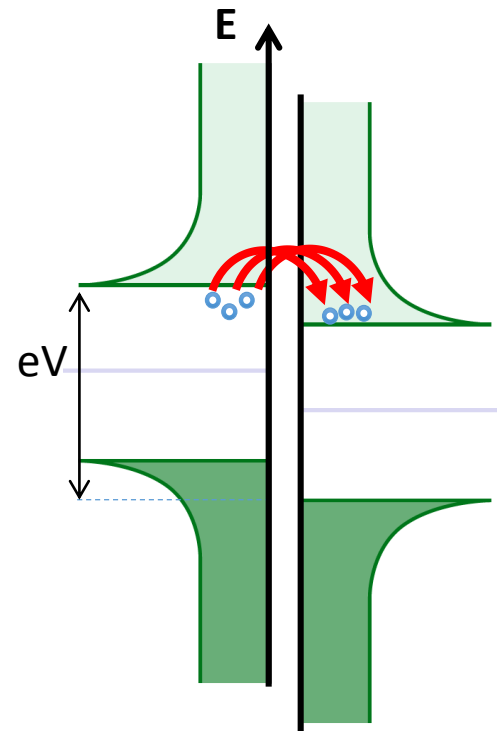
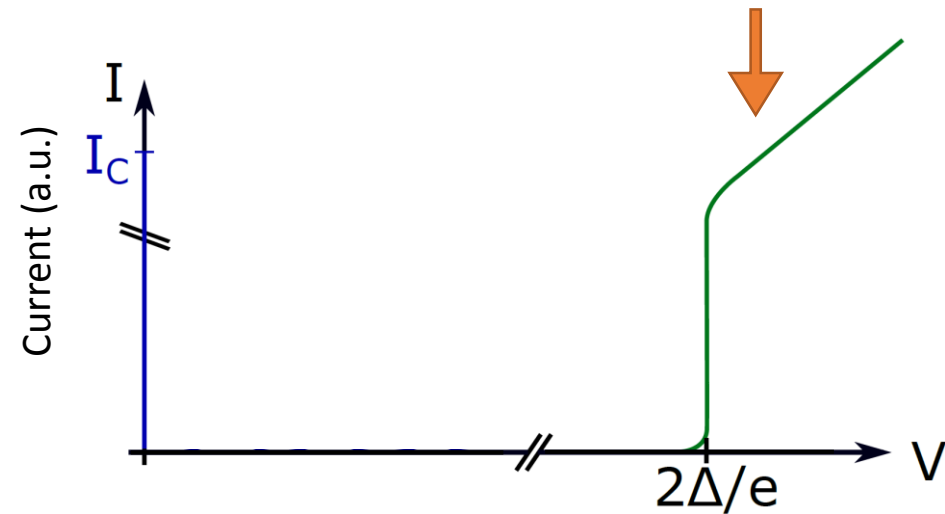
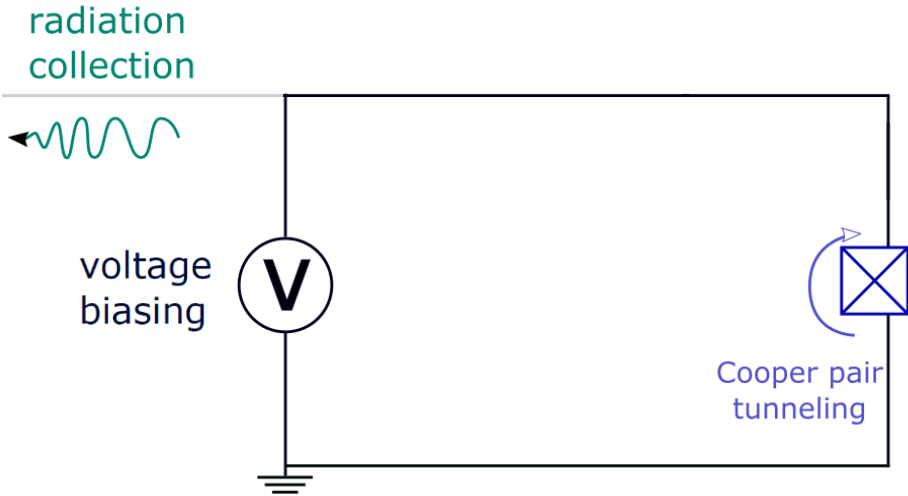
Circuit



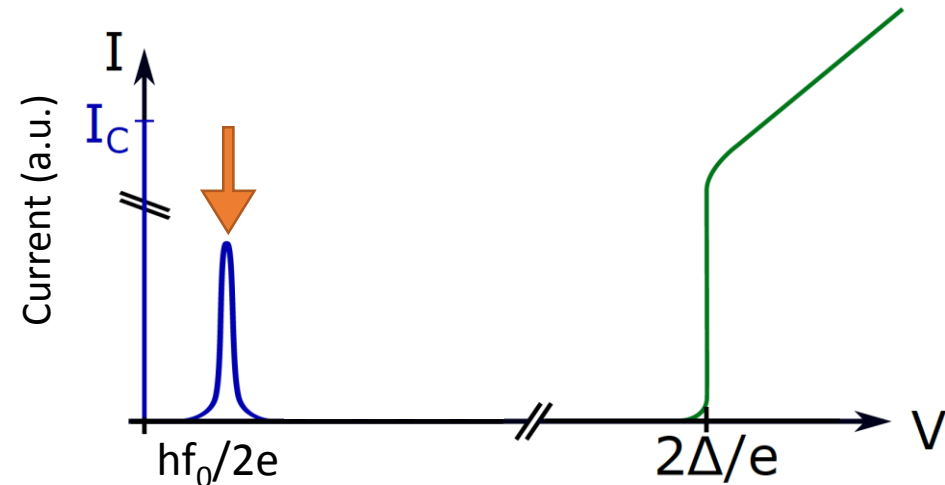
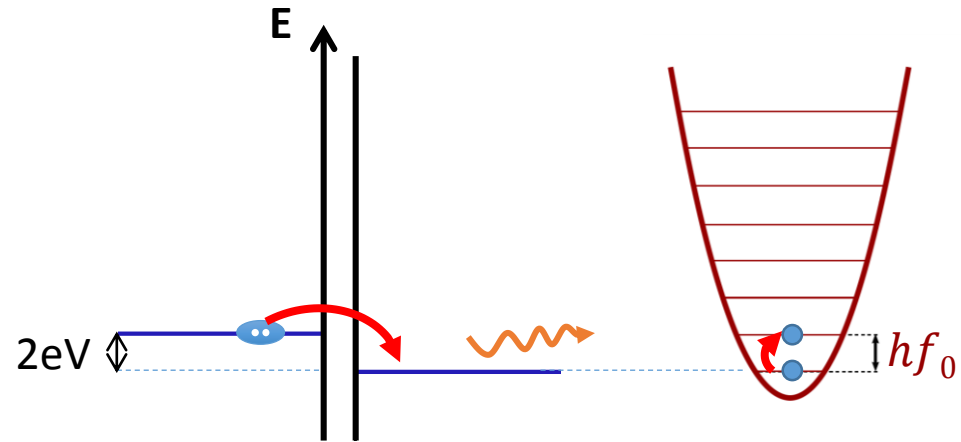
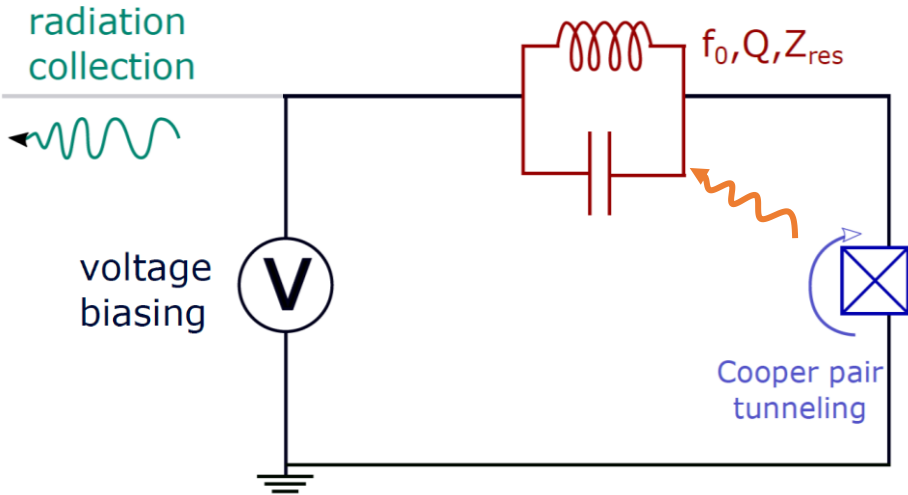
Circuit



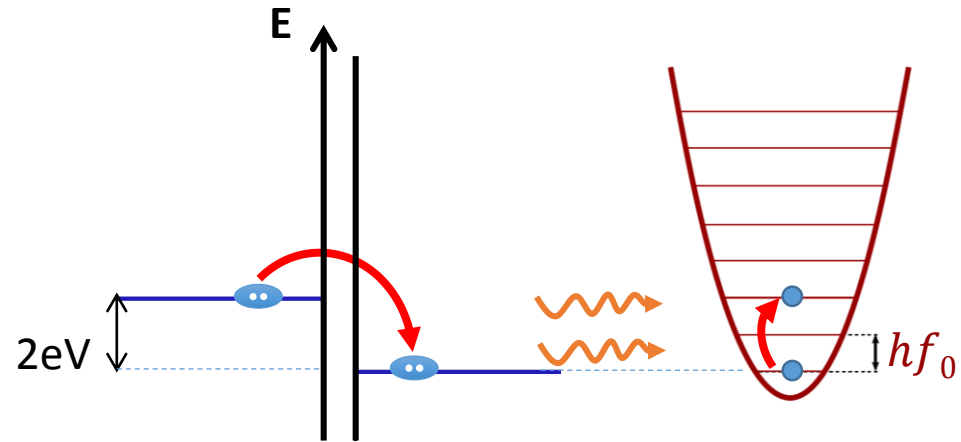
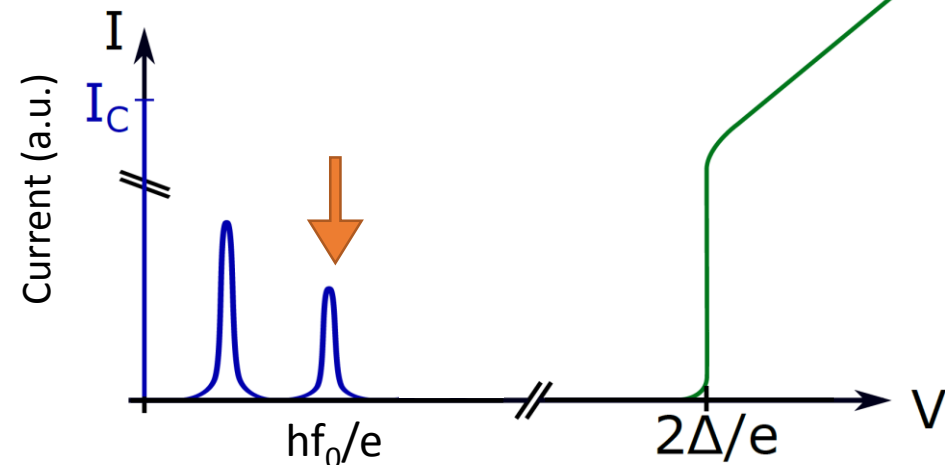
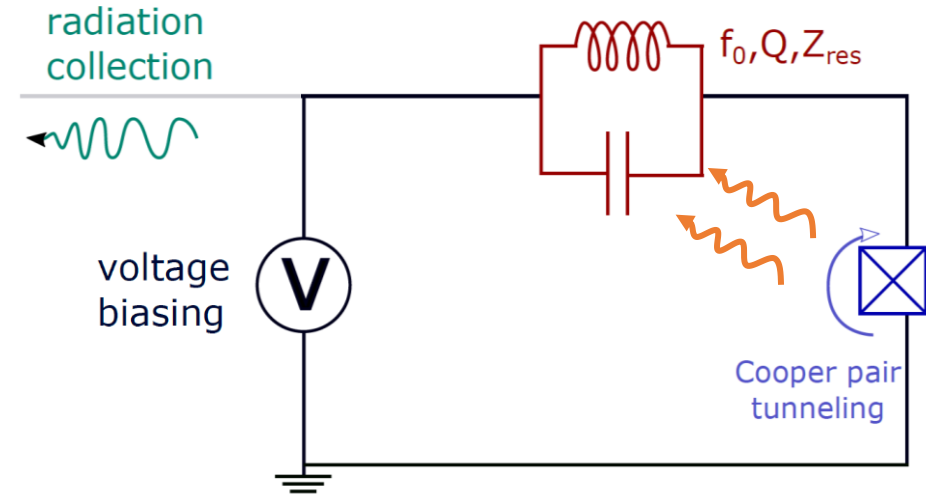
Circuit



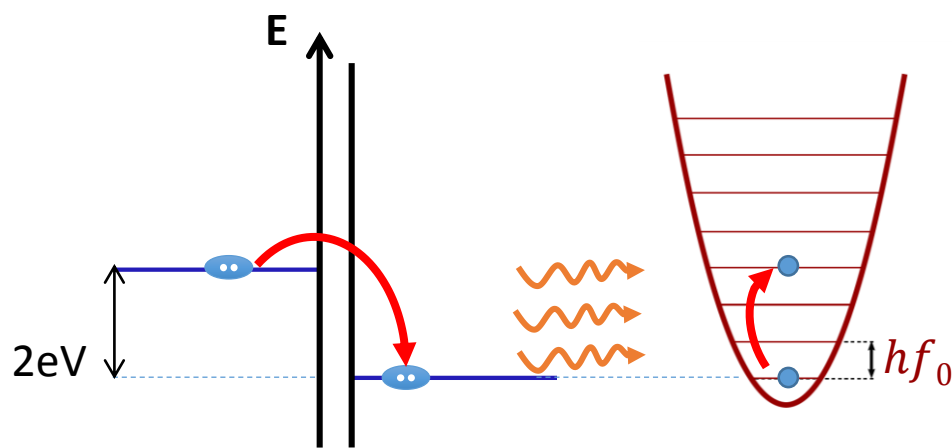
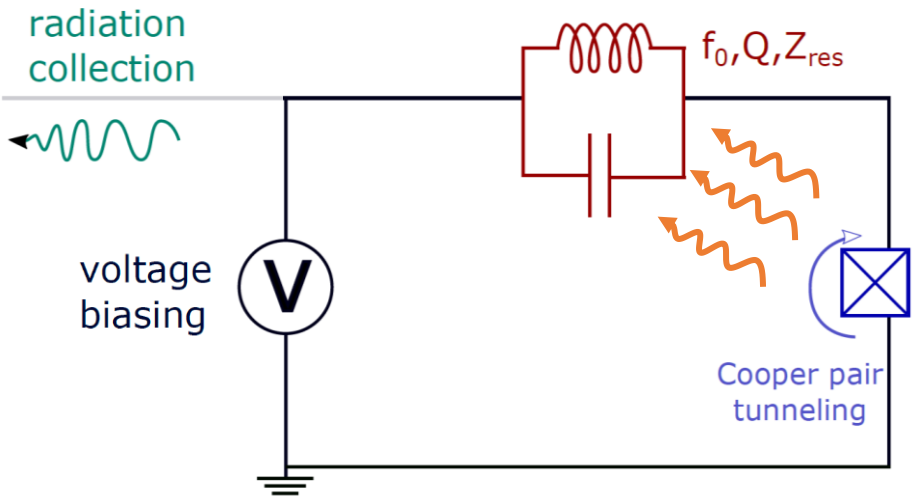
Circuit



Circuit

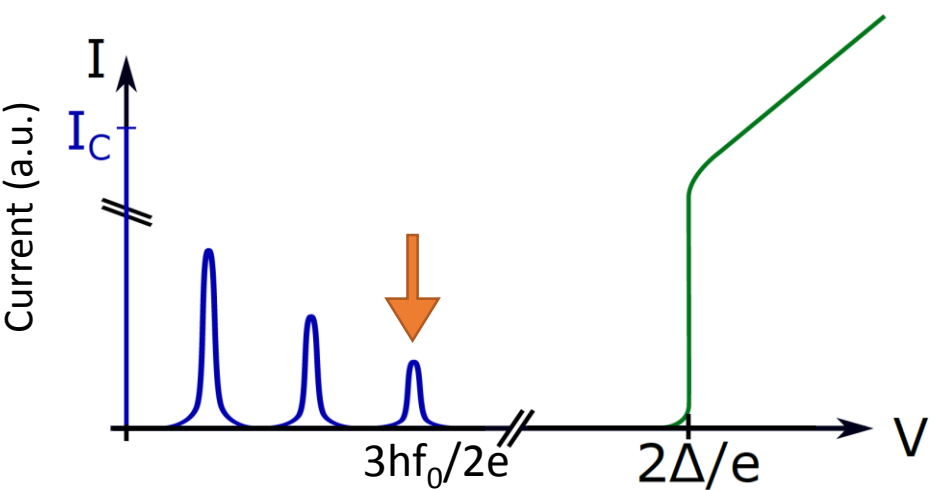


dc current only if photons can be absorbed



$$2eV = khf_0$$

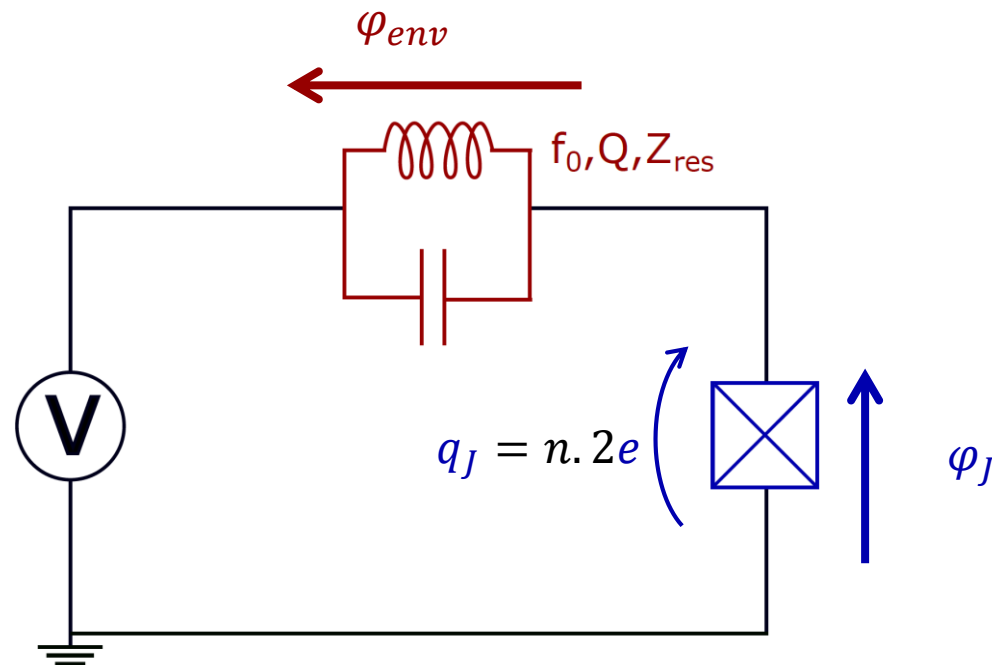
Emission rates ?



Hamiltonian of the system

$$H = \underbrace{hf_0 (a^\dagger a + 1/2)}_{\text{Resonator}} - \underbrace{E_J \cdot \cos(\varphi_J)}_{\text{CP tunneling}}$$

$$[\varphi, q] = i\hbar$$



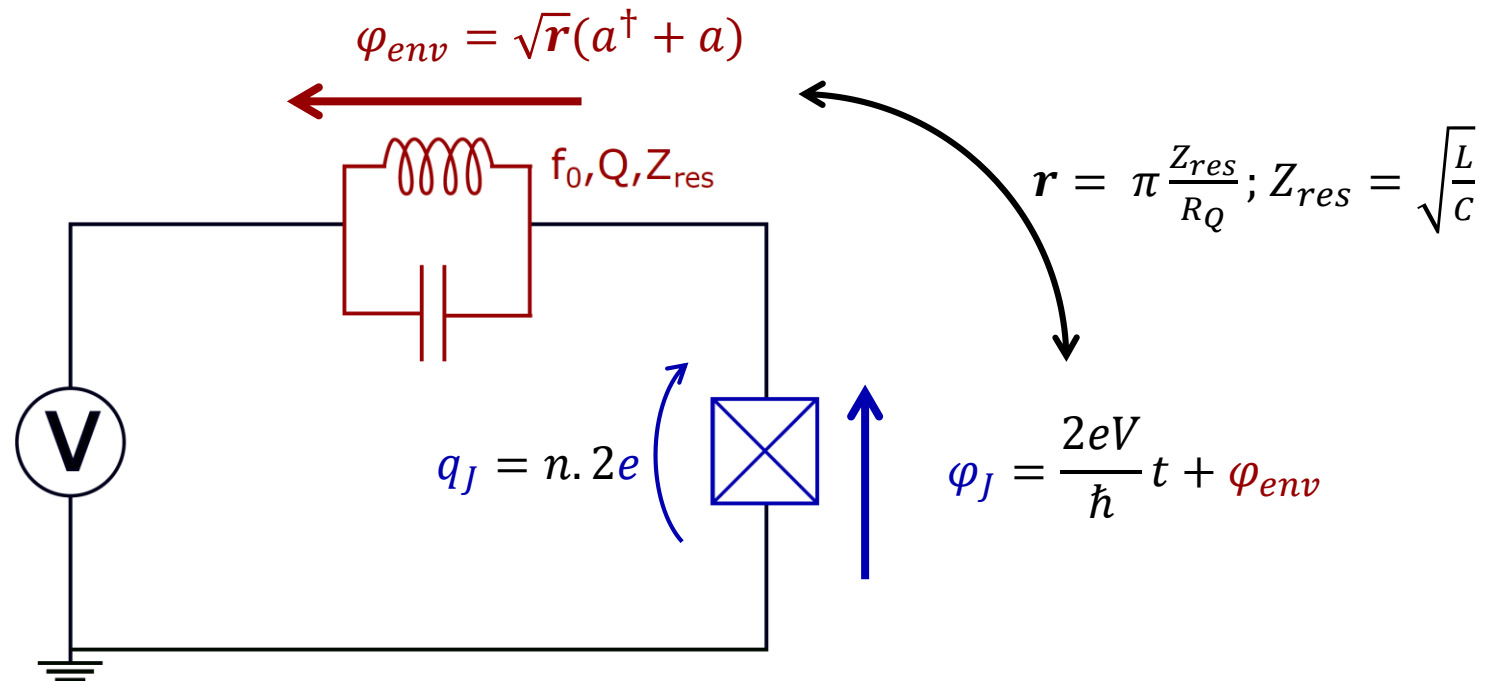
D. Averin, Y. Nazarov, and A. Odintsov, Physica B 165-166, 945 (1990)

Ingold & Nazarov, arxiv:0508728 (1992)

Hamiltonian of the system

$$H = \underbrace{hf_0 (a^\dagger a + 1/2)}_{\text{Resonator}} - \underbrace{\frac{E_J}{2} (e^{i\varphi_J} + e^{-i\varphi_J})}_{\text{CP tunneling}}$$

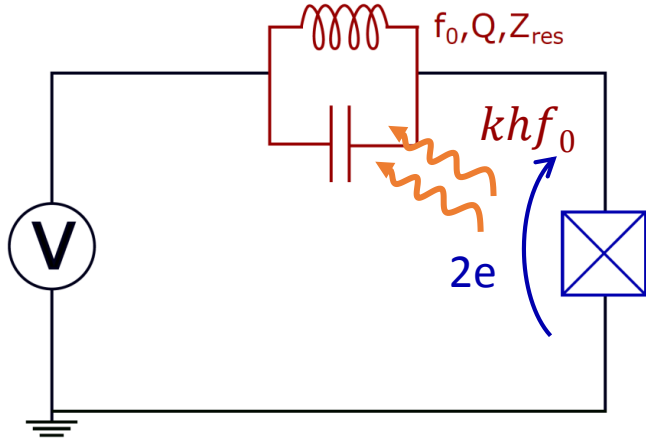
$$[\varphi, q] = i\hbar$$



D. Averin, Y. Nazarov, and A. Odintsov, Physica B 165-166, 945 (1990)

Ingold & Nazarov, arxiv:0508728 (1992)

Cooper pair and photon rates



$$H = hf_0 (a^\dagger a + 1/2) - \frac{E_J}{2} (e^{i\varphi_J} + e^{-i\varphi_J})$$

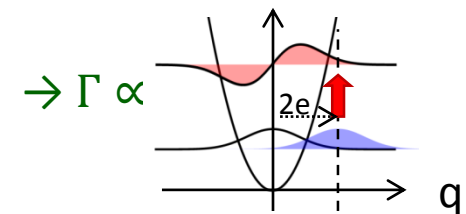
$$r = \pi \frac{Z_{res}}{R_Q}$$

Emission rates (from empty resonator):

$$\Gamma^{ph}(2eV = khf_0) = k \cdot \Gamma^{2e}(2eV = khf_0)$$

$$\begin{aligned} \Gamma^{2e}(2eV) &= \frac{\pi E_J^2}{2\hbar} \sum_k |\langle k | e^{-2i\varphi_J} | 0 \rangle|^2 \delta(2eV - khf_0) \\ &= \frac{\pi E_J^2}{2\hbar} \sum_k e^{-r} \frac{r^k}{k!} \delta(2eV - khf_0) \end{aligned}$$

→ k photons / Cooper pair

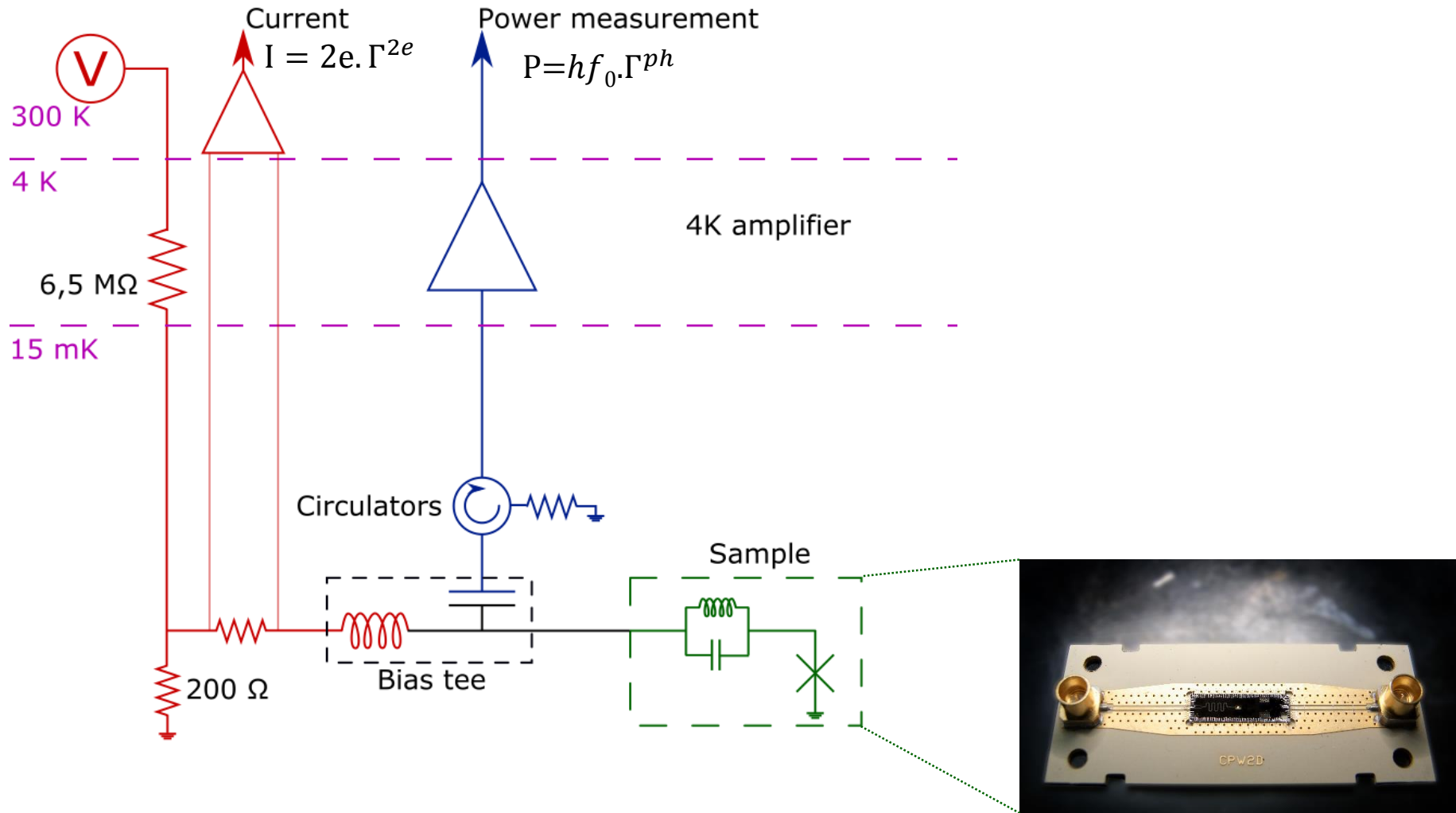


→ Large r favors multiphoton processes

D. Averin, Y. Nazarov, and A. Odintsov, Physica B 165-166, 945 (1990)

H. Pothier, PhD dissertation (1991)

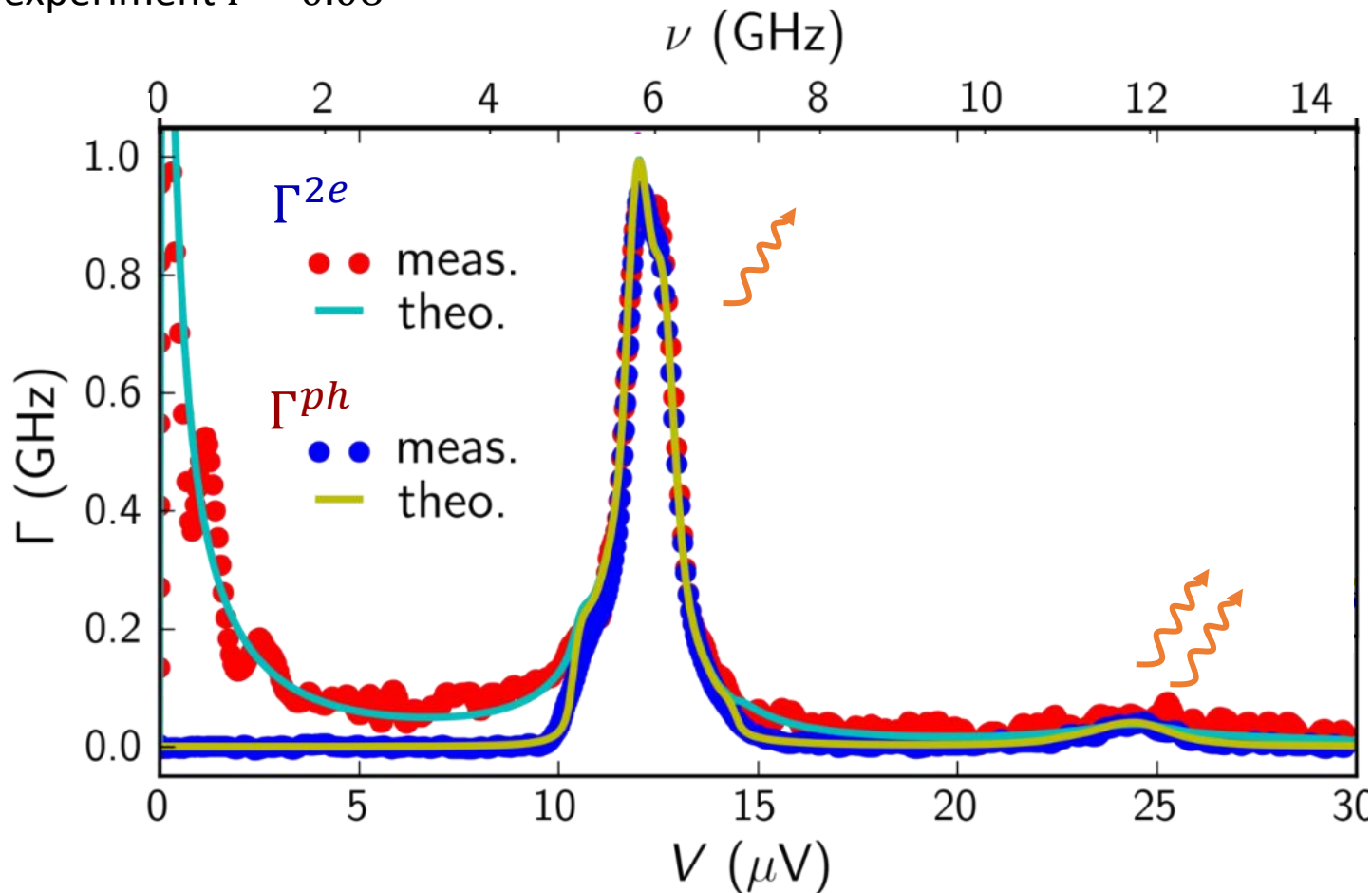
Measurement setup



Hofheinz *et al.*, PRL **106**, 217005 (2011)

Current vs Power measurements

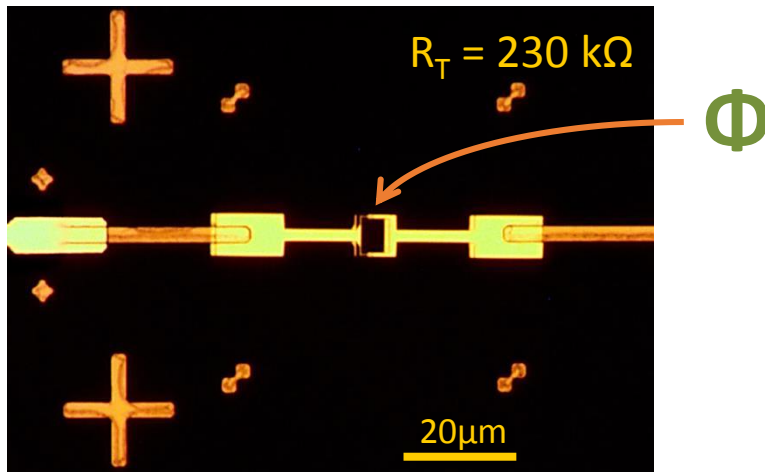
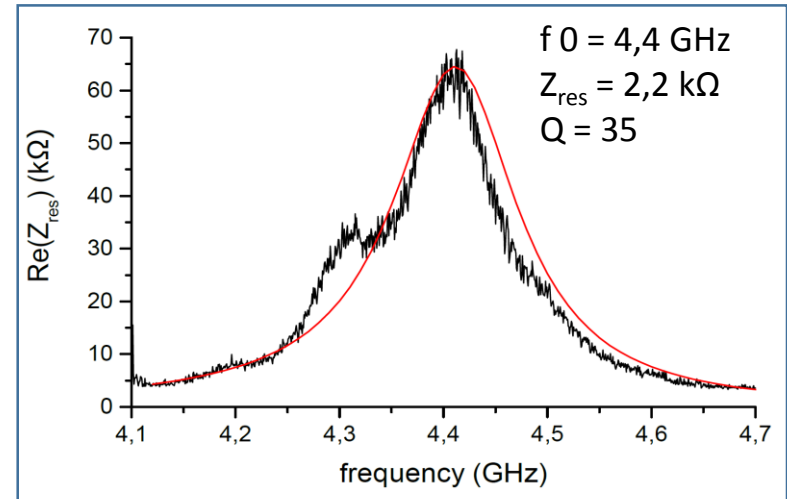
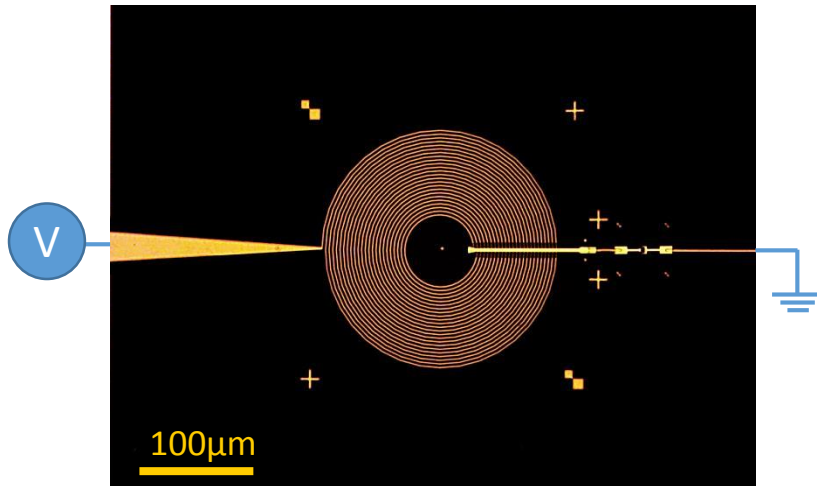
First experiment $r = 0.08$



Cooper pair and photon mean rates match!

High impedance resonator

$$\omega_0 = \frac{1}{\sqrt{LC}} \in [4; 8] \text{ GHz and } Z_{res} = \sqrt{\frac{L}{C}} \approx 2 \text{ k}\Omega \rightarrow \begin{cases} \text{High inductance : } L \sim 60 \text{ nH} \\ \text{Low capacitance : } C \sim 15 \text{ fF} \end{cases}$$



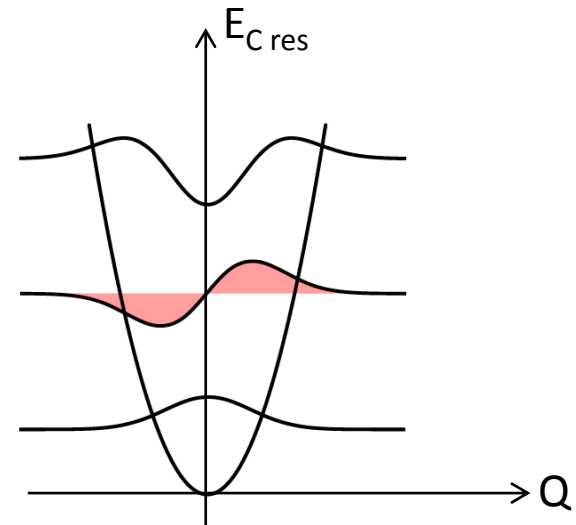
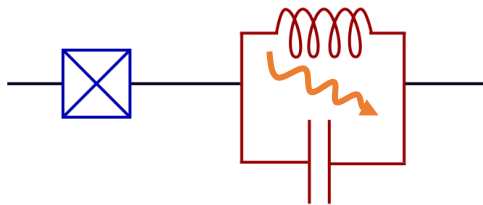
$$E_j(\Phi) = E_j^0 \left| \cos(\pi \Phi / \Phi_0) \right|$$

Large coupling achieved : $r=1.1$

Resonator back-action?

$$\Gamma^{2e}(2eV = hf_0) \propto |\langle n+1 | e^{-2i\phi_J} | n \rangle|^2$$

successive tunnel events

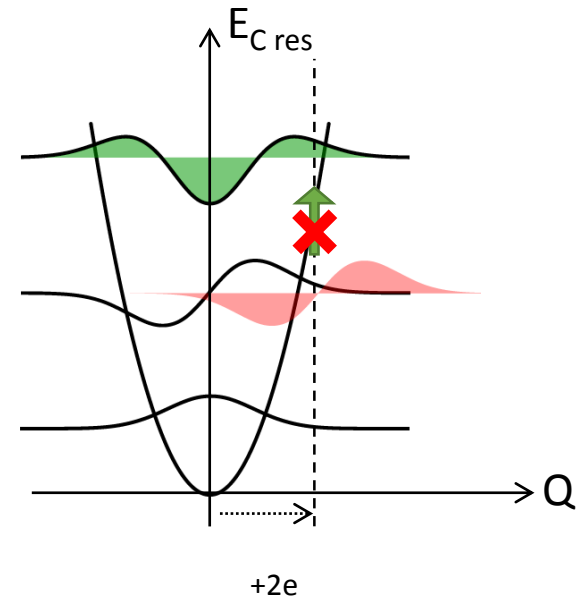
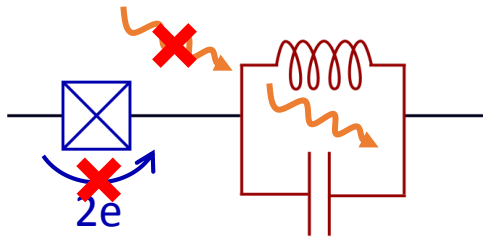


V. Gramich *et al*, Phys.Rev.Lett. **111**, 247002, 2013

Anti-bunching

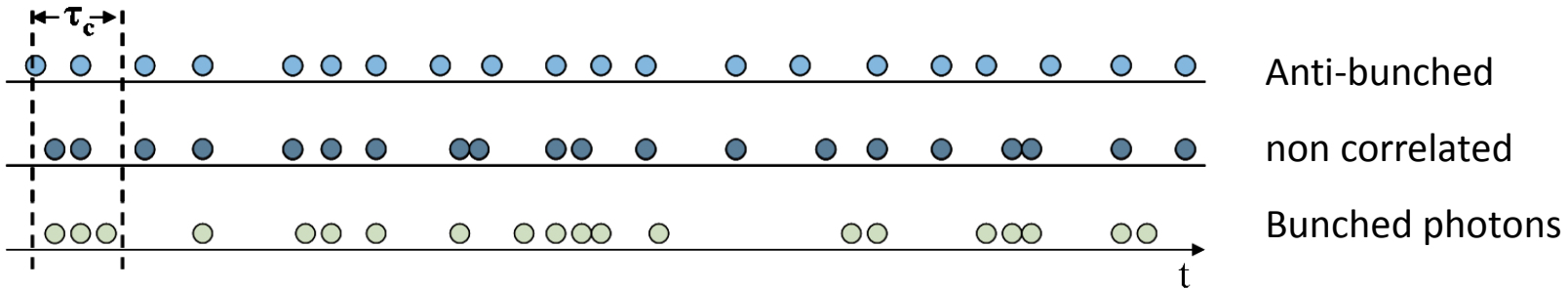
$$\Gamma^{2e}(2eV = hf_0) \propto |\langle 2 | e^{-2i\varphi_j} | 1 \rangle|^2 = 0 \text{ for } r=2$$

successive tunnel events



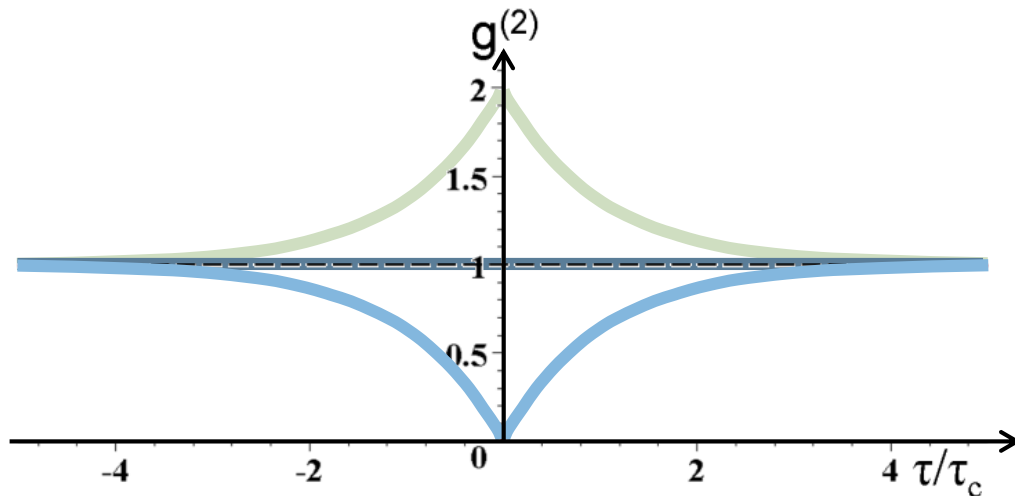
CP can't tunnel when 1 photon in the resonator !

Photon statistics



Statistic of photon emission: $g^{(2)}$ function

$$g^{(2)}(\tau) = \frac{\langle a^\dagger(t)a^\dagger(t+\tau)a(t+\tau)a(t) \rangle}{\langle a^\dagger a \rangle^2}$$



τ_c , averaged photon lifetime in the resonator

Theoretical predictions:

$$g^{(2)}(0) = \left(1 - \frac{r}{2}\right)^2$$

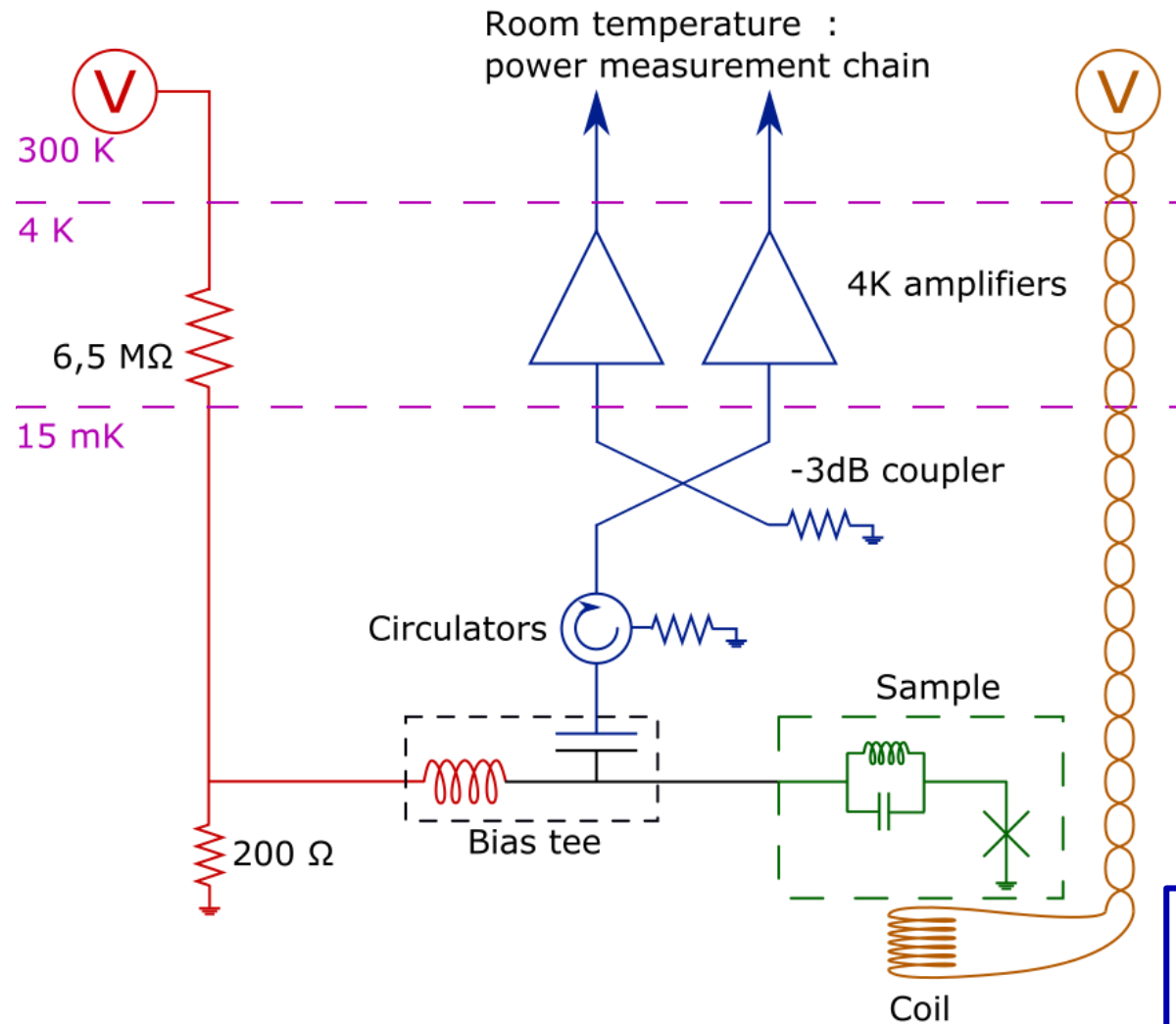
our sample: $r = 1.1 \rightarrow g^{(2)}(0) \approx 0.22$

V. Gramich *et al*, Phys.Rev.Lett. **111**, 247002, 2013

Measurement :

$$g^{(2)}(\tau) = 1 + \frac{\langle \delta P(t)\delta P(t+\tau) \rangle}{P^2}$$

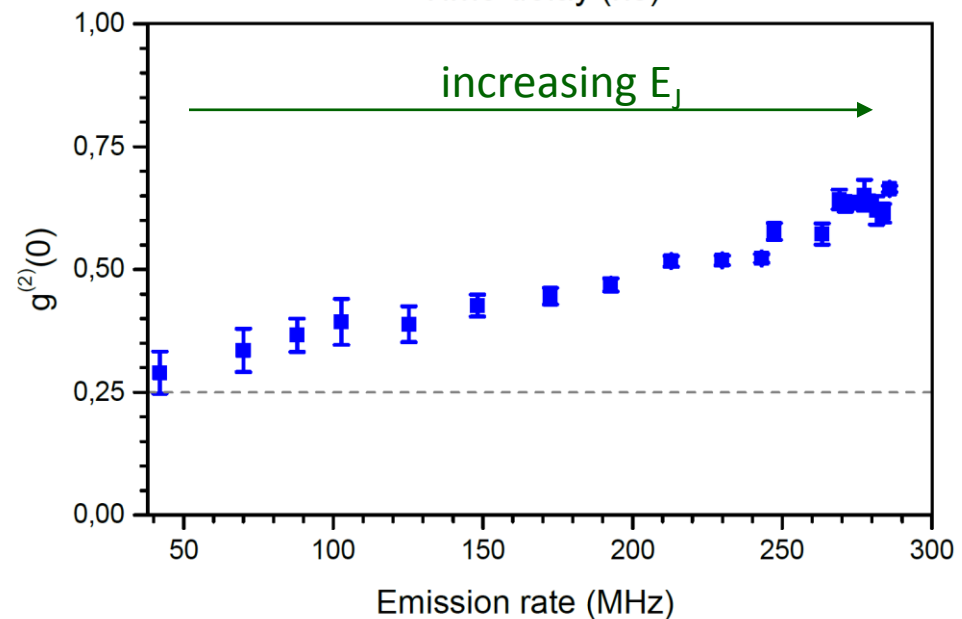
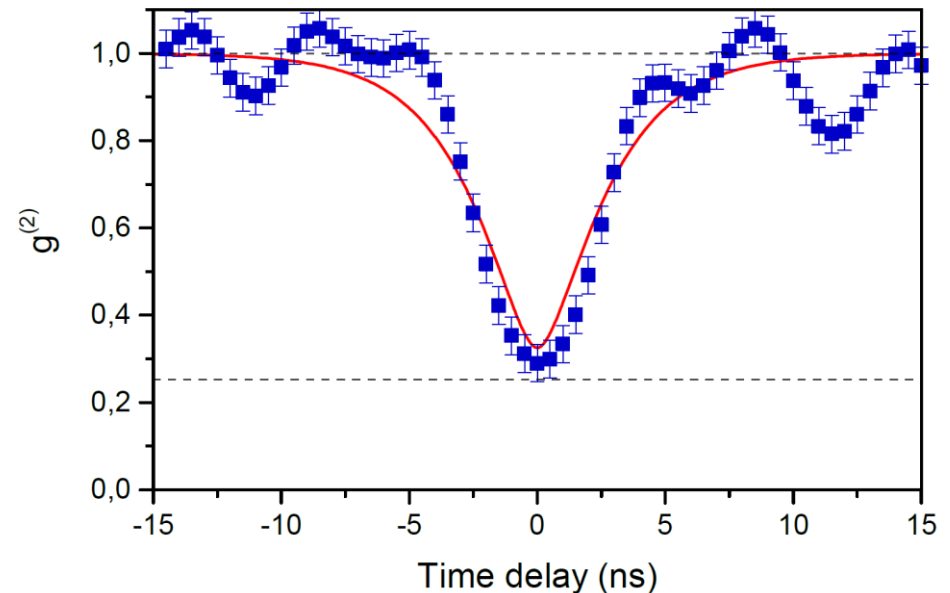
HBT measurement setup



$$g^{(2)}(\tau) = 1 + \frac{\langle \delta P(t) \delta P(t + \tau) \rangle}{P^2}$$

Second order correlator results $g^{(2)}(\tau)$

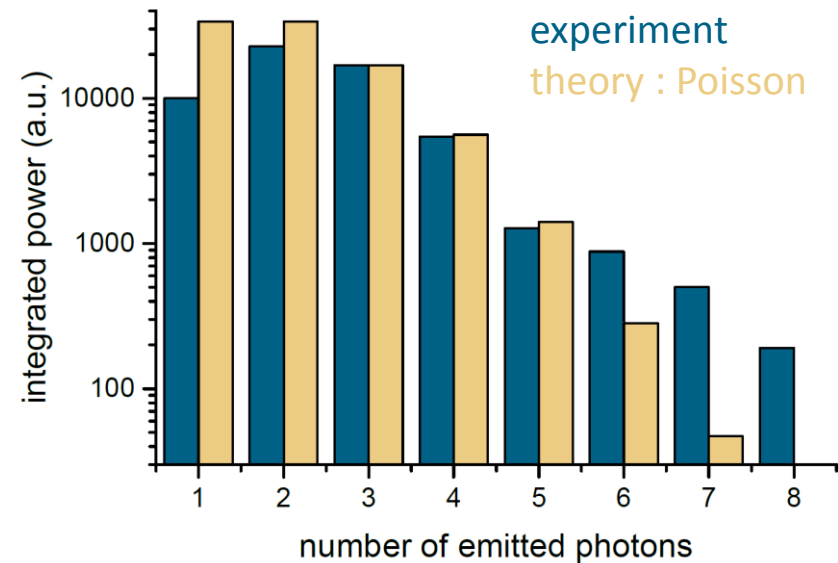
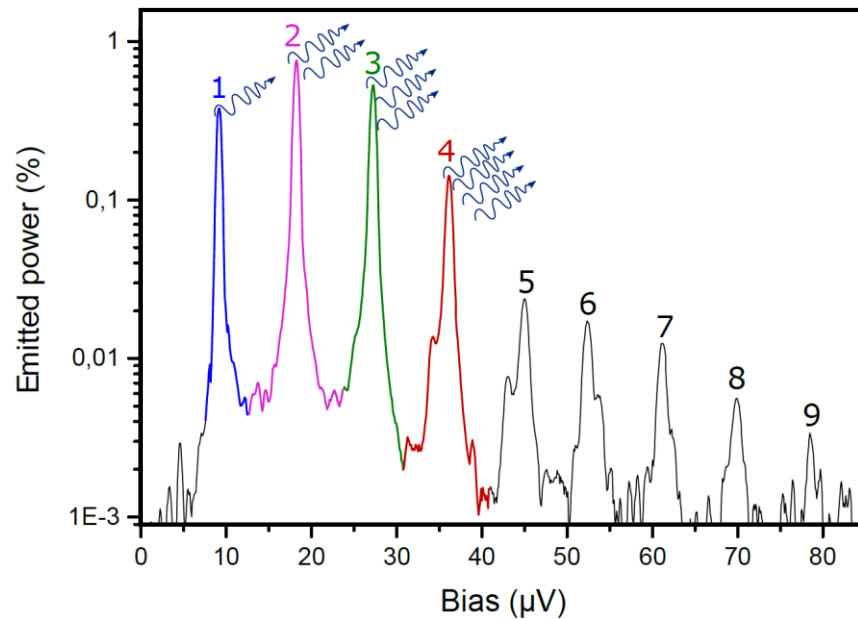
- Anti-bunched photon emission
- Time scale of correlations : photon lifetime in the cavity
- The correlation decreases with mean population of the resonator



Multiple photon emission

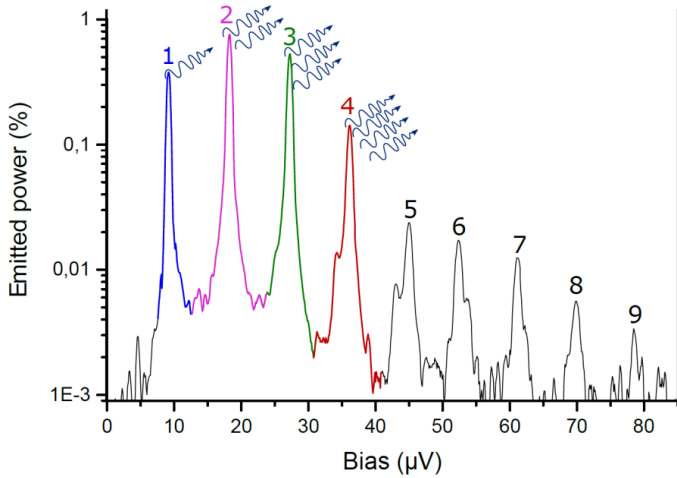
- We can now observe up to 9-photon processes :

$$\Gamma^{ph}(2eV = khf_0) = E_J^2 \frac{r^k}{(k-1)!}$$

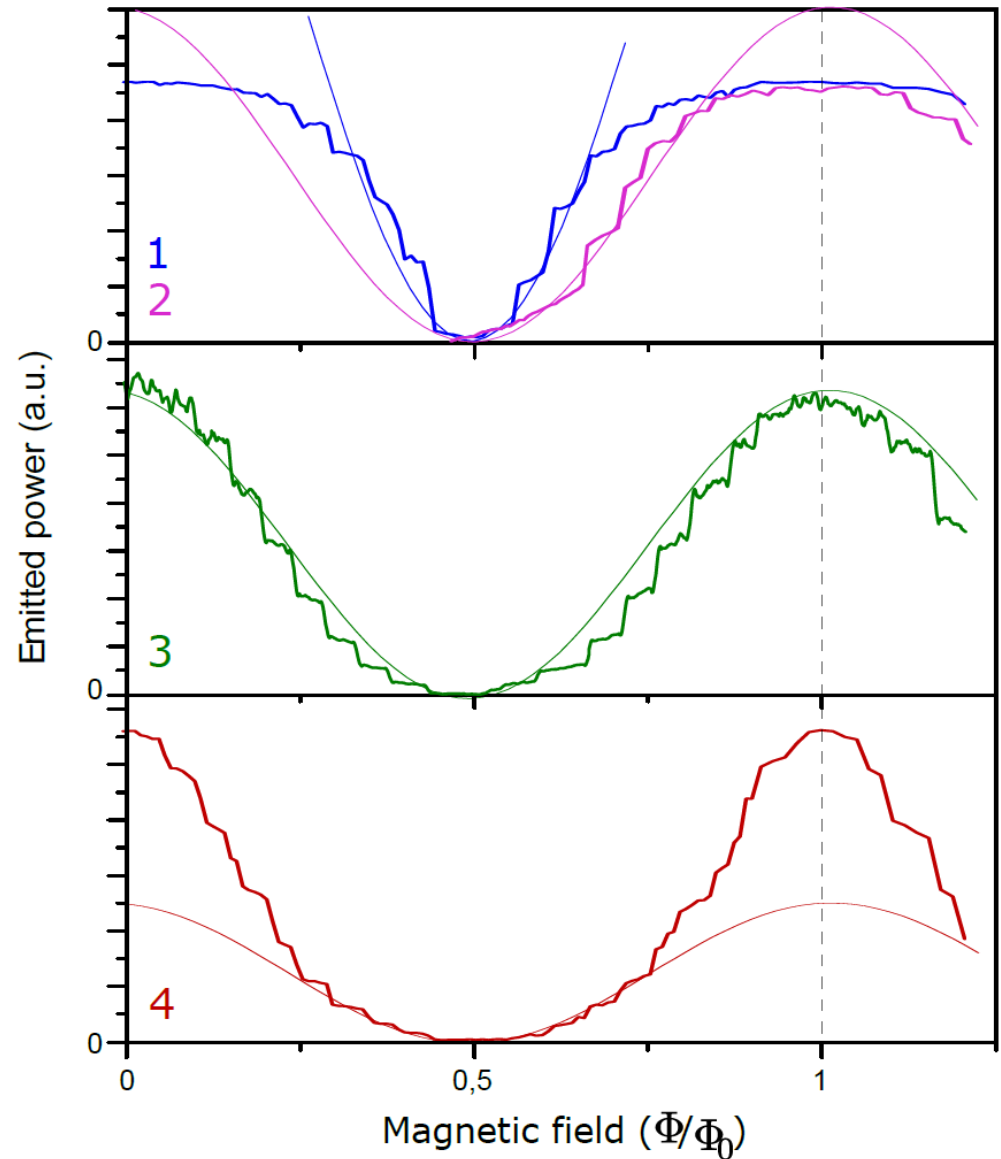


Does the emission follow theoretical predictions ?

Multiple photon emission



$$\Gamma^{ph} \propto E_J^2 \cos^2(\Phi)$$



Conclusion and perspectives

Summary

dc-biased Josephson junction:
simple, compact and bright source of non-classical photons

What's next ?

- Higher coupling : perfect single photon source ?
- Transition to parametric resonance (photon pair bursts)
- Probe for other quantum devices ?

Thanks : Marc Westig, Iouri Moukharski, Denis Vion, Philippe Joyez, Carles Altimiras, Patrice Roche, Daniel Estève

Thanks !

