Shot-Noise and Symmetry of the Kondo Effect

Meydi Ferrier

Université Paris sud, Orsay

R. Delagrange, R. Deblock, H. Bouchiat

Osaka University

T. Arakawa, T. Hata, R. Fujiwara, K. Kobayashi

Theory

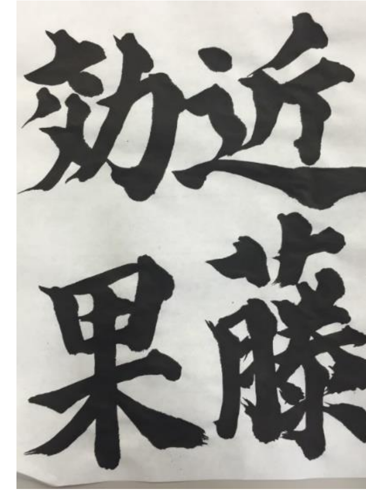
R. Sakano, ISSP, Tokyo

Y. Teratani, A. Oguri, Osaka city University

Quick reminder

Shot noise

Kondo effect in CNT



Linear shot noise of SU(2) and SU(4) Kondo effect

Direct signature of the symmetry class

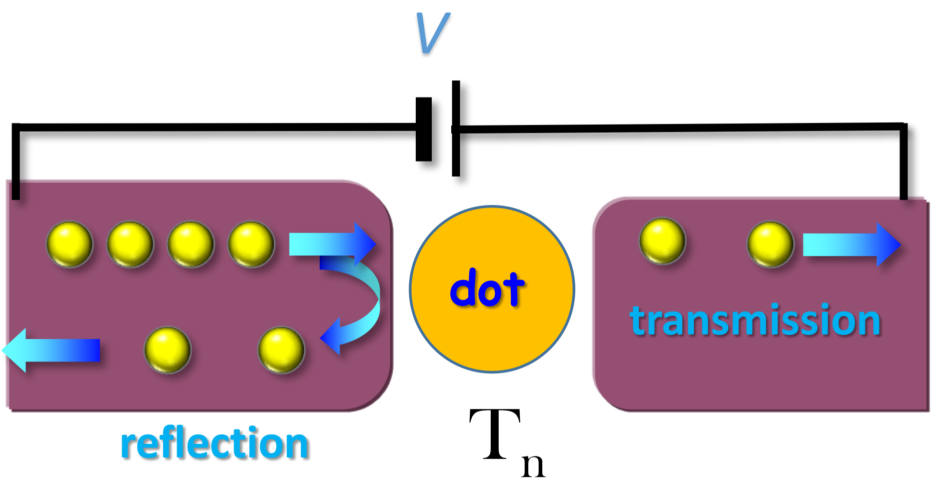
Non-linear noise:

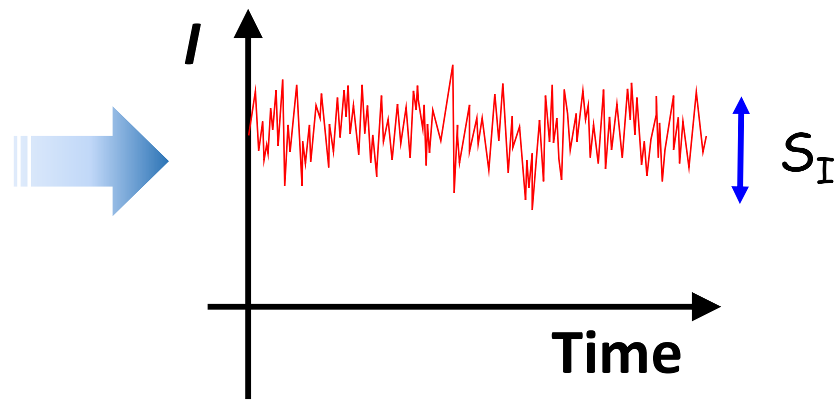
Observation of 2-particle scattering induced by interactions out of equilibrium

SU(4) to SU(2) crossover in magnetic field

Continuous evolution of effective charge e*

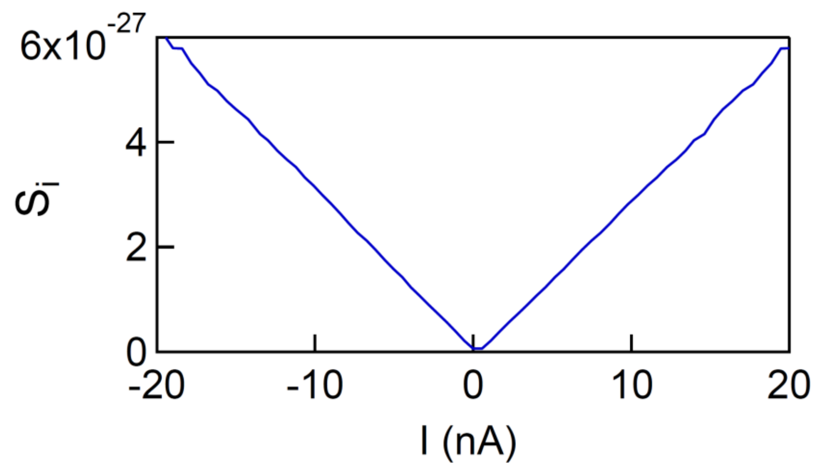
Origin of shot noise





Fluctuations due to the partition of scattered particles

$$S_I = 2eF\,I \quad \text{with} \quad F = \frac{\sum T_n(1-T_n)}{\sum T_n}$$

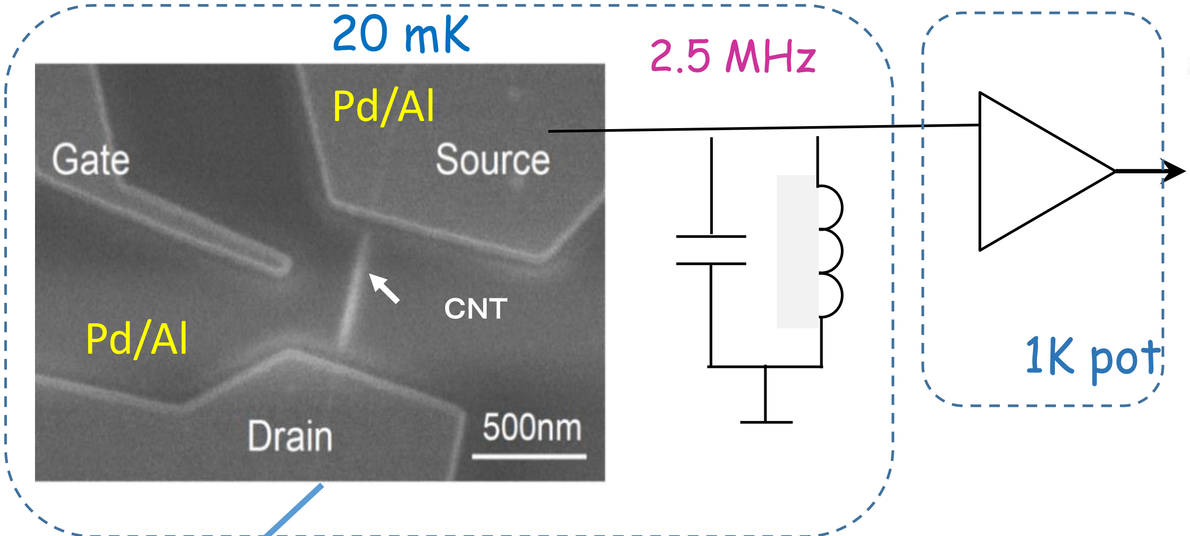


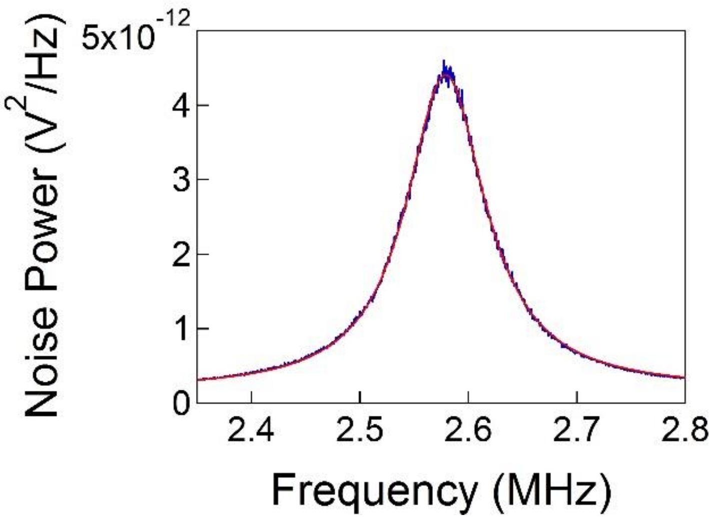
e = charge of the particle

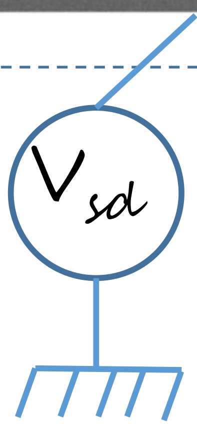
F = Fano factor = statistics of the particles

Nature of the quasiparticle and the scattering mechanism

Experimental set-up for noise measurement





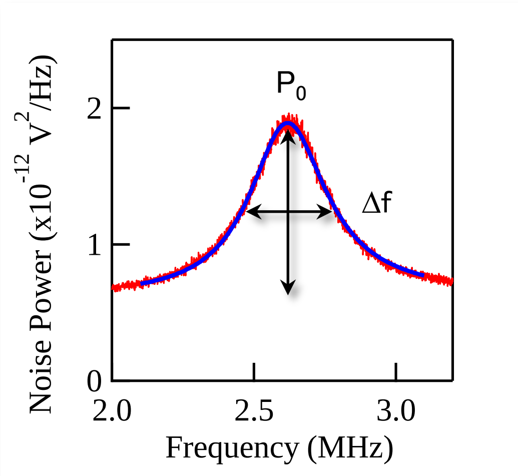


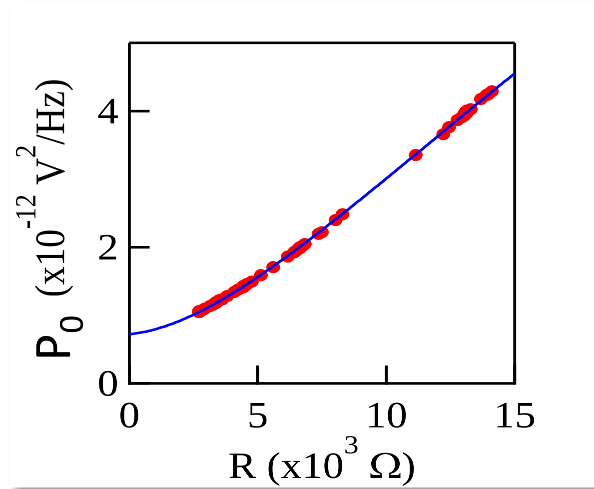
DC measurement +
resonant circuit for noise

Sensitivity ~ 10^{-29} A^2/Hz

Hashisaka, KK *et al.* Rev. Sci. Inst. 80, 096105 (2009);
Arakawa, KK, et al. Appl. Phys. Lett. 103, 172104 (2013).

Calibration and measurement of the Noise

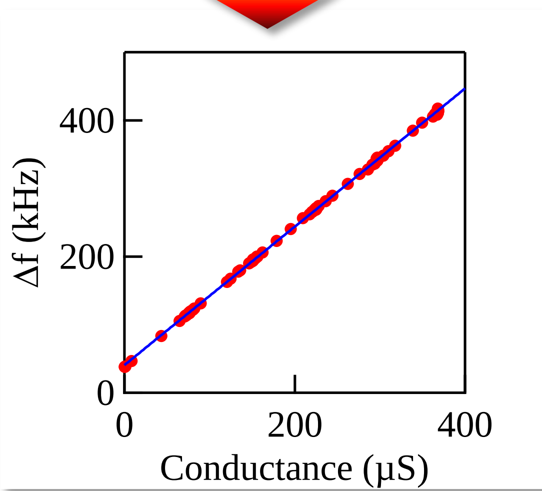




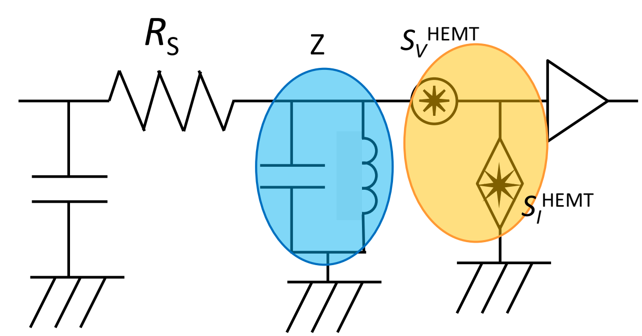
Z : LC circuit Impedance

R : Sample Resistance

$$P_0 = A\left\{\left(\frac{4k_BT}{R} + S_I^{(HEMT)}\right)\left(\frac{ZR}{Z+R}\right)^2 + S_V^{(HEMT)}\right\}$$



$$\Delta f = \frac{1}{2\pi C}\left(G + \frac{1}{Z}\right)$$

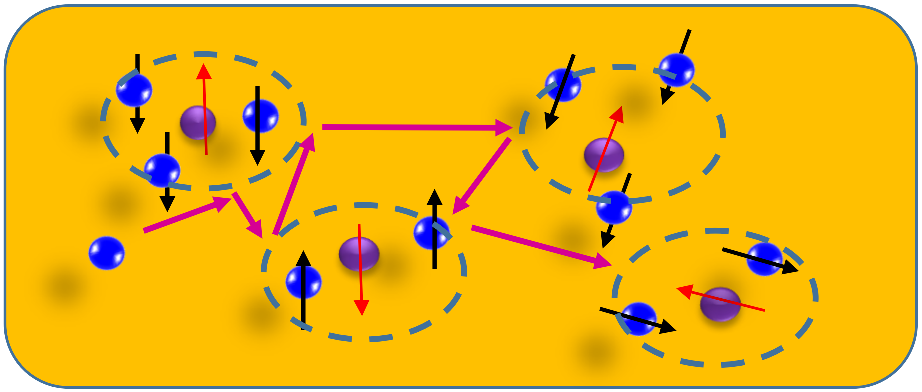


L. DiCarlo *et al.*, Rev. Sci. Instrum. **77**, 073906 (2006).

Signature of the Kondo effect

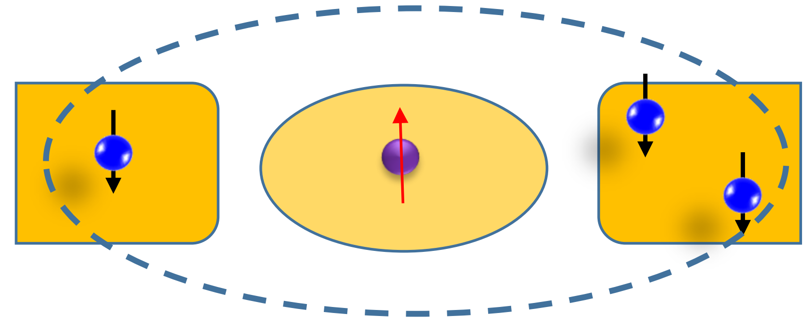
→ Spin screening induced by electron interaction

Macroscopic sample

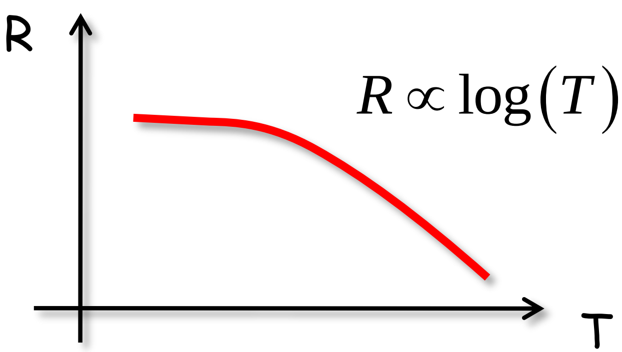


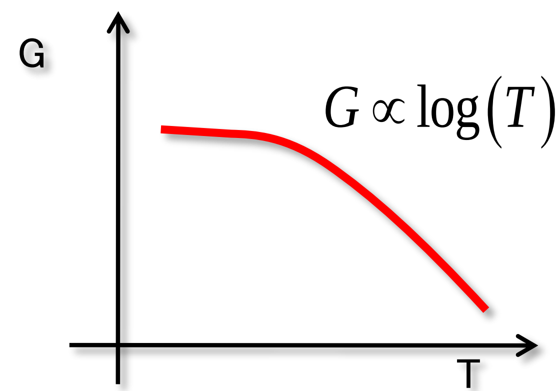
Scattering enhanced

Quantum dot (1 electron)



Delocalization enhanced





Probe locally a many body state

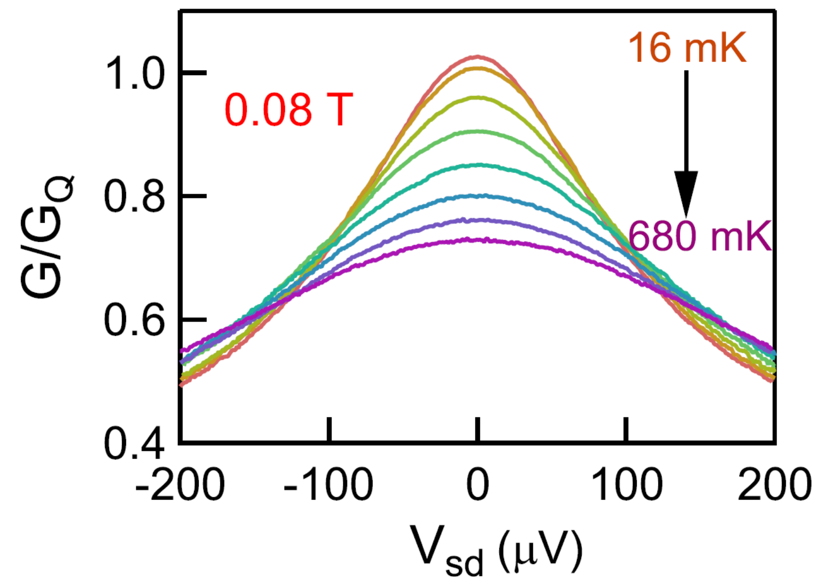
Interaction and Non-equilibrium Kondo physics

« First order » correction:

Non-linear conductance

Scaling with T,B,V

Kretinin et al, Phys. Rev. B 84 (2011)



Higher order : 2-particle scattering induced by residual interaction

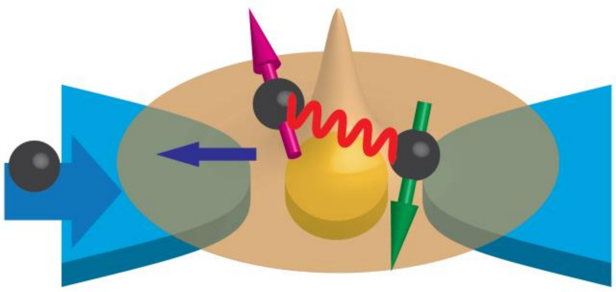
Non-linear noise   Enhanced current fluctuations

Zarchin et al, PRB (2008)

Delattre et al, Nature Physics (2009)

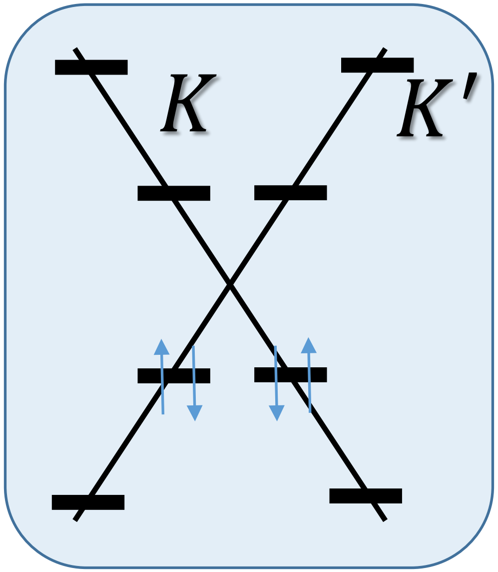
Yamauchi et al, PRL (2011)

not quantitative

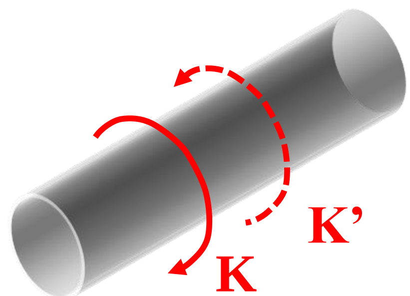


Nanotube dot = 2 different Kondo states

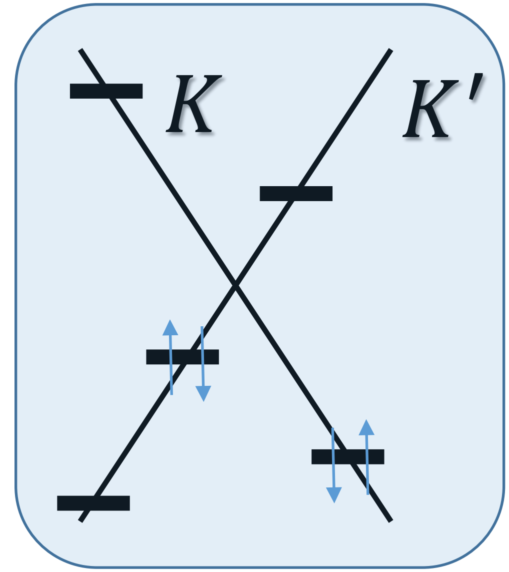
nanotube band structure



4 e$^-$ per shell



Disorder, spin-orbit = splitting



2 e$^-$ per shell

2 « spin » are screened

2 transport channels

SU(4) symmetry

Only the usual spin is screened

1 transport channel

SU(2) symmetry

Signature of interaction depends on the symmetry class

Part 1

Noise in the linear regime

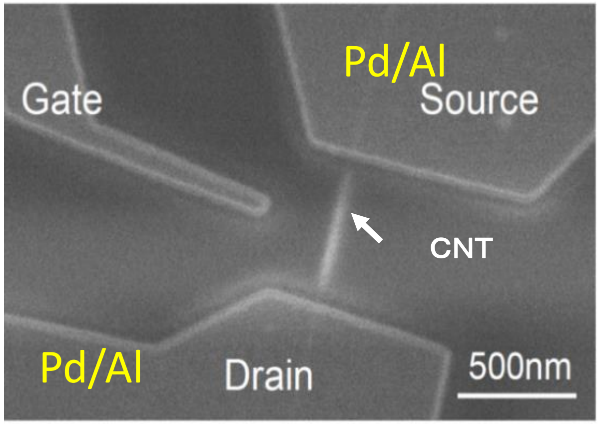
$$eV \ll k_B T_K$$

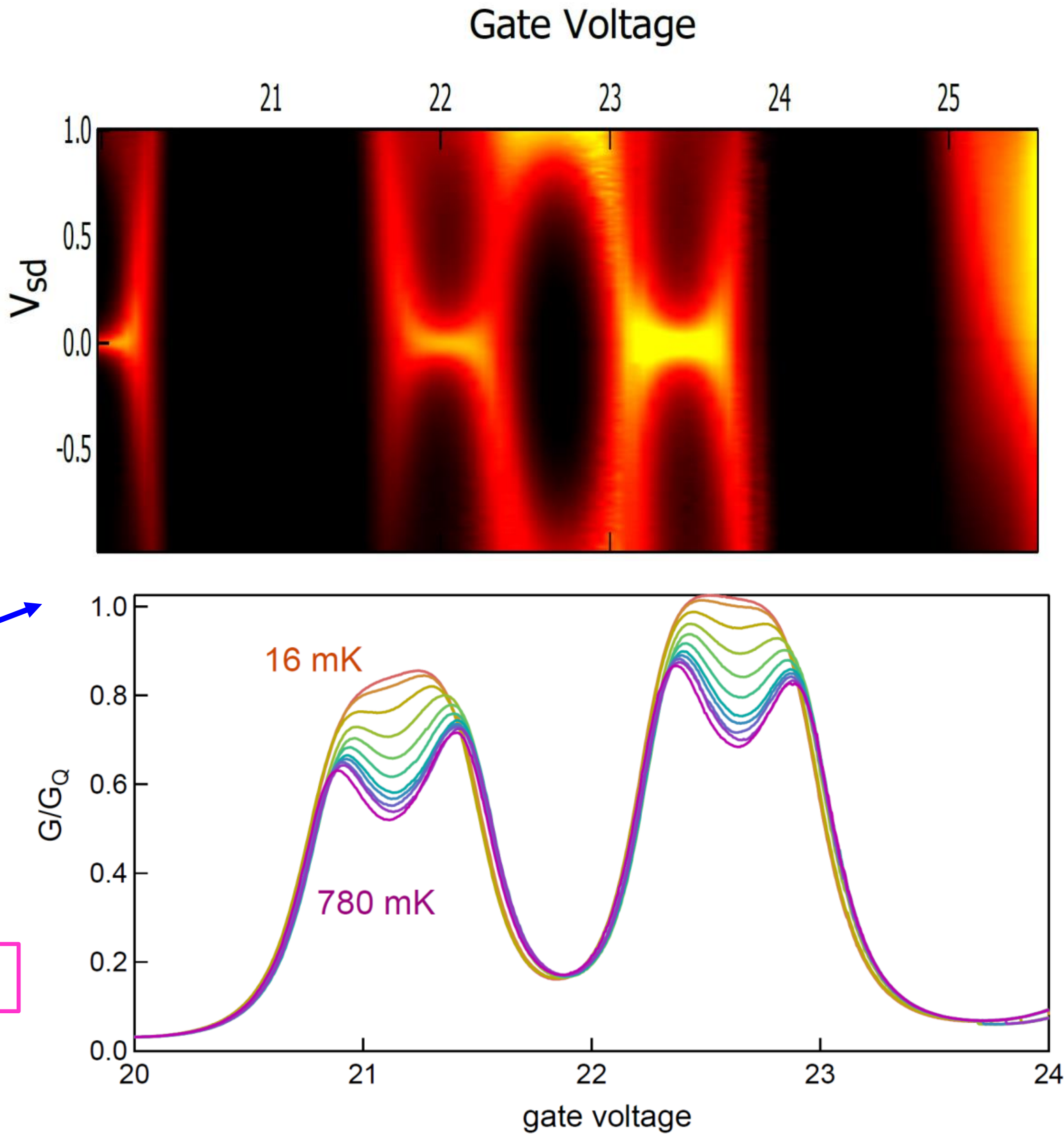
➡ Kondo state = Fermi-liquid constituted of non-interacting quasi-particles

➡ Signature of the symmetry class of the Fermi-liquid

SU(2) or SU(4)

Carbon Nanotube in the SU(2) Kondo state



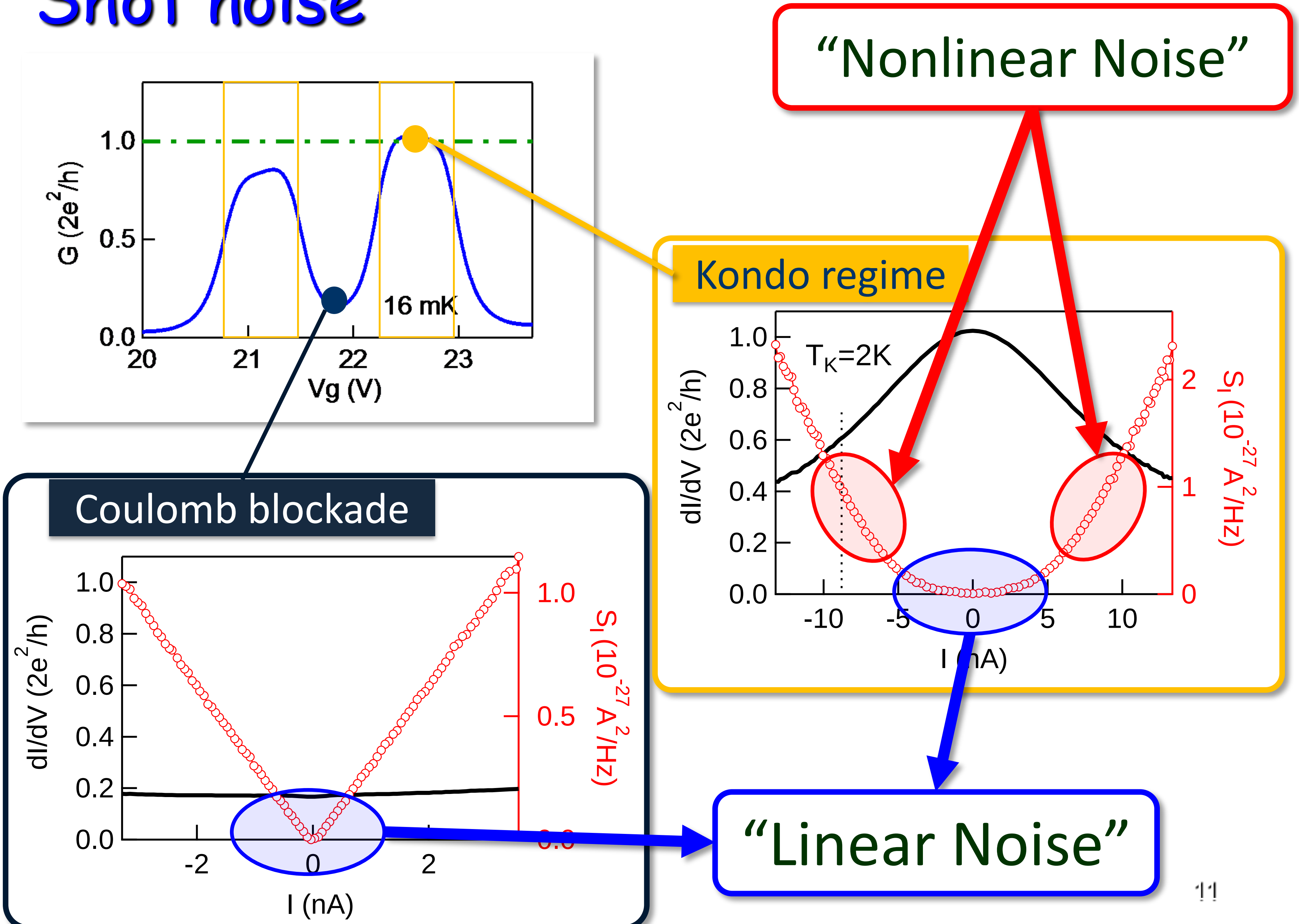


$$G = \frac{2e^2}{h}$$

B=0.08T to destroy superconductivity

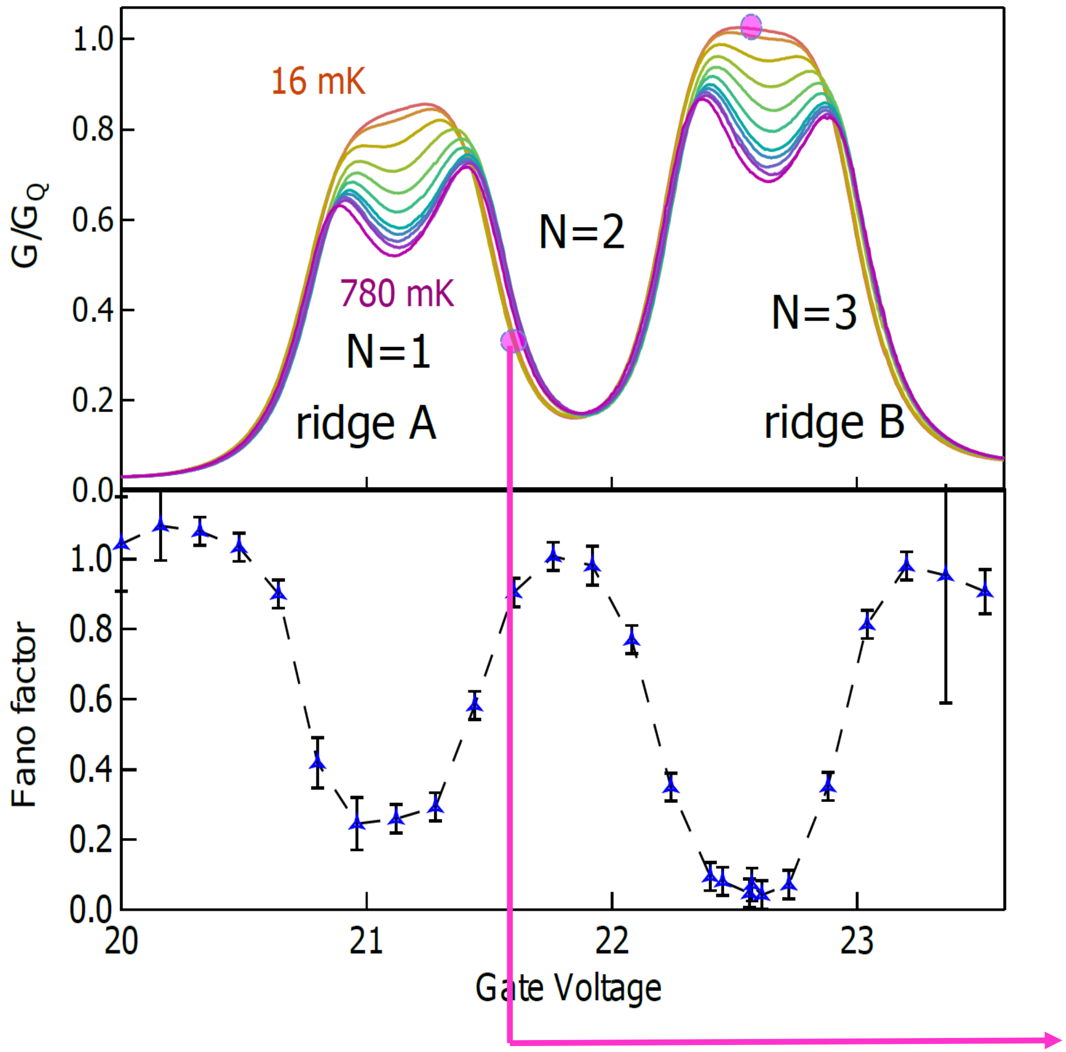
Unitary limit reached at low T

Shot noise



"Nonlinear Noise"
1.0
0.5
0.0
G (2e2/h)
20
21
22
23
Vg (V)
16 mK
Kondo regime
TK=2K
dI/dV (2e2/h)
SI (10-27 A2/Hz)
I (nA)
Coulomb blockade
"Linear Noise"

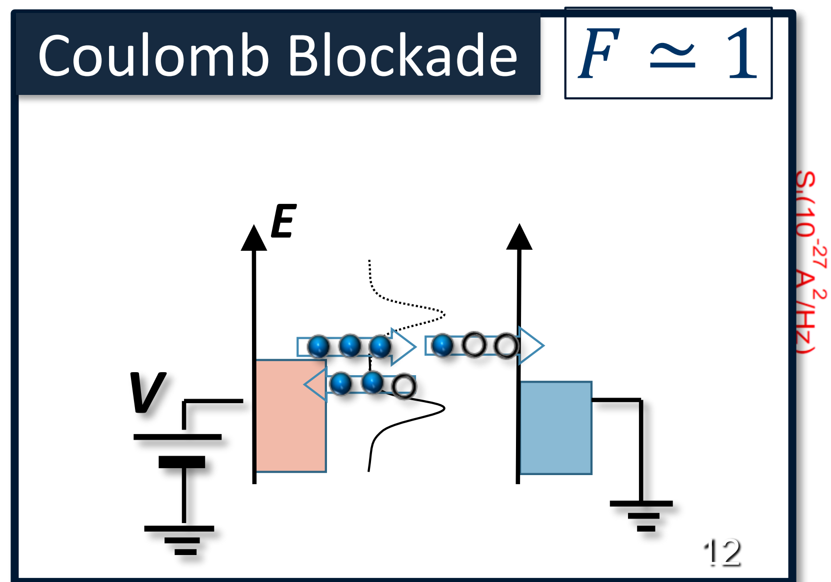
Linear Noise ($eV \ll T_K$) in the Coulomb Valley



No interaction:

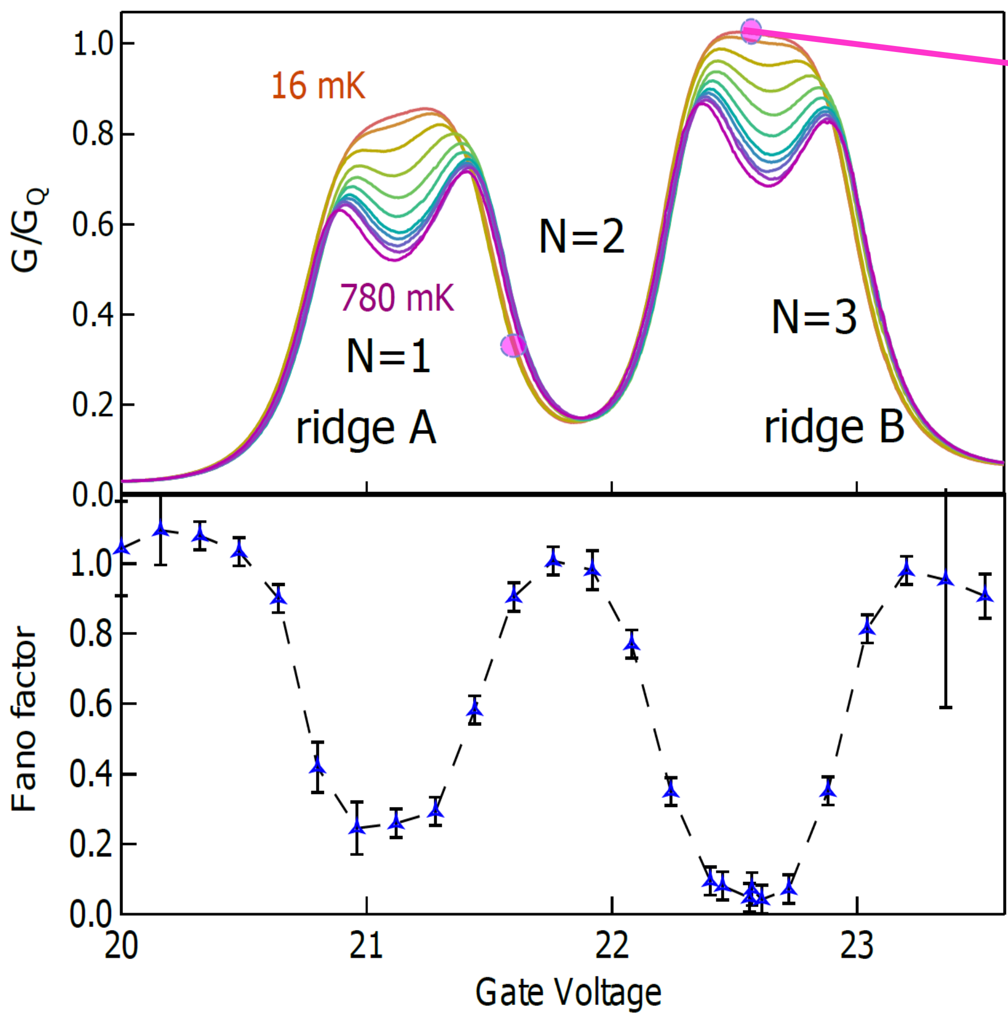
$$G = G_Q T$$

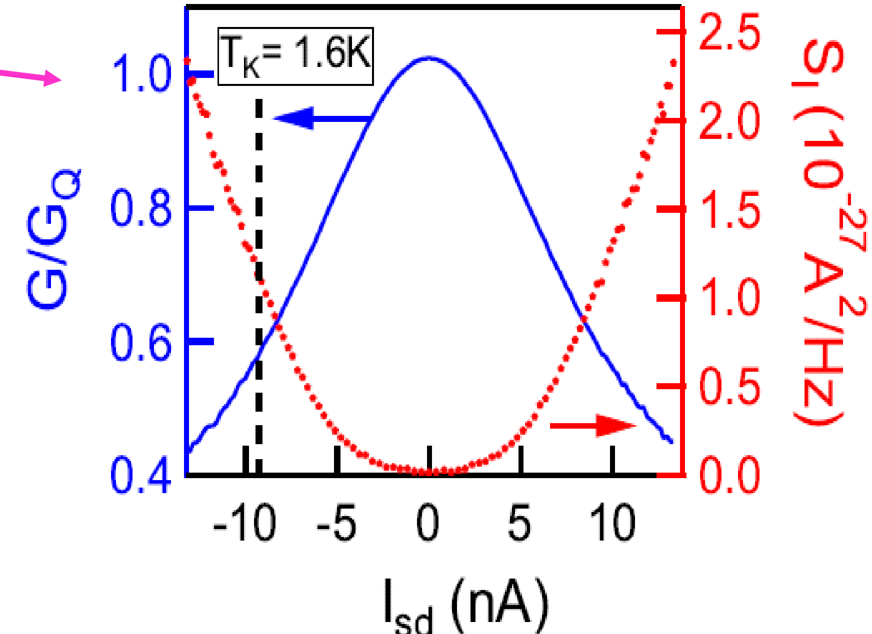
$$F = 1 - T$$



Conventional Poissonian tunneling

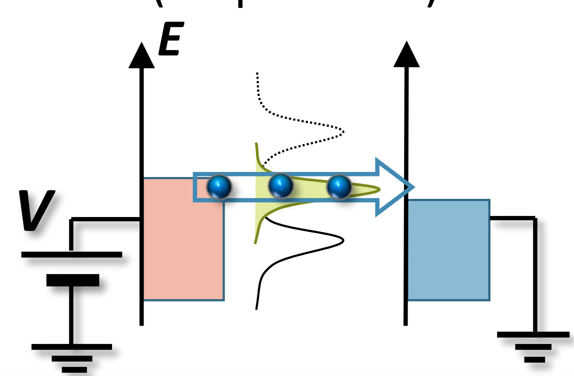
Linear Noise (eV « T_K) on the Kondo ridge





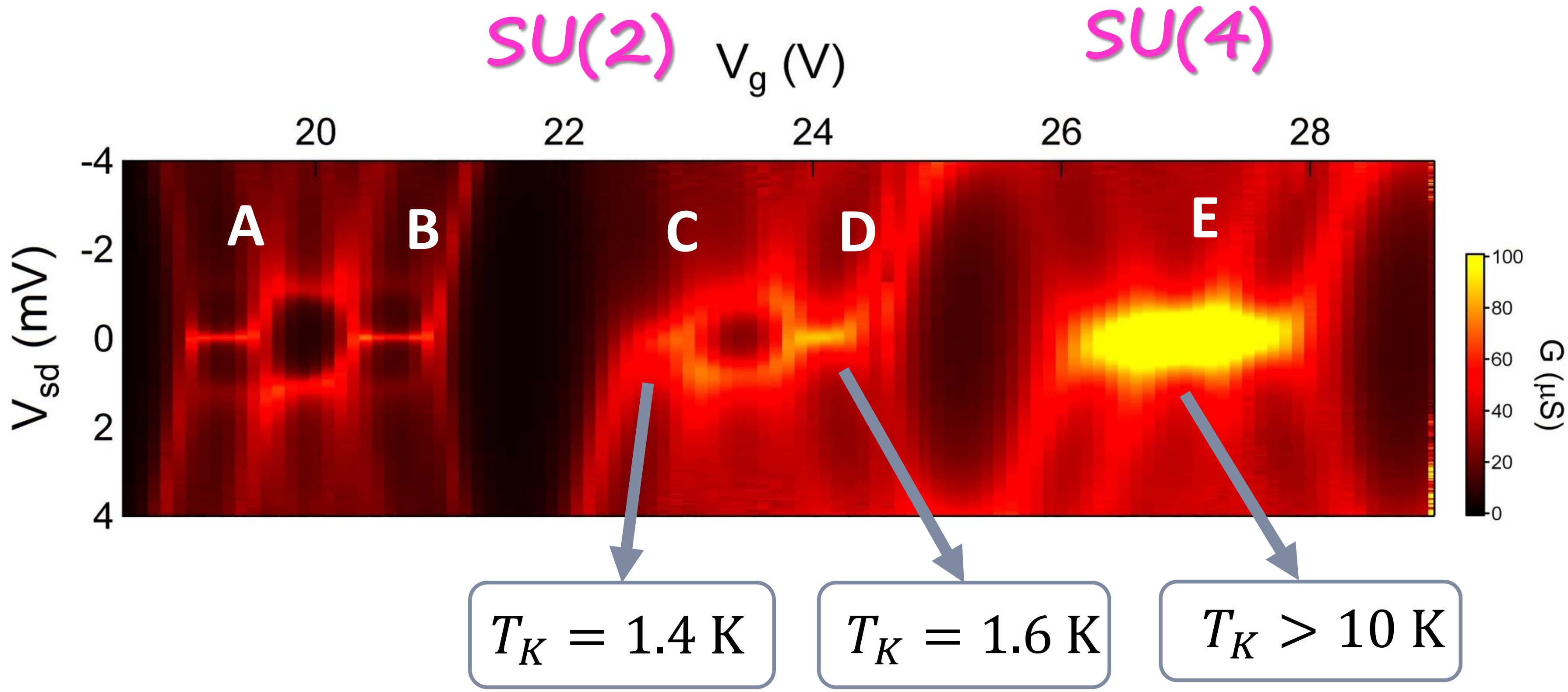
Kondo regime $F \ll 1$

Perfect transmission through resonance (no partition)

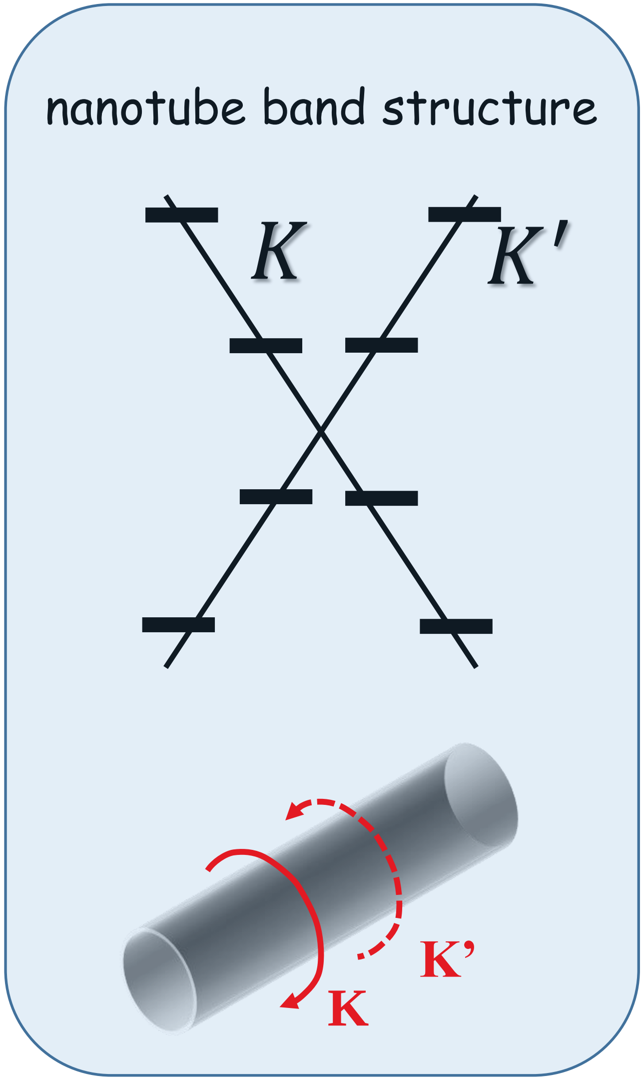


silent impurity around V=0

Different Kondo effect in the same Nanotube



Screening of spin & orbit degrees of freedom

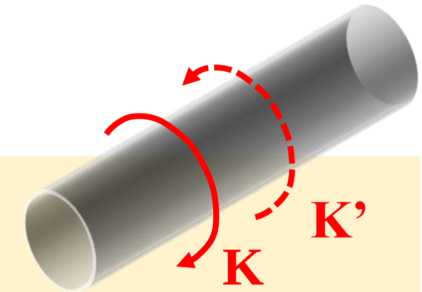


2 channels participates for transport

Kondo screening for odd and even number of electrons

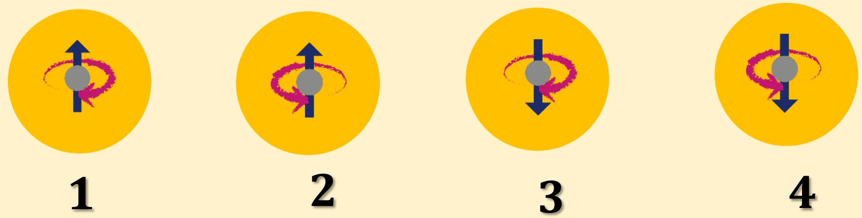
N=1

SU(4) Kondo state in CNT



2 quantum numbers : K, σ

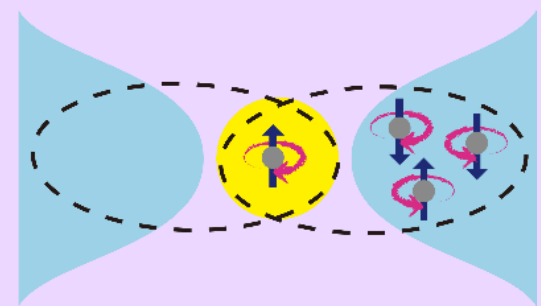
isolated ground states degenerate



Interaction + leads

Kondo singlet ground state

$$G = 2\frac{e^2}{h}$$



Ψ

Screening of the total magnetic moment

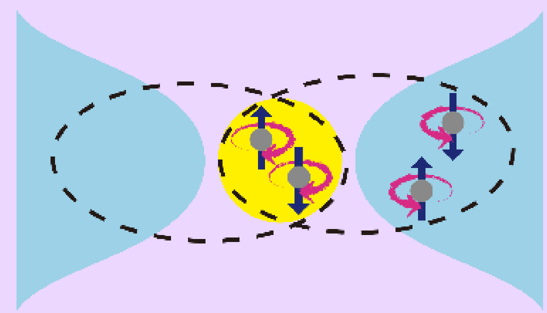
SU(4) Kondo effect at N=2

Degenerate isolated ground states

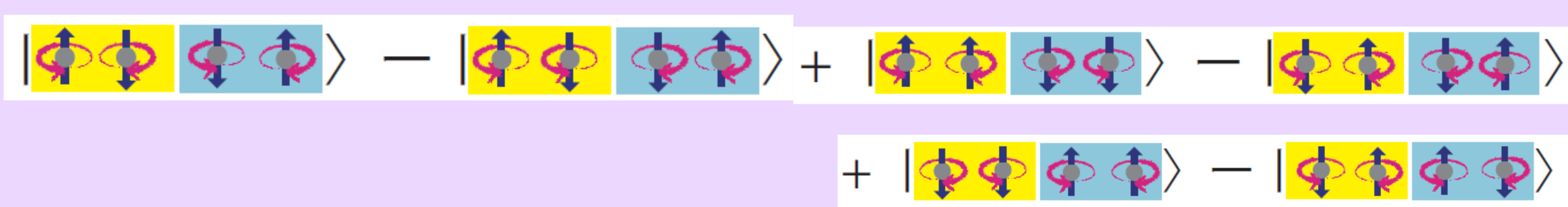
Interaction + leads

Kondo singlet ground state

$$G = 4\frac{e^2}{h}$$

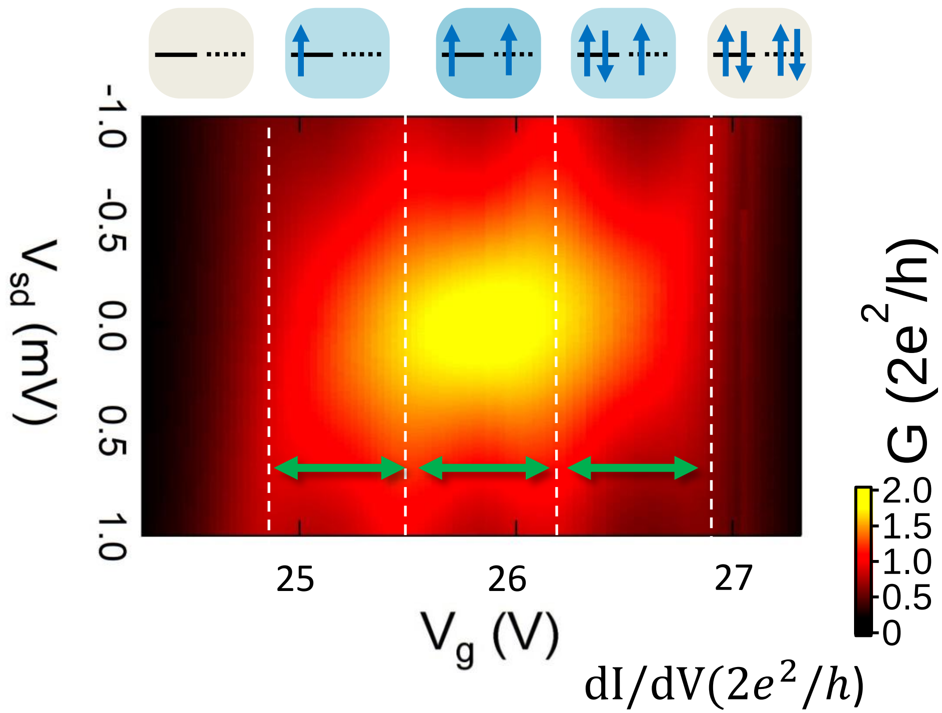


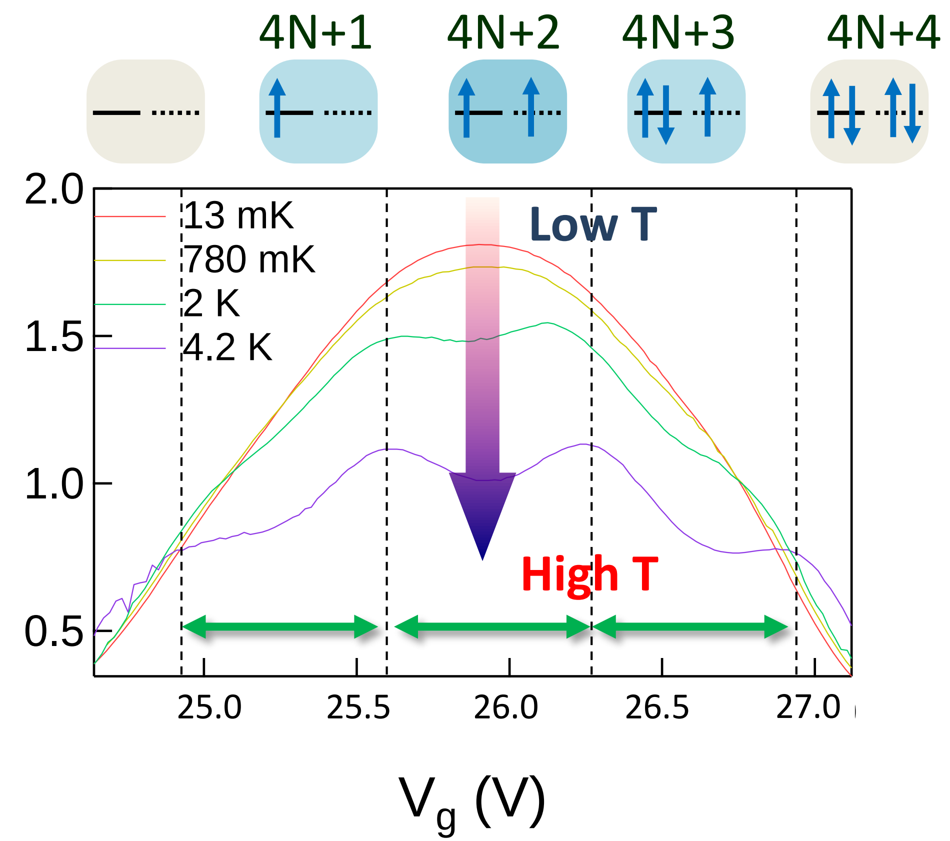
$\Psi =$



The SU(4) Kondo state

Coulomb diamond



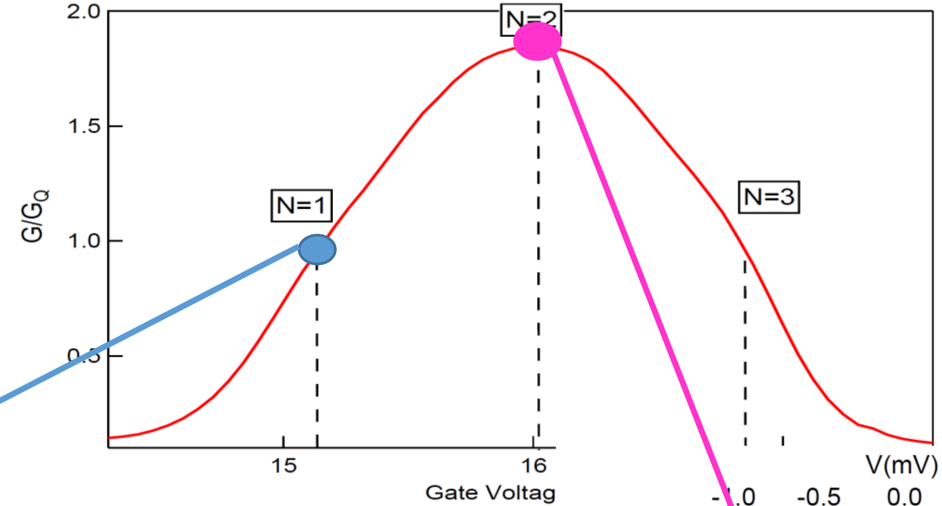


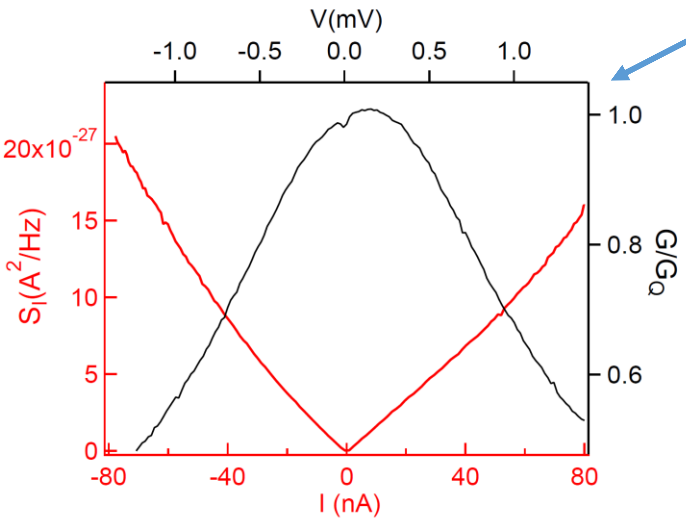
Kondo resonance for N=1,2 and 3 electrons

Shot noise and SU(4) Kondo effect

$$G = G_Q(T_1 + T_2)$$

$$F = \frac{\sum T_n(1-T_n)}{\sum T_n}$$

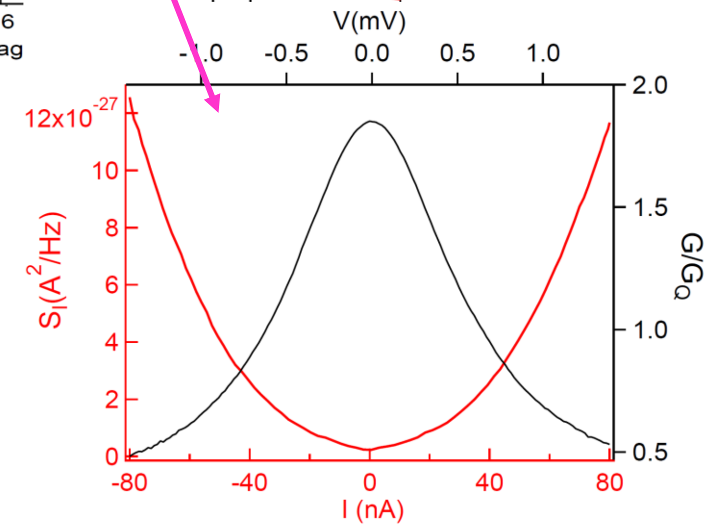




2 channels with T= 1/2

$G = G_Q \qquad F \approx 0.5$

Delattre et al, Nature Physics (2009)



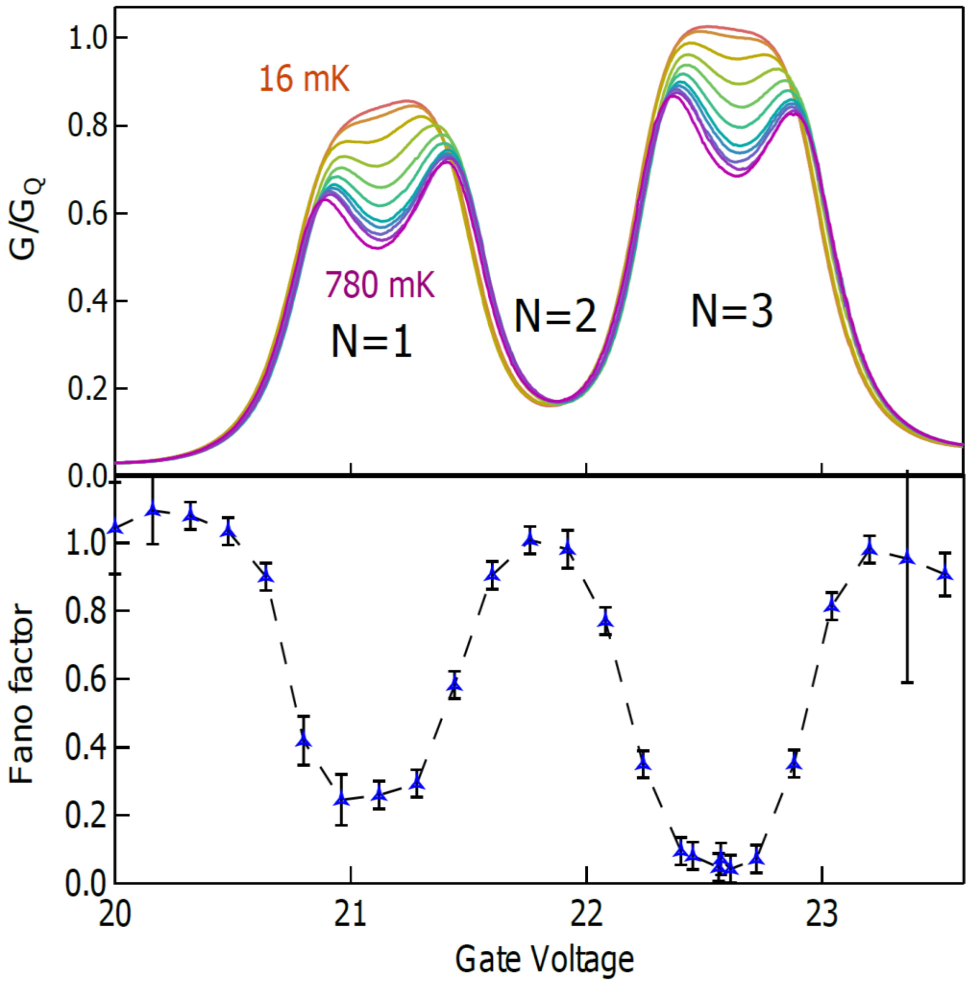
2 perfect channels : Noise ≈ 0

$G = 1.82\, G_Q \qquad F = 0.15$

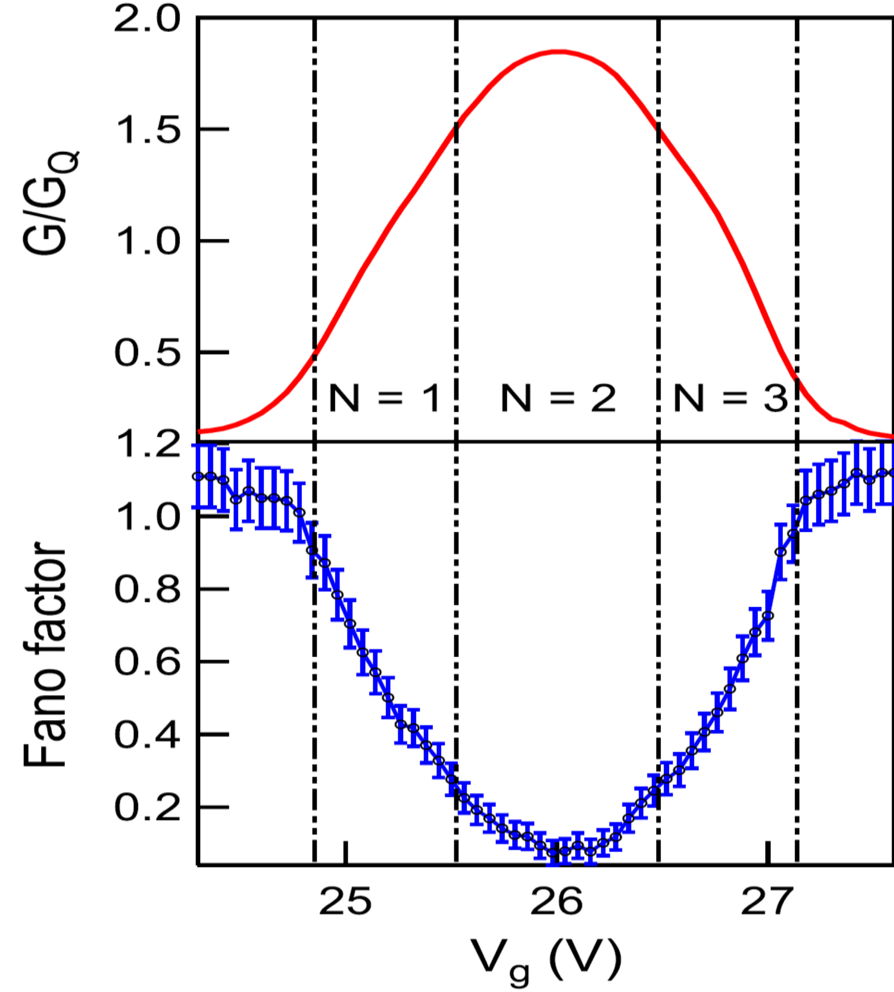
Odd and Even SU(4) Kondo effect unambigously observed

Summary for the linear shot noise

SU(2)

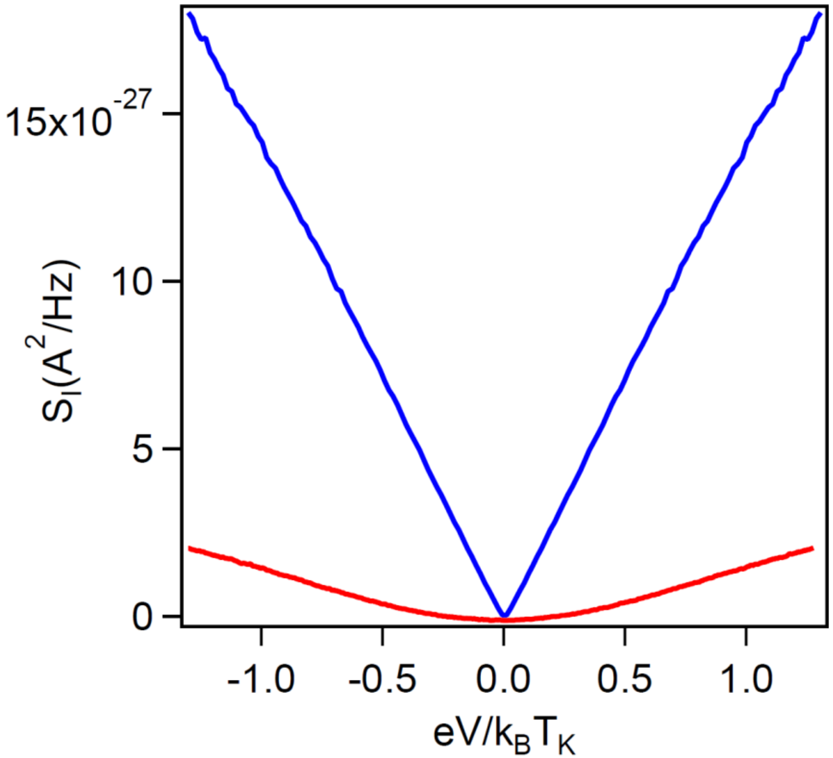


SU(4)



Symmetry distinguished by the linear shot noise

Noise contains more information than conductance !



SU(2) odd

$$G = G_Q \times 1$$

SU(4) odd

$$G = G_Q \times \left(\frac{1}{2} + \frac{1}{2}\right)$$

SU(2) = 1 perfect channel = NO NOISE

$F = 1 - 1$

SU(4) = 2 channels with T = 1/2 ⇒ strong partition = strong shot noise

Very difficult to distinguish experimentally

$F = 2 \times \left(1 - \frac{1}{2}\right)$

Scattering is fundamentally different

Linear noise completely described by non-interacting quasi-particles

↓

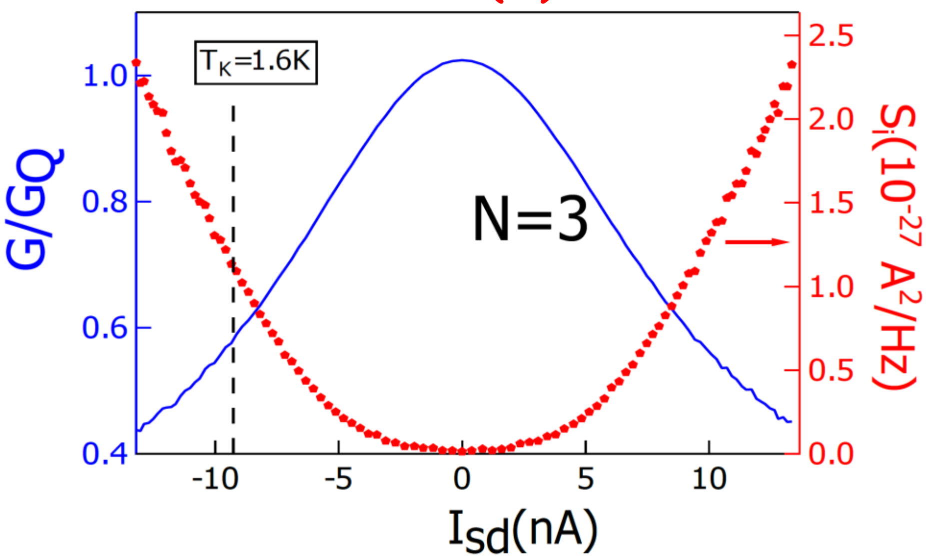
Part 2

What about non-linear Noise?

Observation of 2-quasi-particle scattering induced by interaction

Is non-linear noise only due to non-linear conductance?

SU(2)



Kondo effect : Transmission depends strongly on energy

$$\frac{dI}{dV} = G_Q T(V)$$

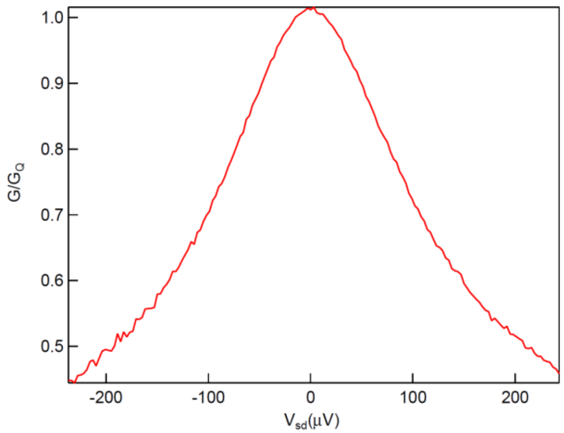
Without interaction non-linearities appear in noise:

$$S(V) = 2G_Q \int_0^{eV} T(\epsilon)\left(1 - T(\epsilon)\right) d\epsilon = 2\int_0^{eV} G(\epsilon)\left(1 - \frac{G}{G_Q}(\epsilon)\right) d\epsilon$$

Non-linear Fano factor for non interacting particles

$$S(V^3) = 2eF_K I(V^3)$$

Non-linear conductance :



@ low voltage

$$\begin{cases} T(\varepsilon) = 1 - \alpha\varepsilon^2 \\ I(V) = G_Q V - \dfrac{\alpha}{3} G_Q V^3 \end{cases}$$

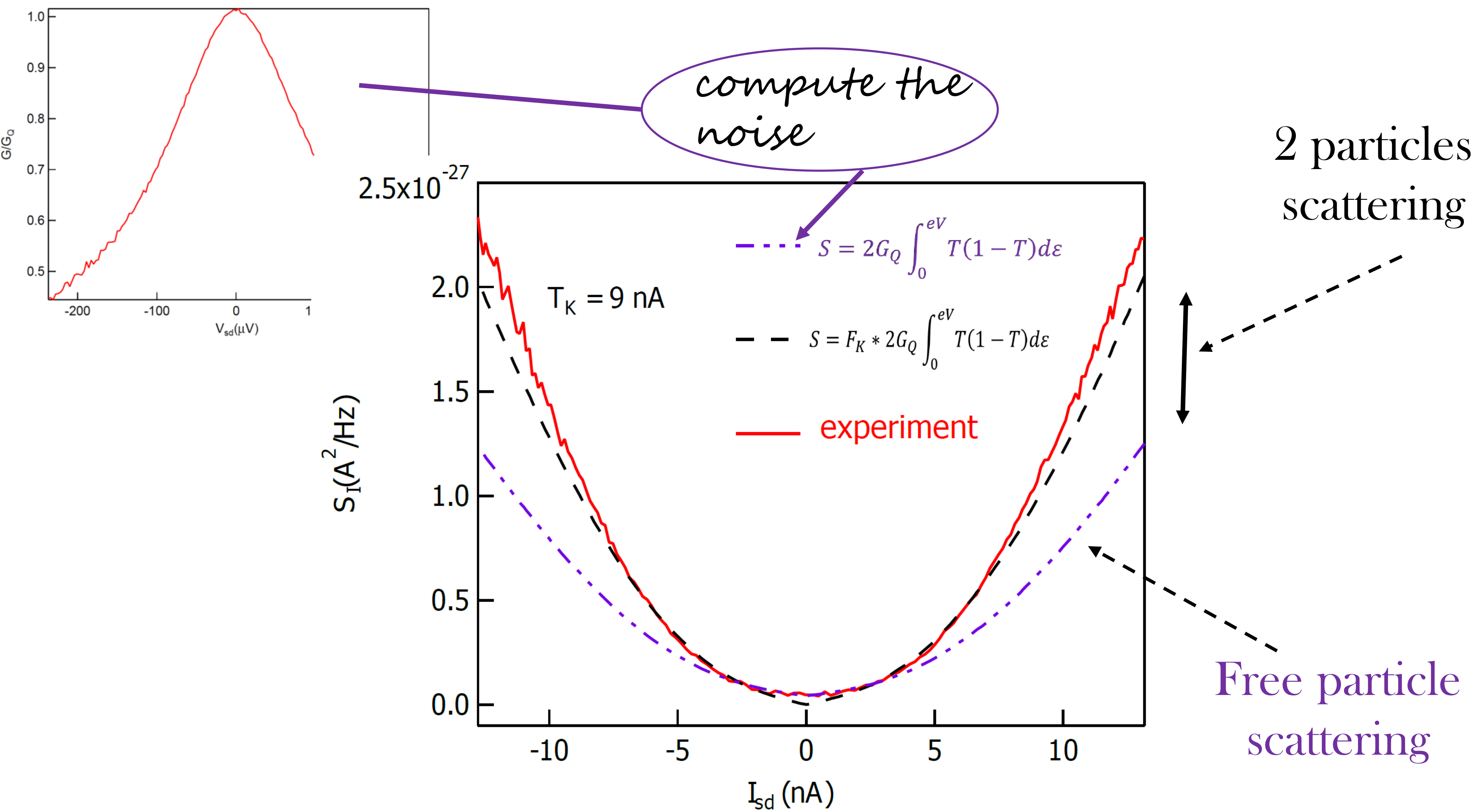
Non interacting quasi-particles picture :

$$S(V) = 2G_Q \int_0^{eV} T(\epsilon)\left(1 - T(\epsilon)\right) d\epsilon = 2e\frac{\alpha e^2 G_Q}{3} V^3$$

$F_K = 1$

F_K measures the probability for 2-particle scattering

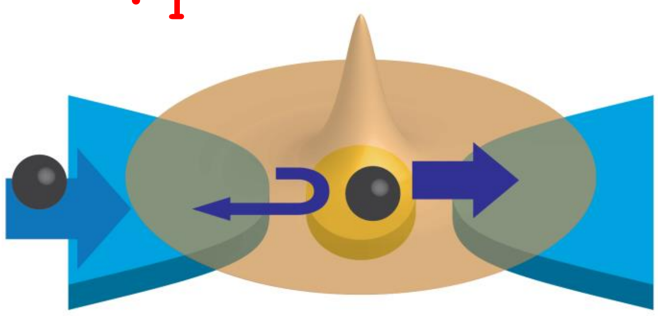
Direct observation of the many-body effect



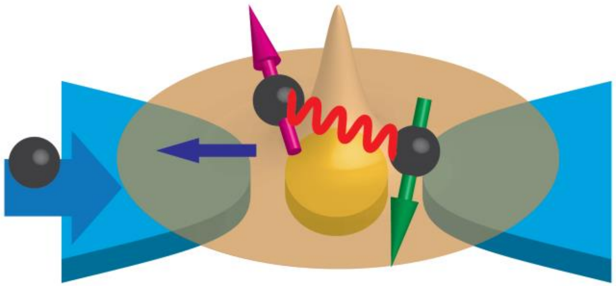
Shot noise contains signature of 2 e scattering which is not in the dI/dV

Non-linear Fano factor : « effective charge »

$$S(V^3) = 2eF_K I(V^3)$$



$$F_K = \frac{P_1 + 4P_2}{P_1 + 2P_2}$$



Not normalized ! $P_1 + P_2 \neq 1$

no interaction $P_2=0$

$F_K=1$

SU(2) $P_1=P_2$

$F_K=5/3$

Sela *et al*, PRL **97** 086601 (2006)

SU(4) $P_1=2P_2$

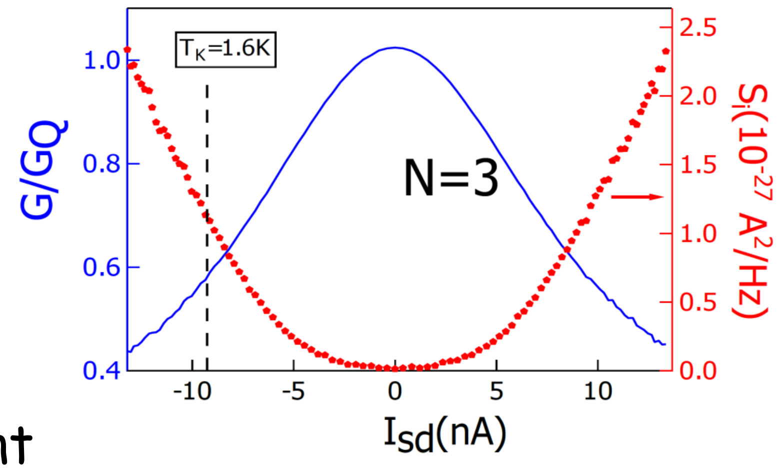
$F_K=3/2$

Mora et al, PRB **80** (2009)

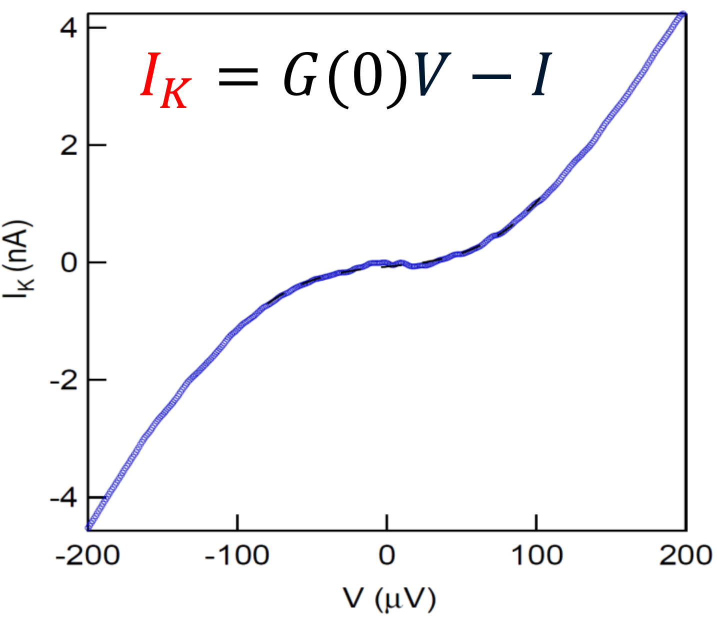
Sakano et al PRB **83** (2011)

F_K measures the probability for 2-particle scattering

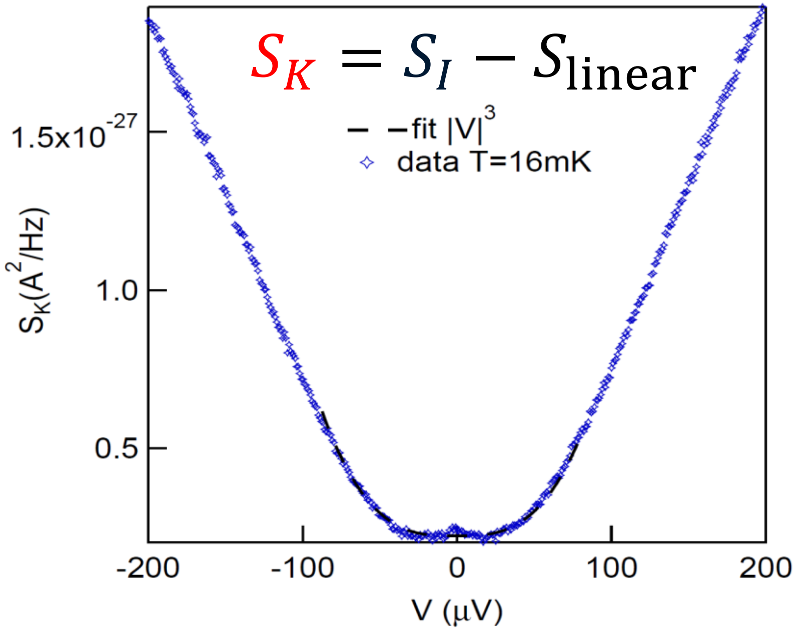
Extraction of non-linear Fano factor



Back scattered current



Nonlinear noise

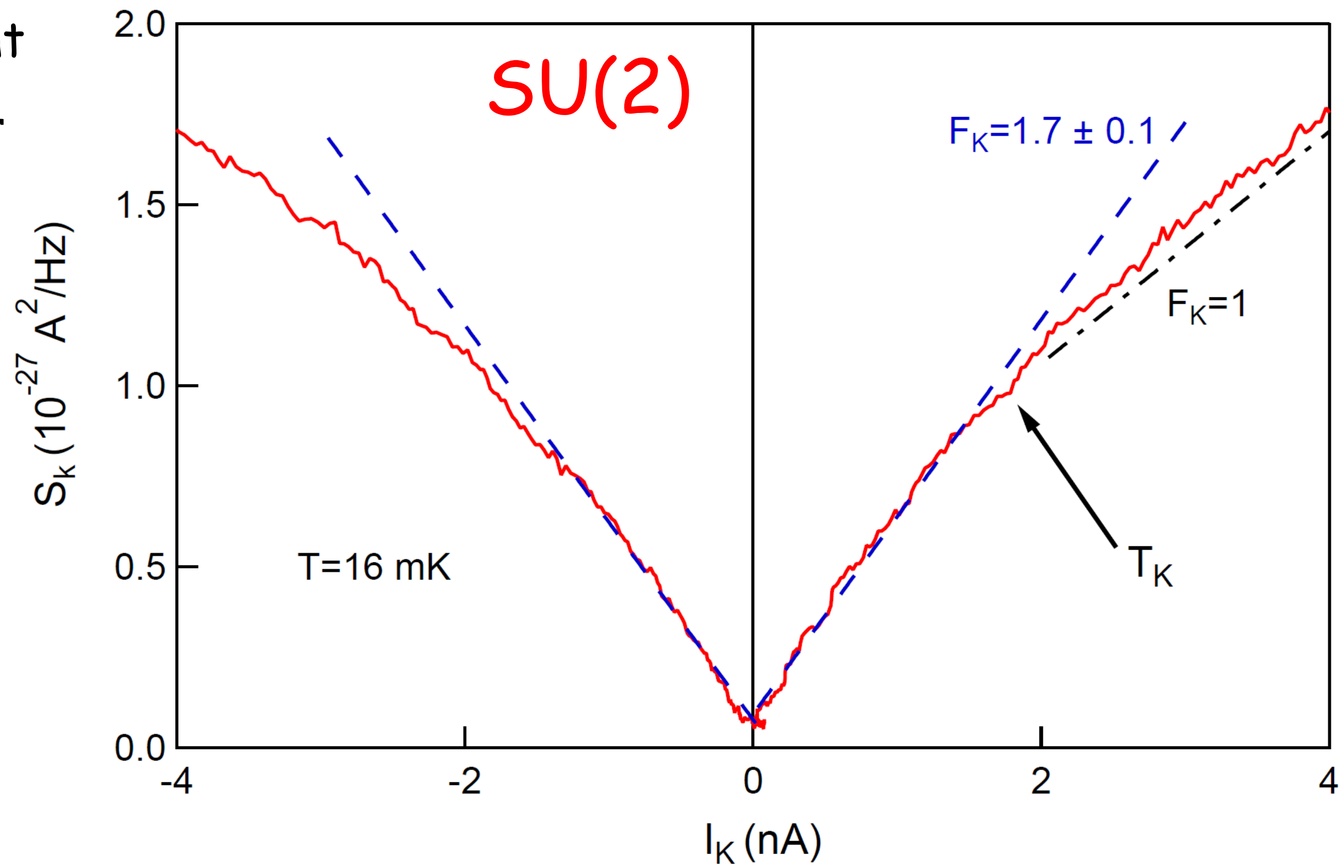


Non equilibrium Fano factor: $S_K = 2eF_K I_K$

Measurement of Kondo Fano factor

Nonlinear component

$$S_K = S_I - S_{\text{linear}}$$



Other experiments

Zarchin et al, PRB (2008)

Yamauchi et al, PRL (2011)

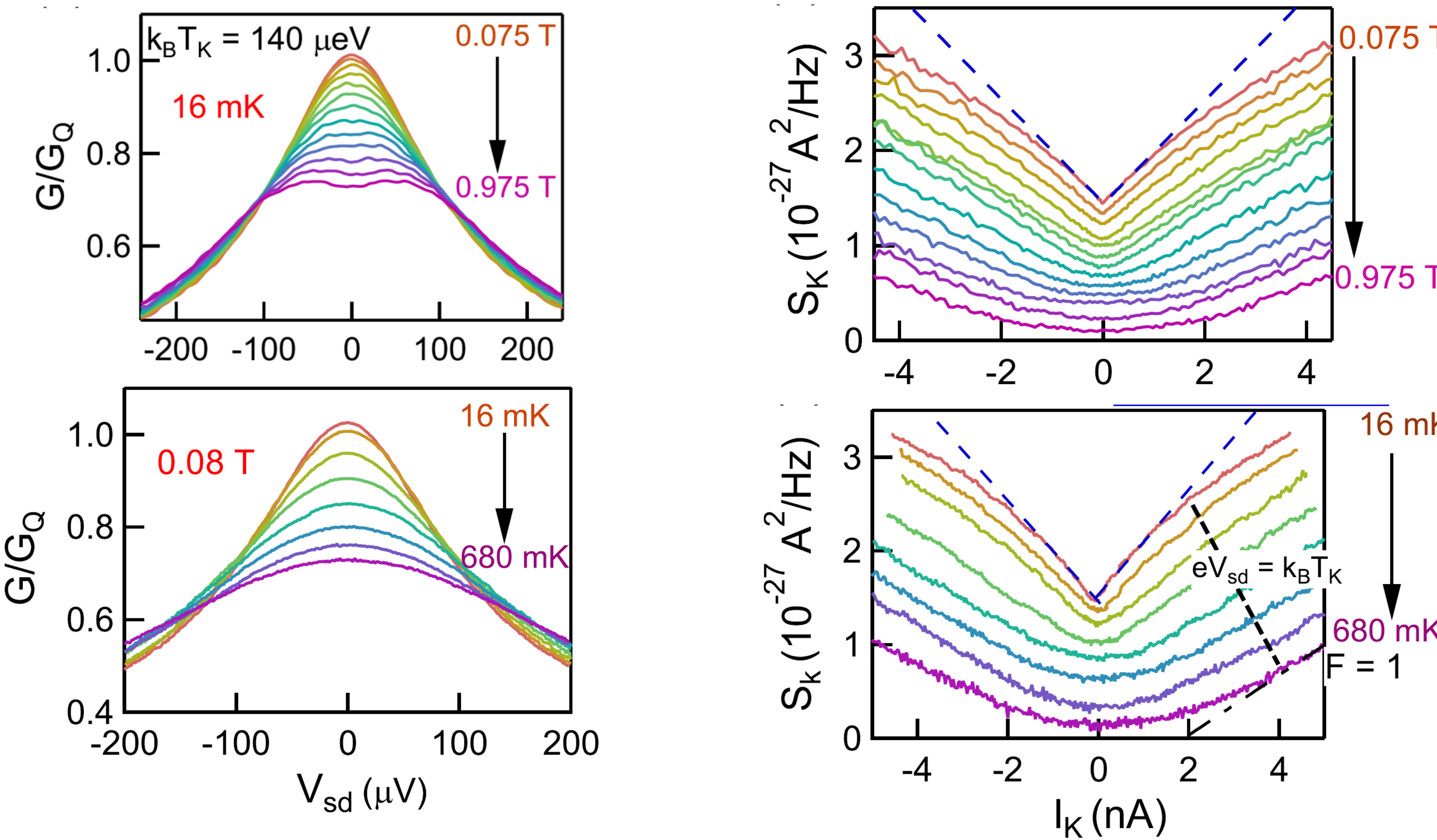
not quantitative

Asymetric dot, low T_K

Back scattered current $I_K = G(0)V - I$

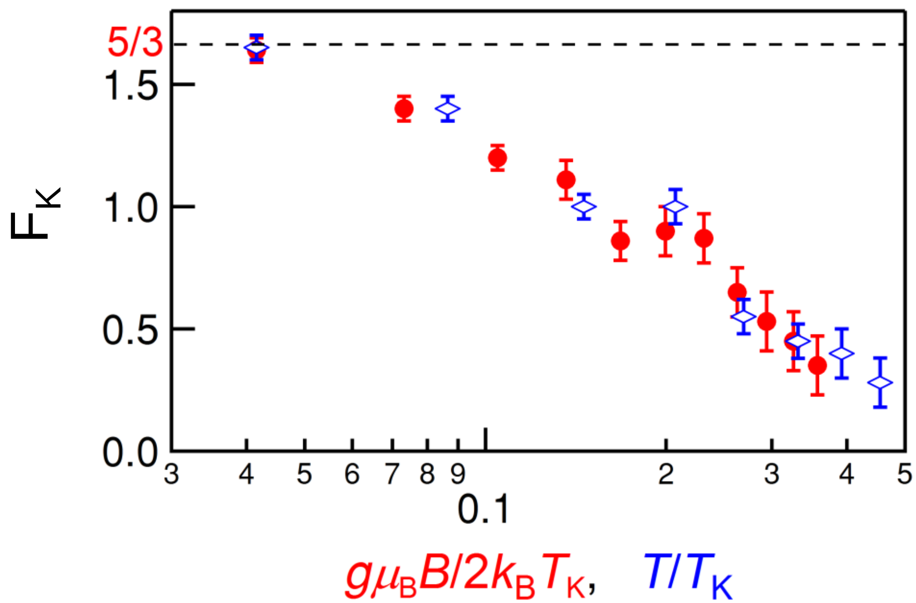
Quantitative agreement with theory

Evolution of Kondo shot noise



2 particles scattering destroyed by magnetic field and temperature

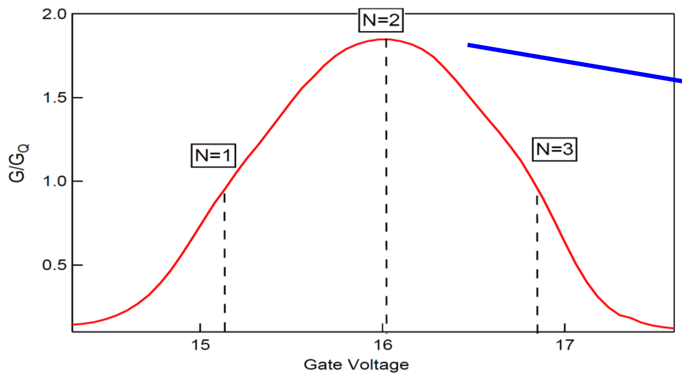
Scaling properties of F_K

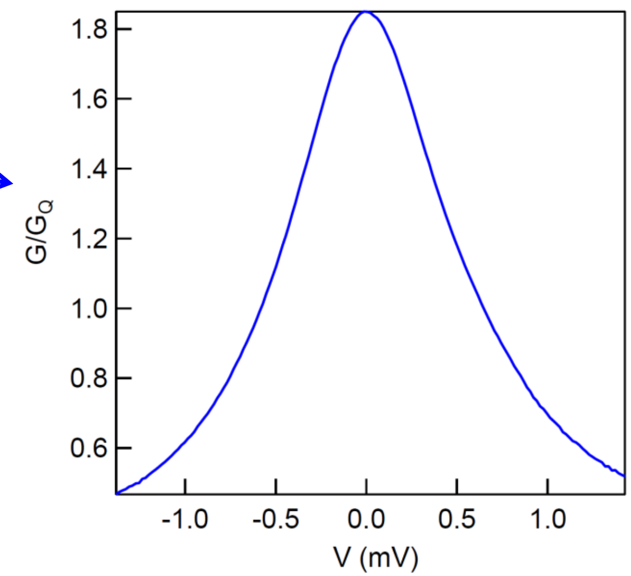


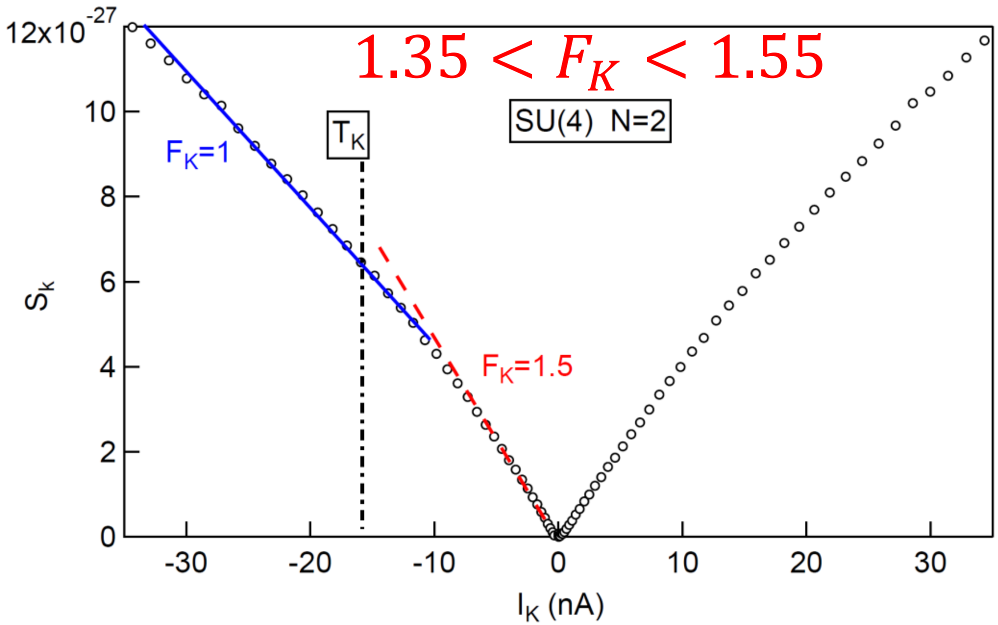
Seems to be logarithmic

Same scaling properties as the conductance

F_K for SU(4) N=2 electrons







3/2 predicted for SU(4):

Mora *et al* PRB **80**, 155322 (2009)

Sakano *et al* PRB **83,** 075440 (2011).

Interaction decreases when degeneracy increases

Significance of these experiments

Around equilibrium

Kondo state = non interacting quasi-particles

Noise = Landauer-Buttiker theory

Out of equilibrium

Interaction between quasi-particles shows up

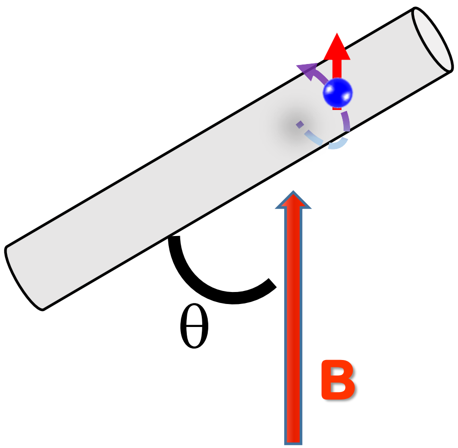
Noise is non-linear and strongly enhanced

→ $F_K > 1$ appears

Very good quantitative agreement with theory

Extension of Fermi-liquid theory out of equilibrium demonstrated experimentally

Effect of magnetic field



Spin splitting $E^{spin} = \pm g\mu_B BS$

Valley splitting $E^{valley} = \pm g_{orb}\mu_B B_{\parallel}$

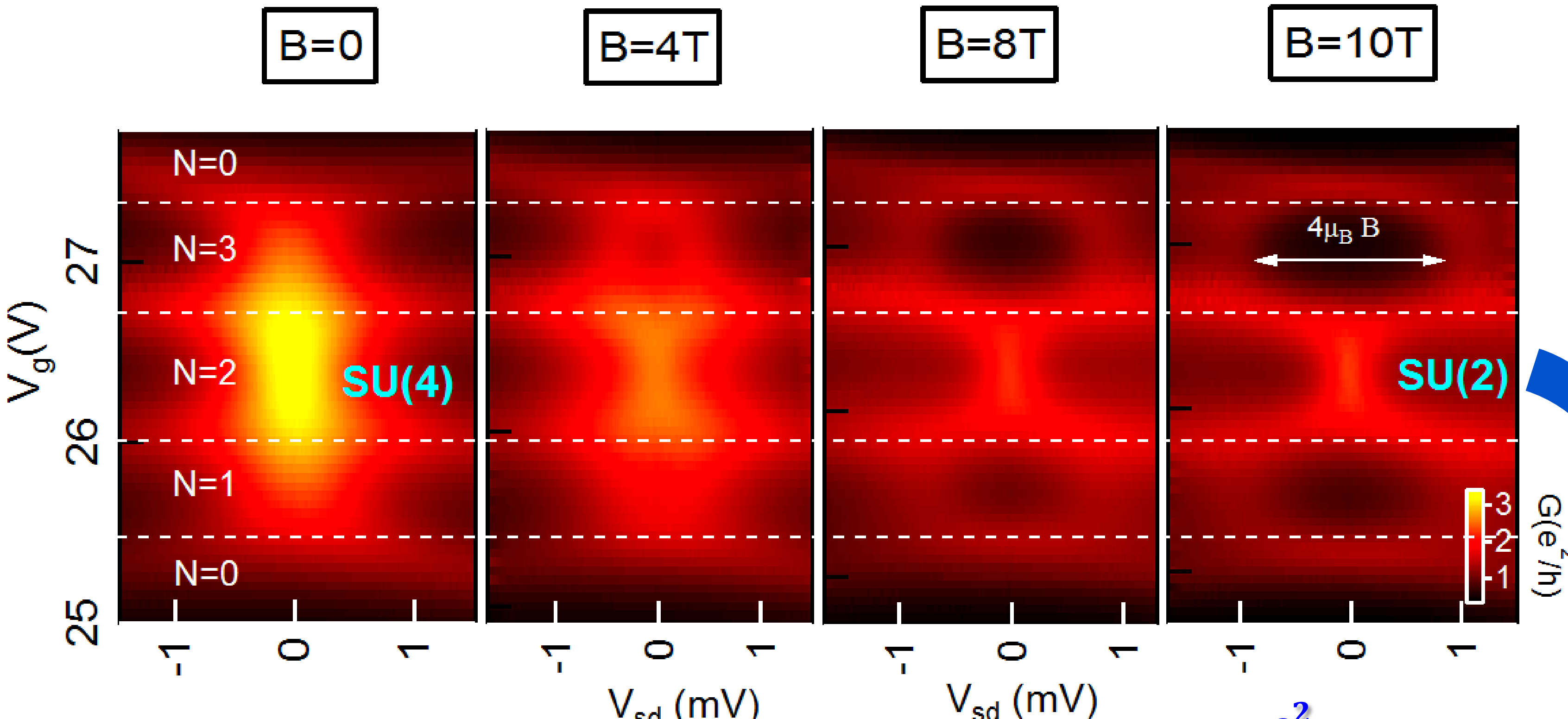
luckily… g_{orb}=4 θ=75

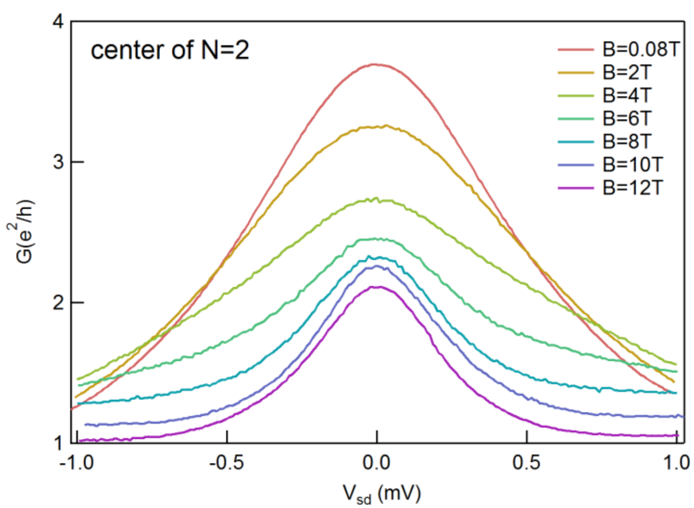
(From NRG + B dependence of excited states)

In this experiment $g_{orb}\cos\theta \sim gS=1$ → $E^{spin} \sim E^{valley}$

$$E^{tot} = (\sigma + \tau)\mu_B B$$

SU(4) to SU(2) crossover at N=2





$4\frac{e^2}{h}$

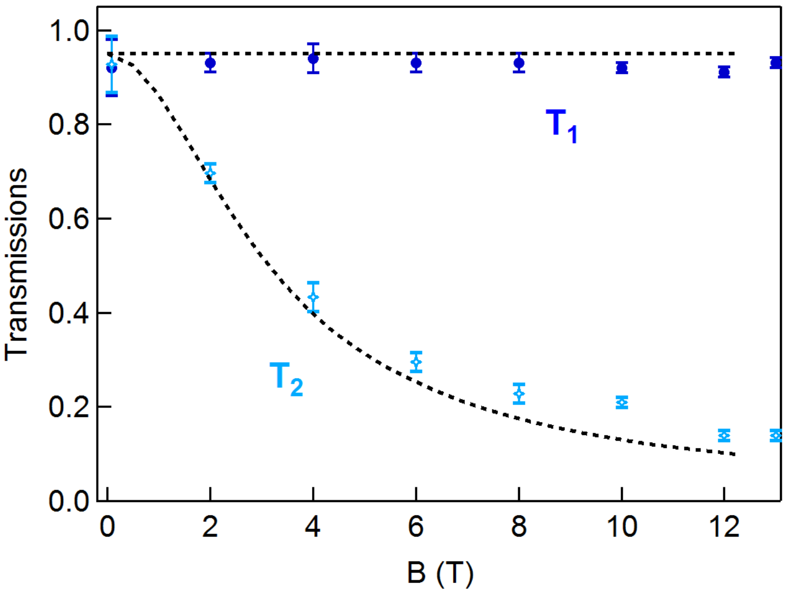
$2\frac{e^2}{h}$

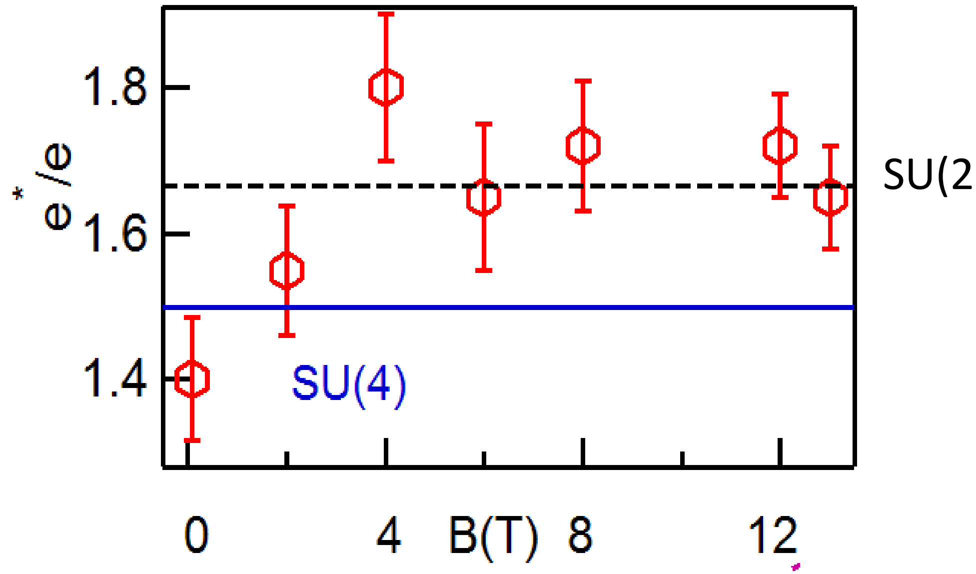
At N=2

G remains maximum at V=0

Continuous crossover for fixed filling N=2

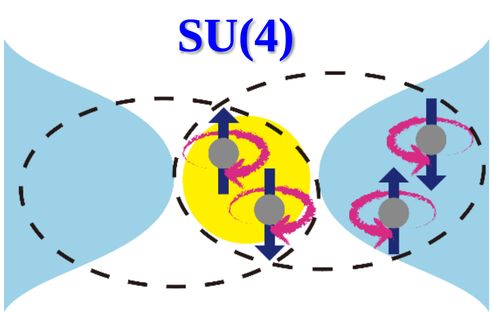
tuned by the magnetic field





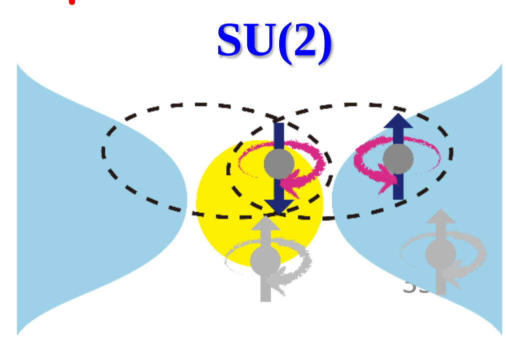
$T_1 = T_2 \approx 1$

2 perfect channels

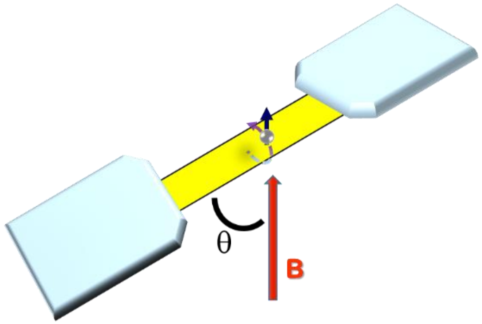


$T_1 \approx 1$ $T_2 \approx 0$

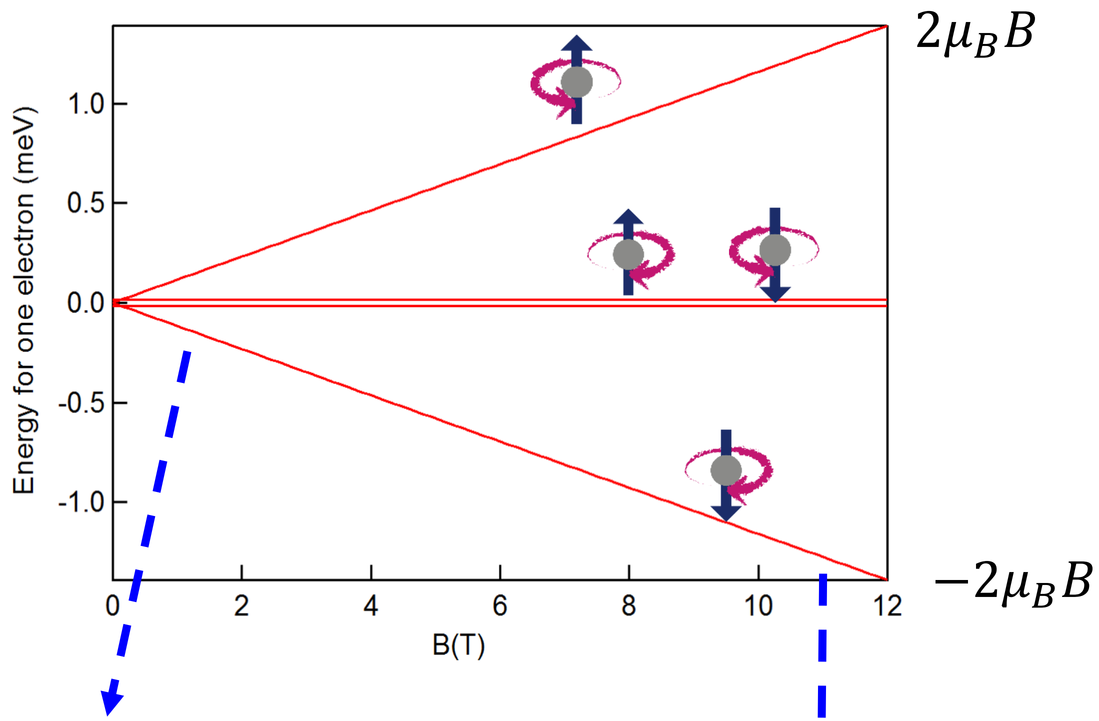
1 perfect channel



Single particle spectrum

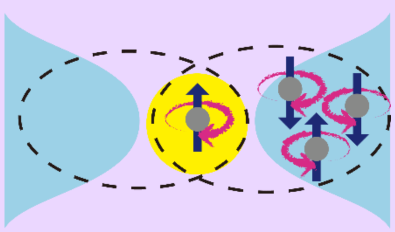


$$E^{tot} = (\sigma + \tau)\mu_B B$$



Kondo ground state

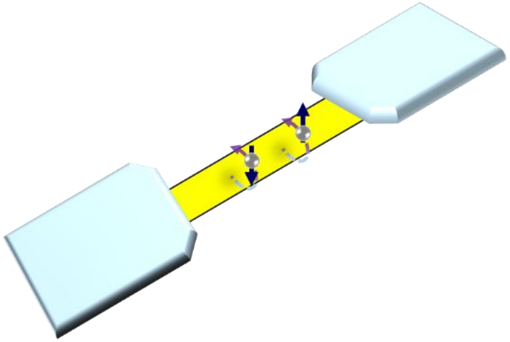
SU(4)

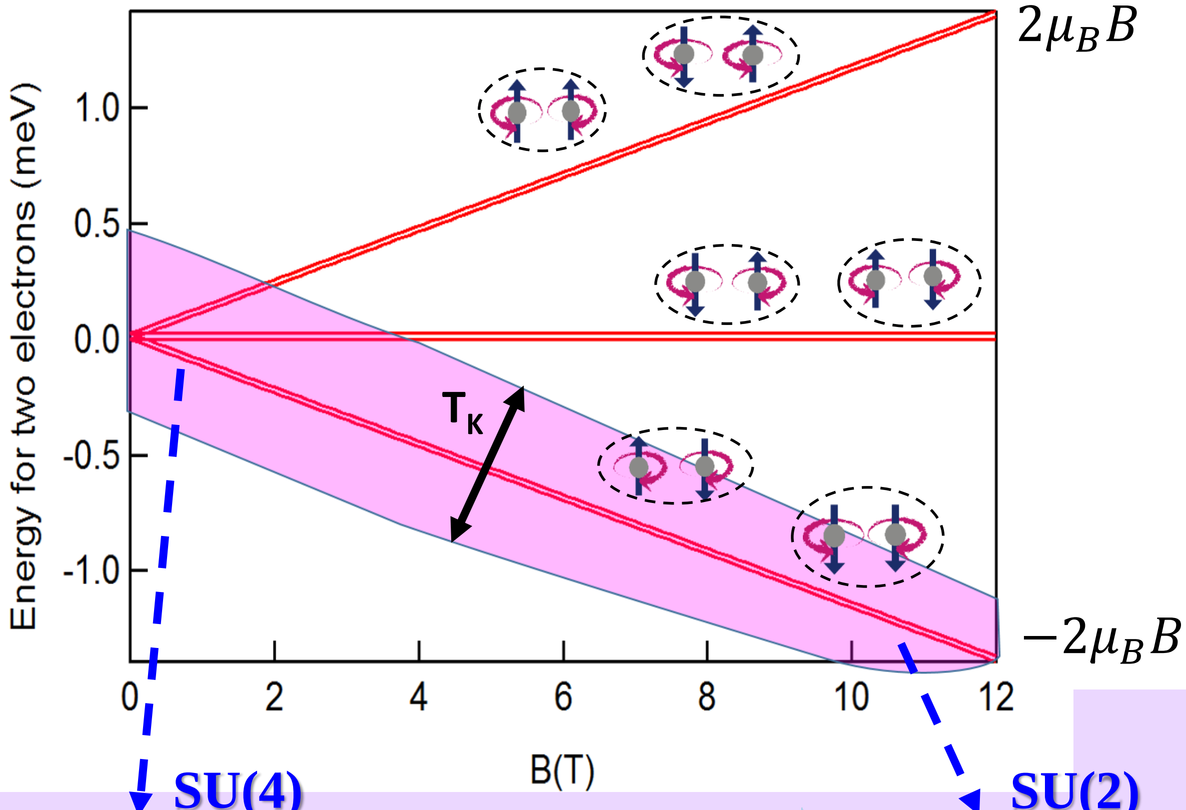


No Kondo effect

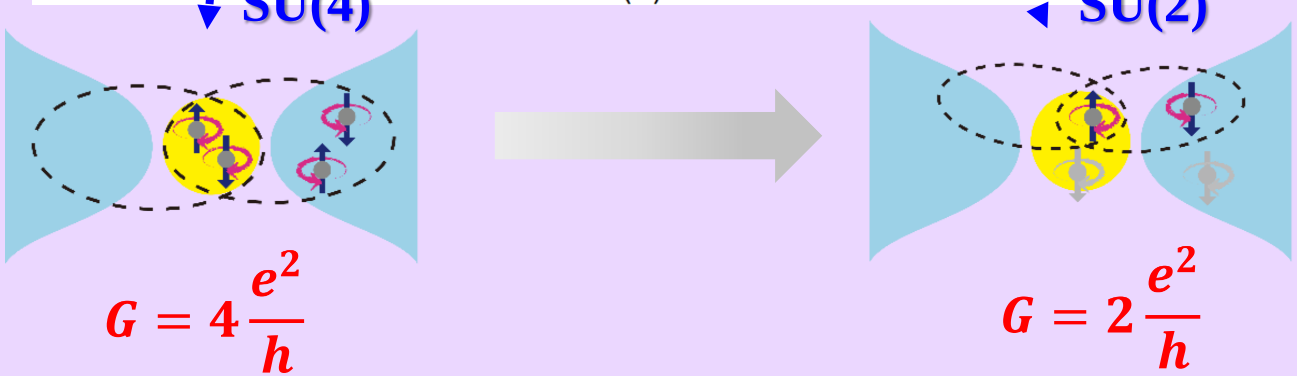
Degeneracy lifted : Kondo state destroyed

Two-particle spectrum: SU(4)-SU(2) crossover





Kondo ground states



Crossover between two different many-body states

CONCLUSION

On-chip collision experiment: Probe dynamical behaviors of a quantum many body system

Noise shows the symmetry class

Direct evidence of 2 quasi-particles scattering due to interaction

F_K~1.7 for two different kind of SU(2)

F_K~1.5 for SU(4) @ N=2

Cross-over in the symmetry class monitored by shot noise

e* is a reliable measure of residual interaction/quantum fluctuations in a many-body state

Next…

Effect of Superconducting leads:

Mixing Kondo and Andreev states

SU(2) SU(4) symmetry also detected

No « real » supercurrent = disspative environment

Noise in the inelastic cotunneling regime :

Conductance peak enhanced by interaction

What about the noise?