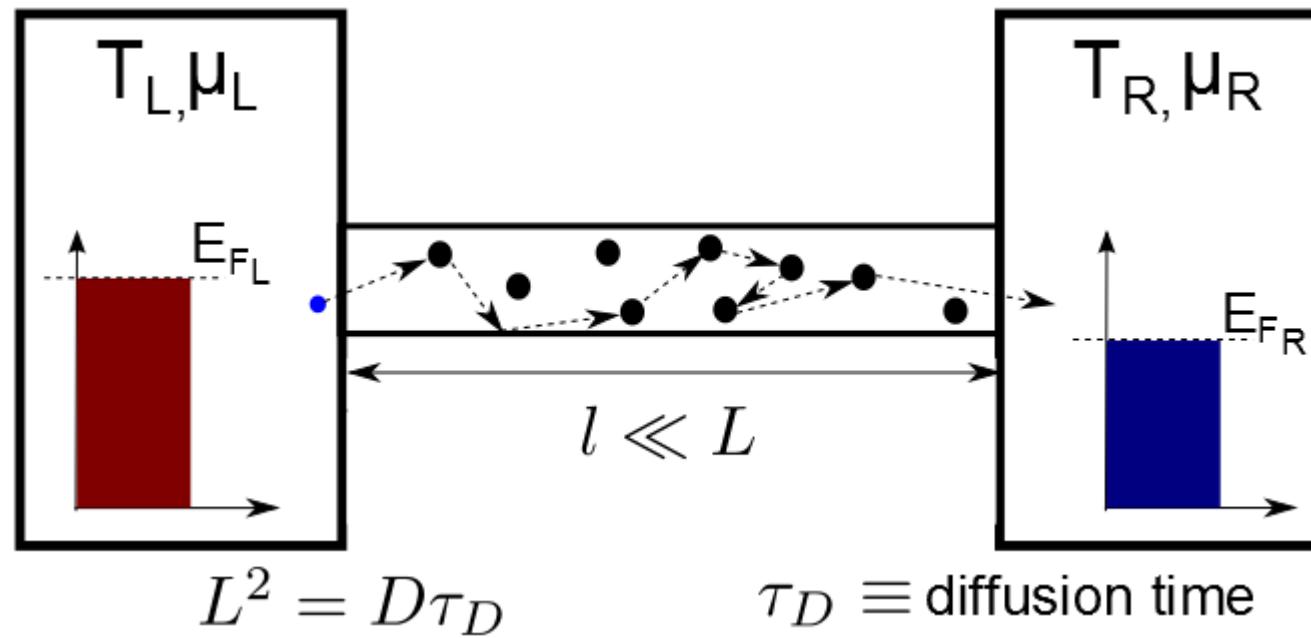


Electron Energy Relaxation Dynamics

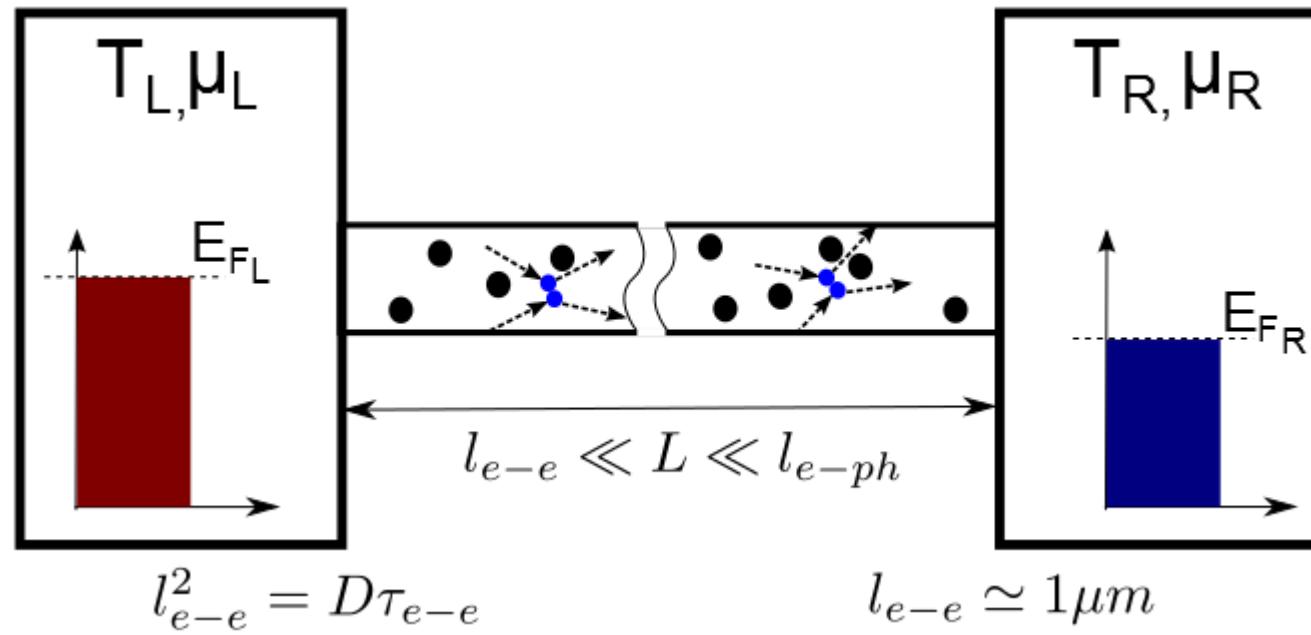
Edouard Pinsolle
Université de Sherbrooke
Reulet's group

GDR

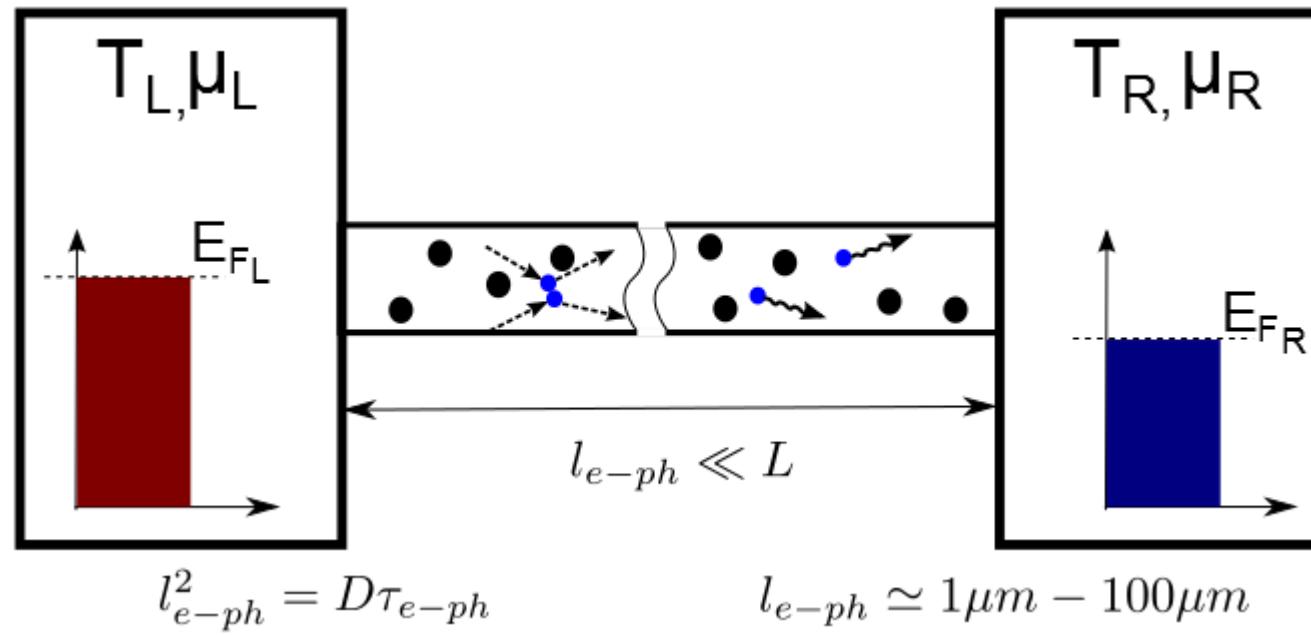
Different Regimes: Diffusive/Elastic



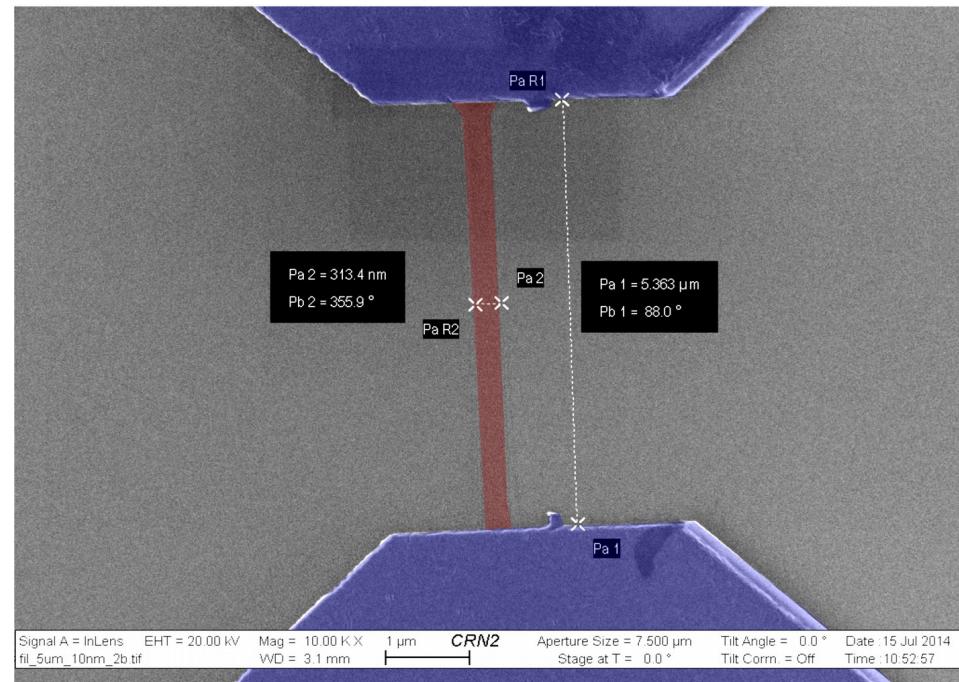
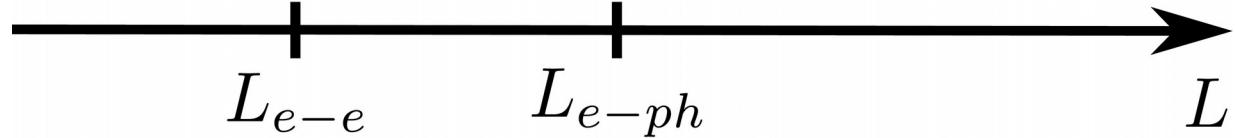
Different Regimes: Hot Electrons



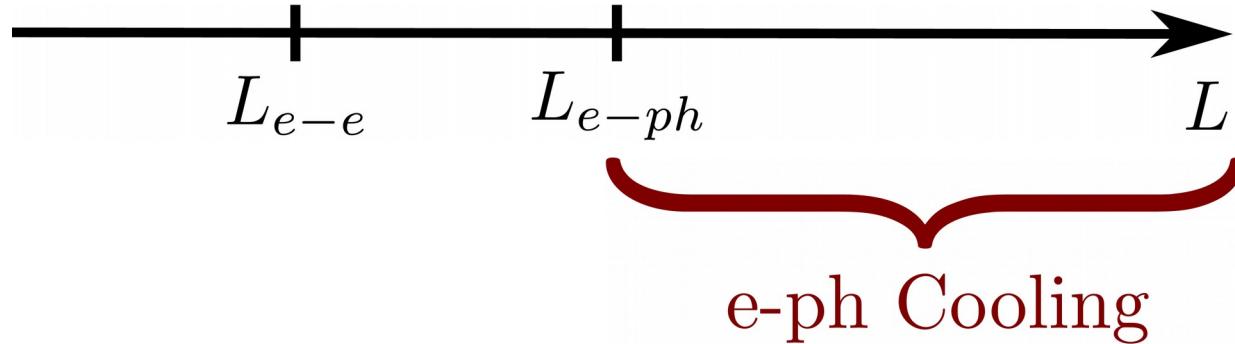
Different Regimes: Macroscopic



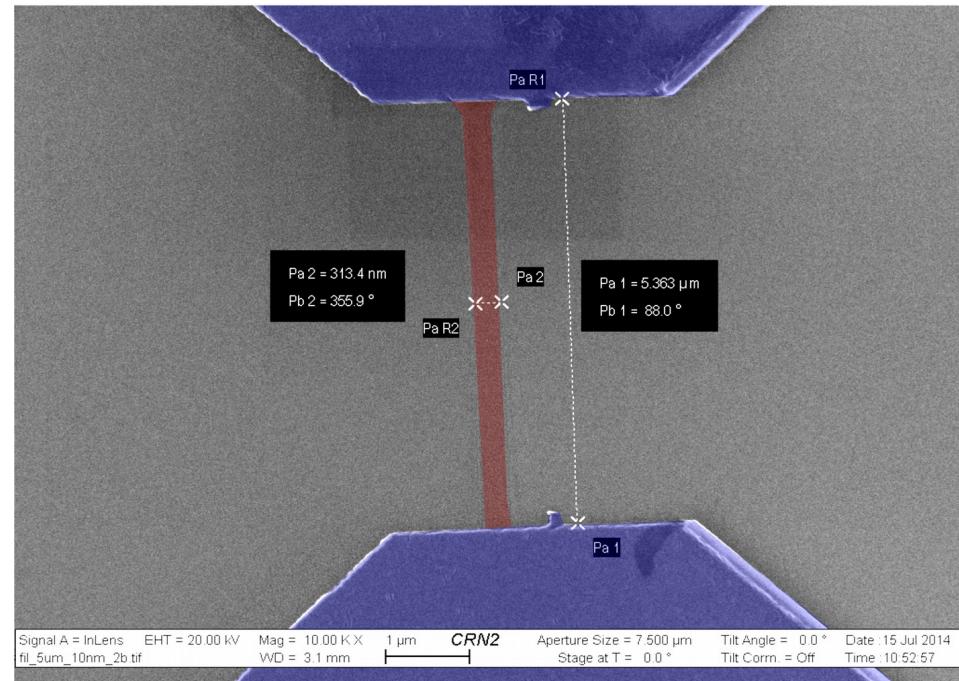
Cooling Processes in Metals



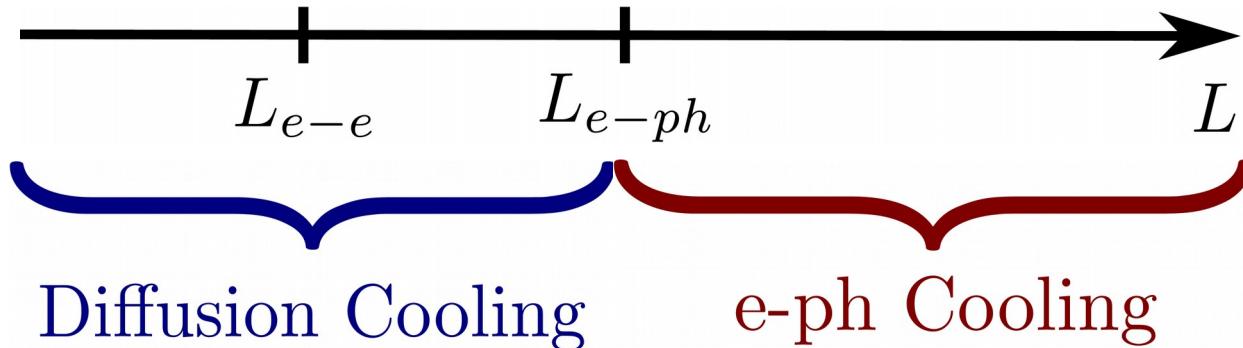
Cooling Processes in Metals



$$\frac{1}{\tau_{e-ph}} = A \times T^3$$

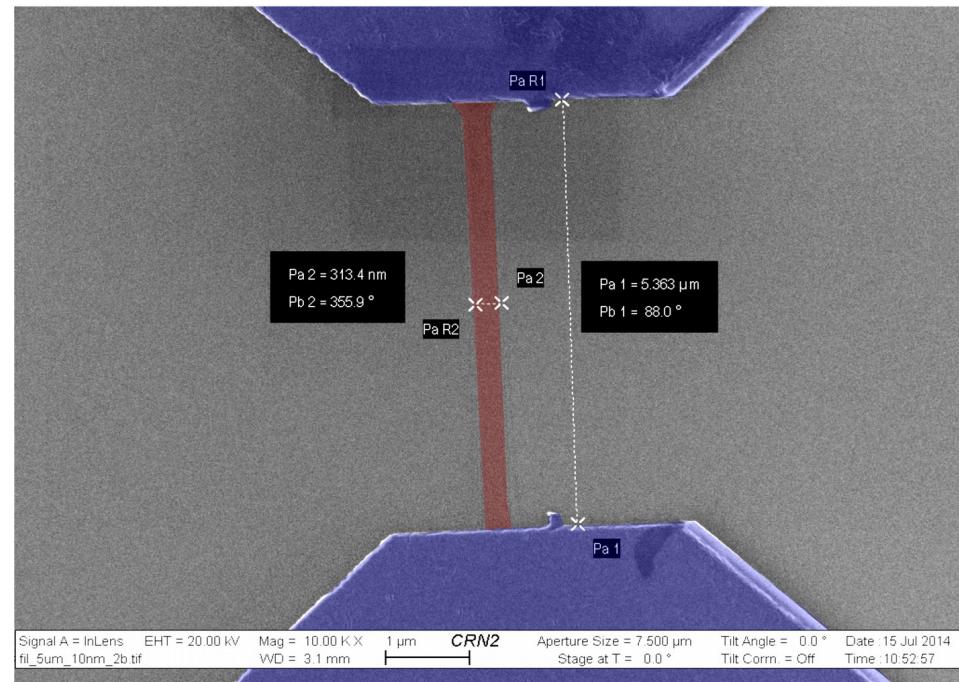


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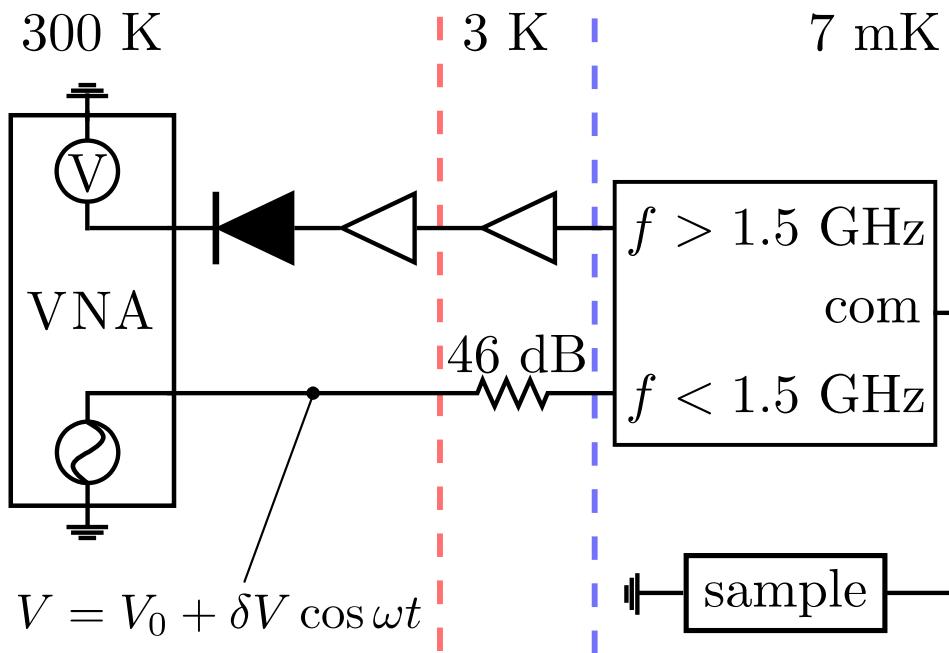


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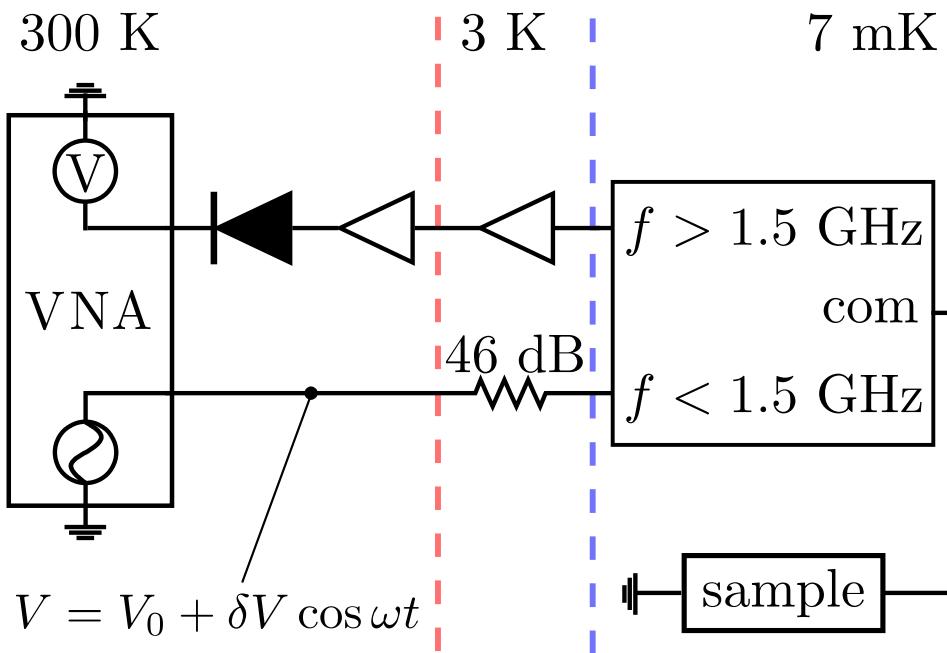
$$\frac{1}{\tau_D} = \frac{D}{L^2}$$



The Experimental Concept



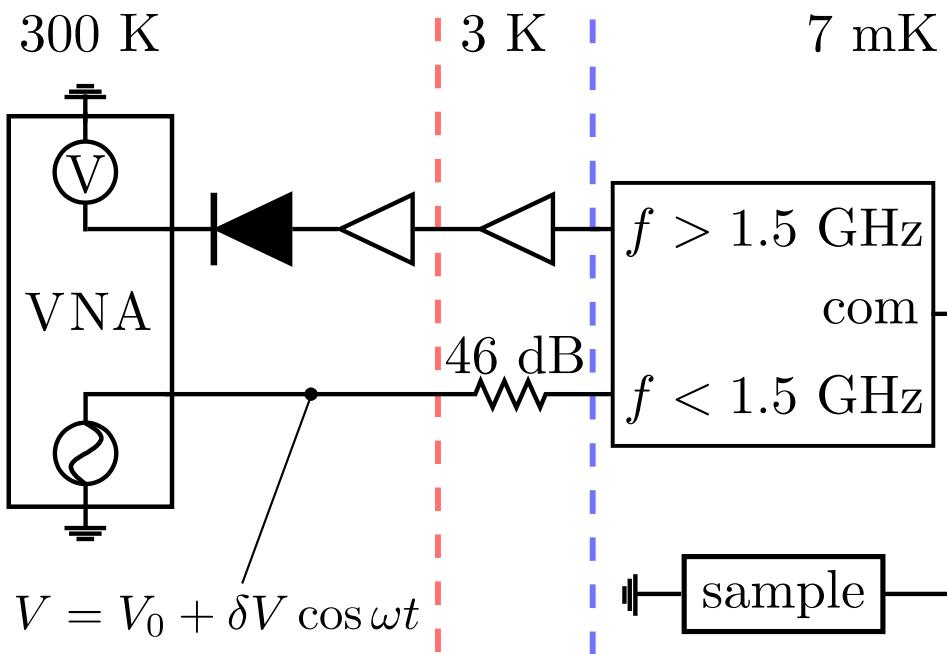
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Voltage fluctuations and Temperature:

$$S = 2Rk_B T_e$$

The Experimental Concept



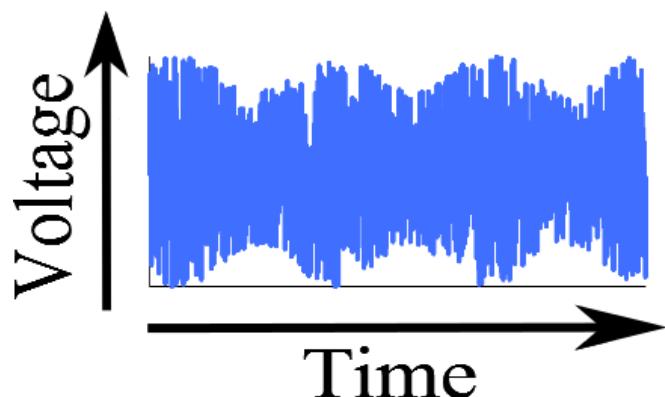
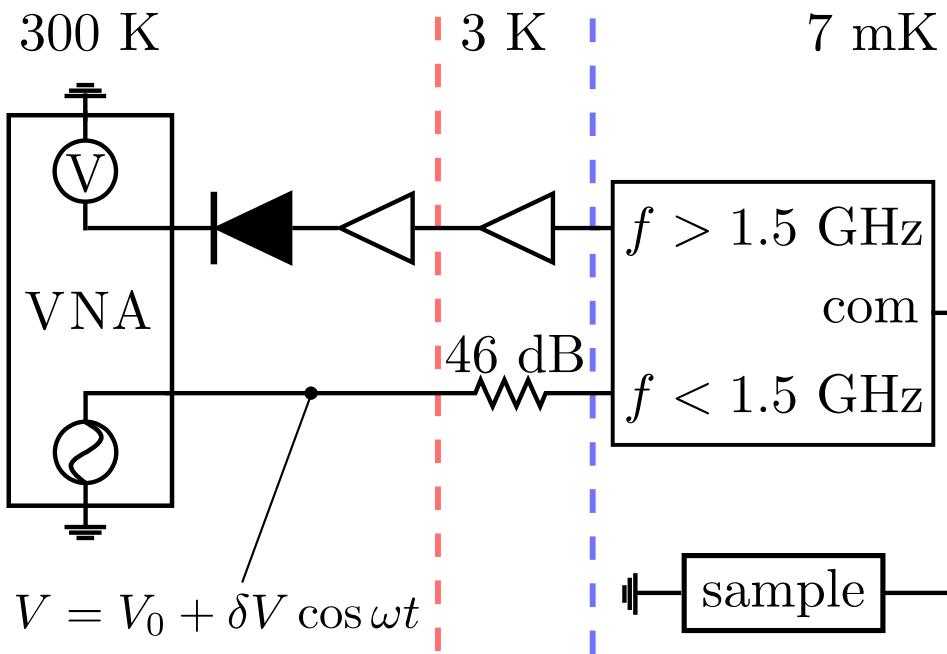
Voltage fluctuations and Temperature:

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$$P = V^2/R$$

The Experimental Concept



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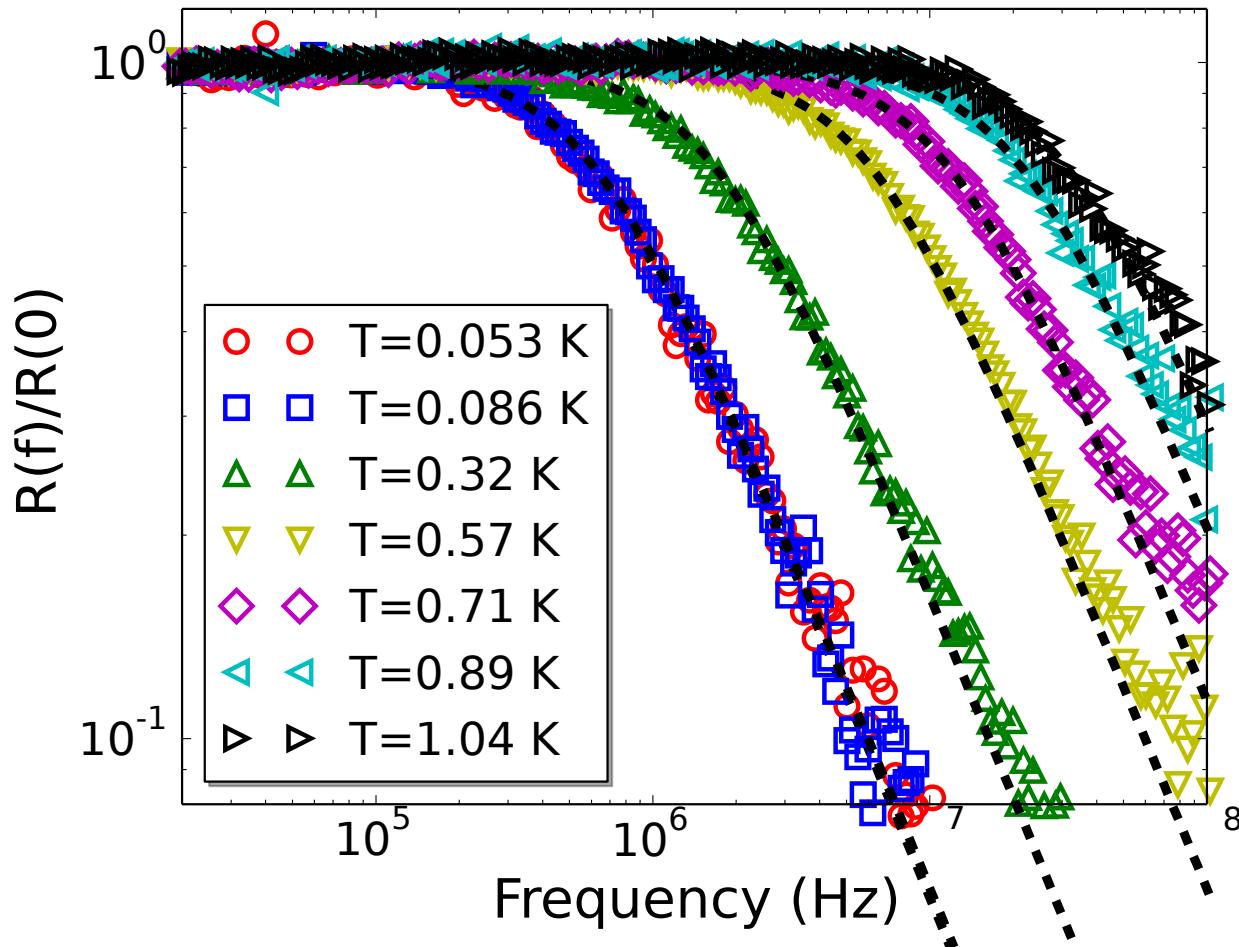
Periodic Joule Heating:

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Thermal Impedance:

$$R(\omega) = \delta T_e(\omega) / \delta P(\omega)$$

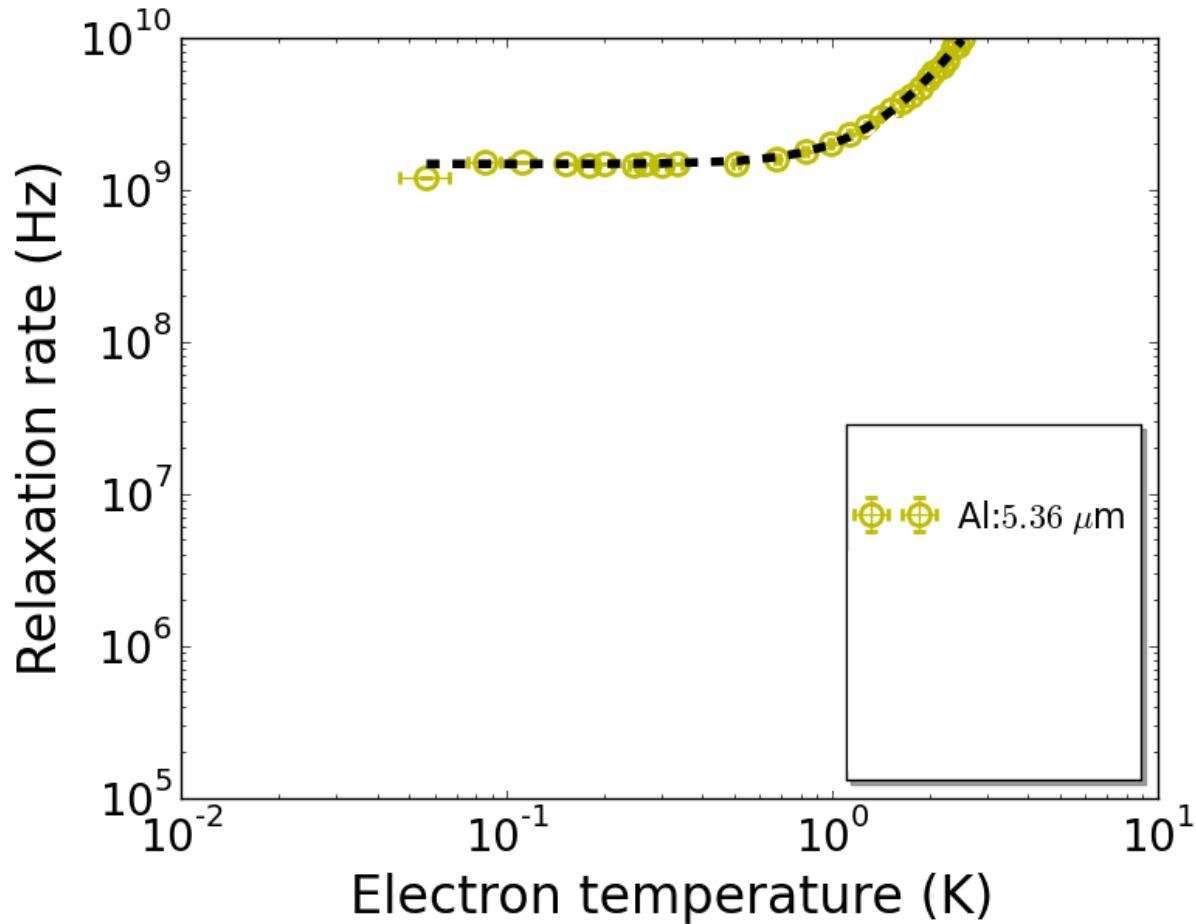
NTI for a 50 micrometers Long Wire



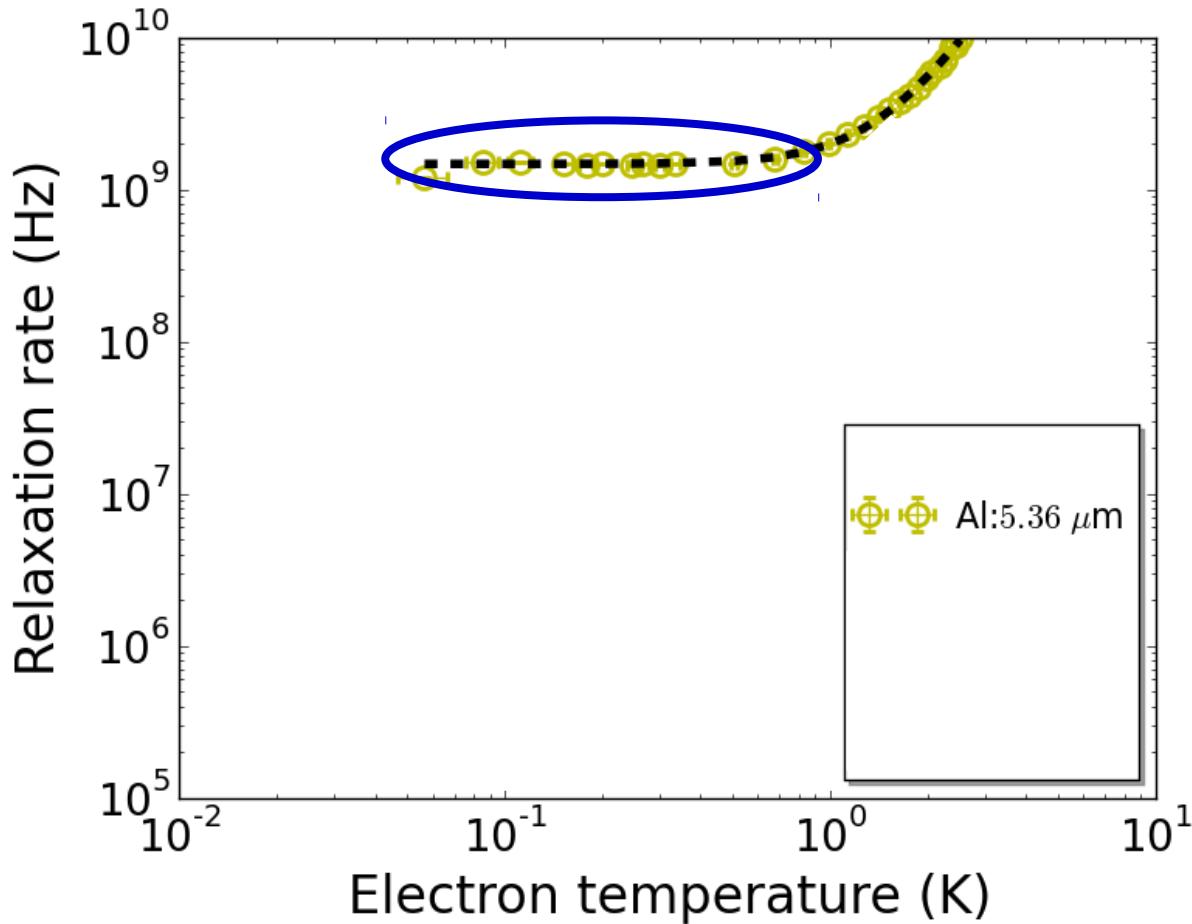
Lorentzian fit:

$$\left| \frac{R(f)}{R(0)} \right|^2 = \frac{1}{1 + (2\pi f \tau)^2}$$

Measured Relaxation Rates



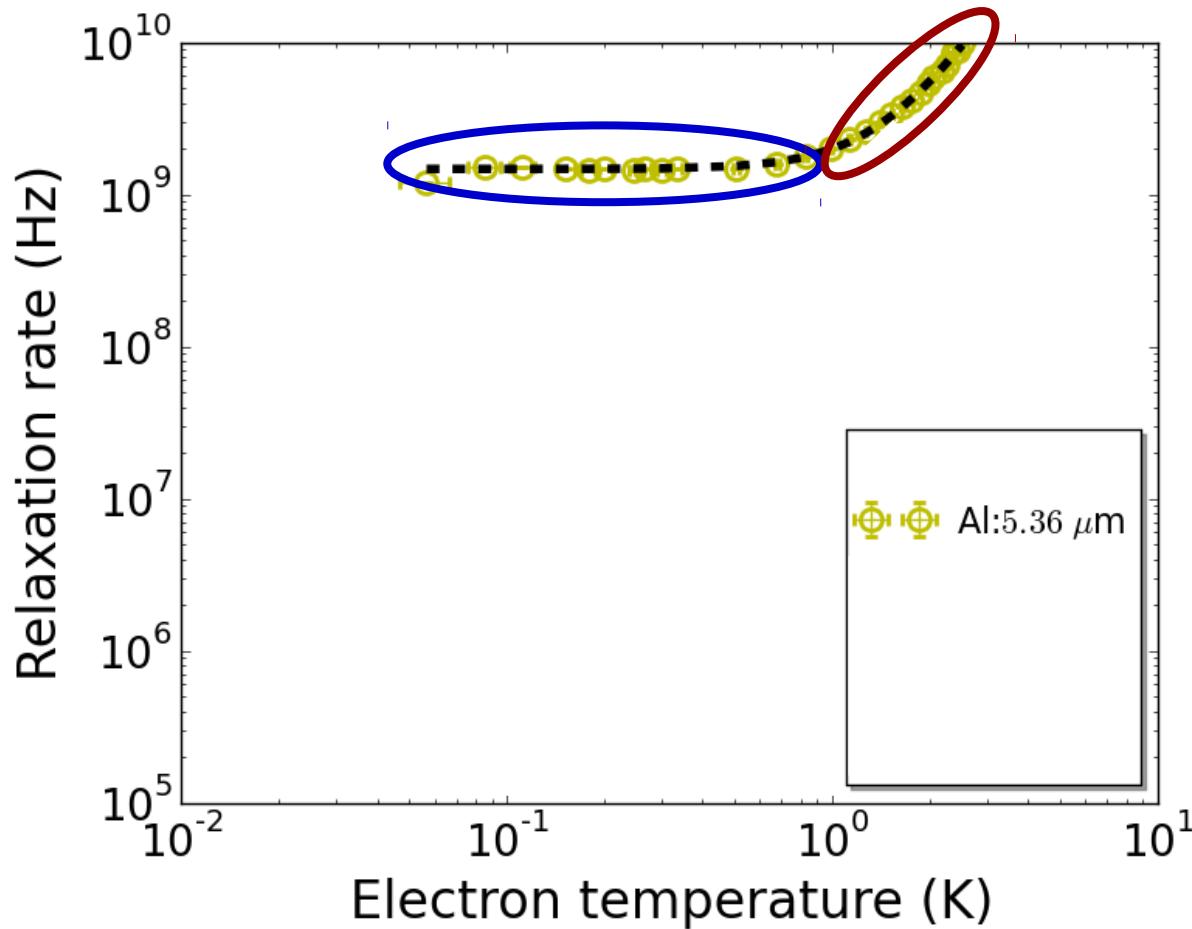
Measured Relaxation Rates



Plateau at low temperature:

$$f_D = \text{cste}$$

Measured Relaxation Rates



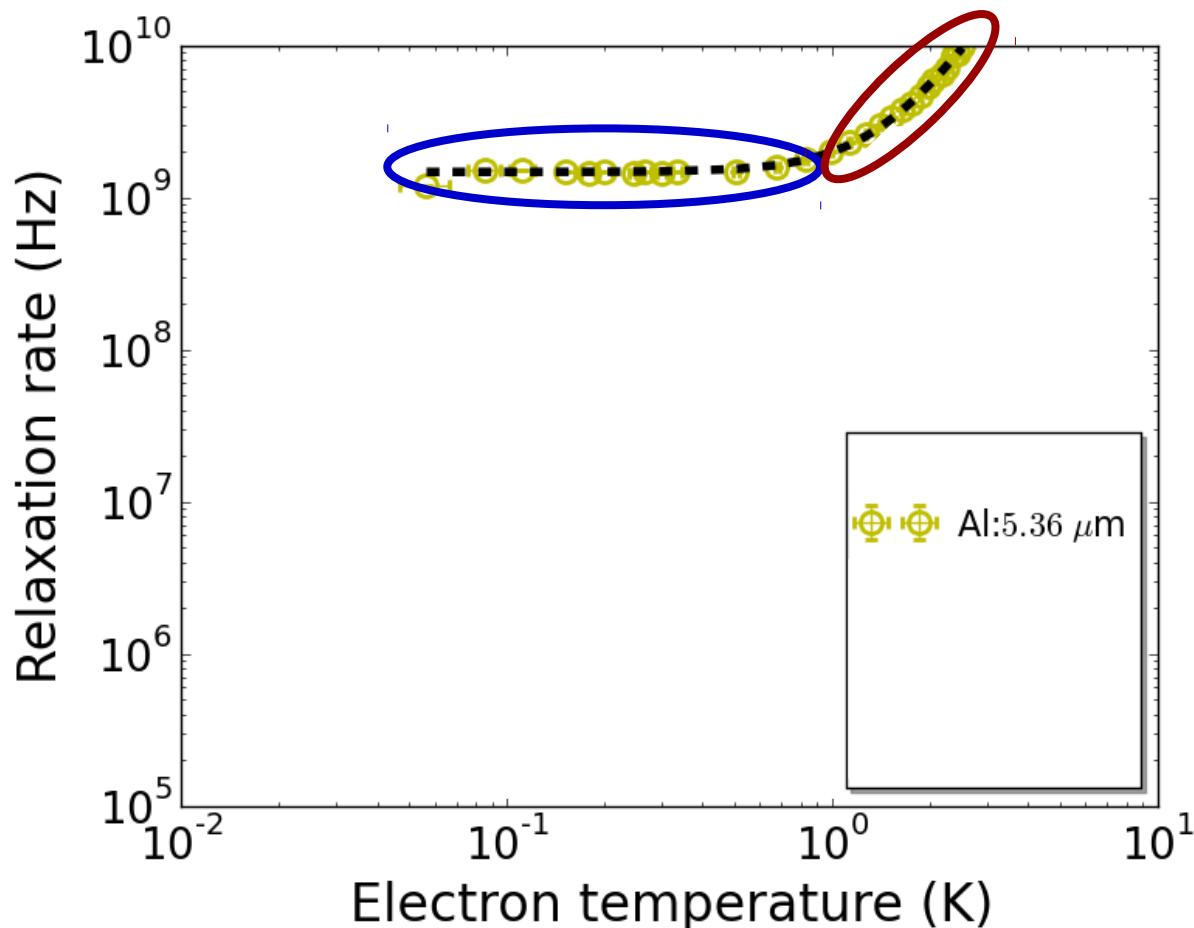
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$$f_{e-ph} = AT^3$$

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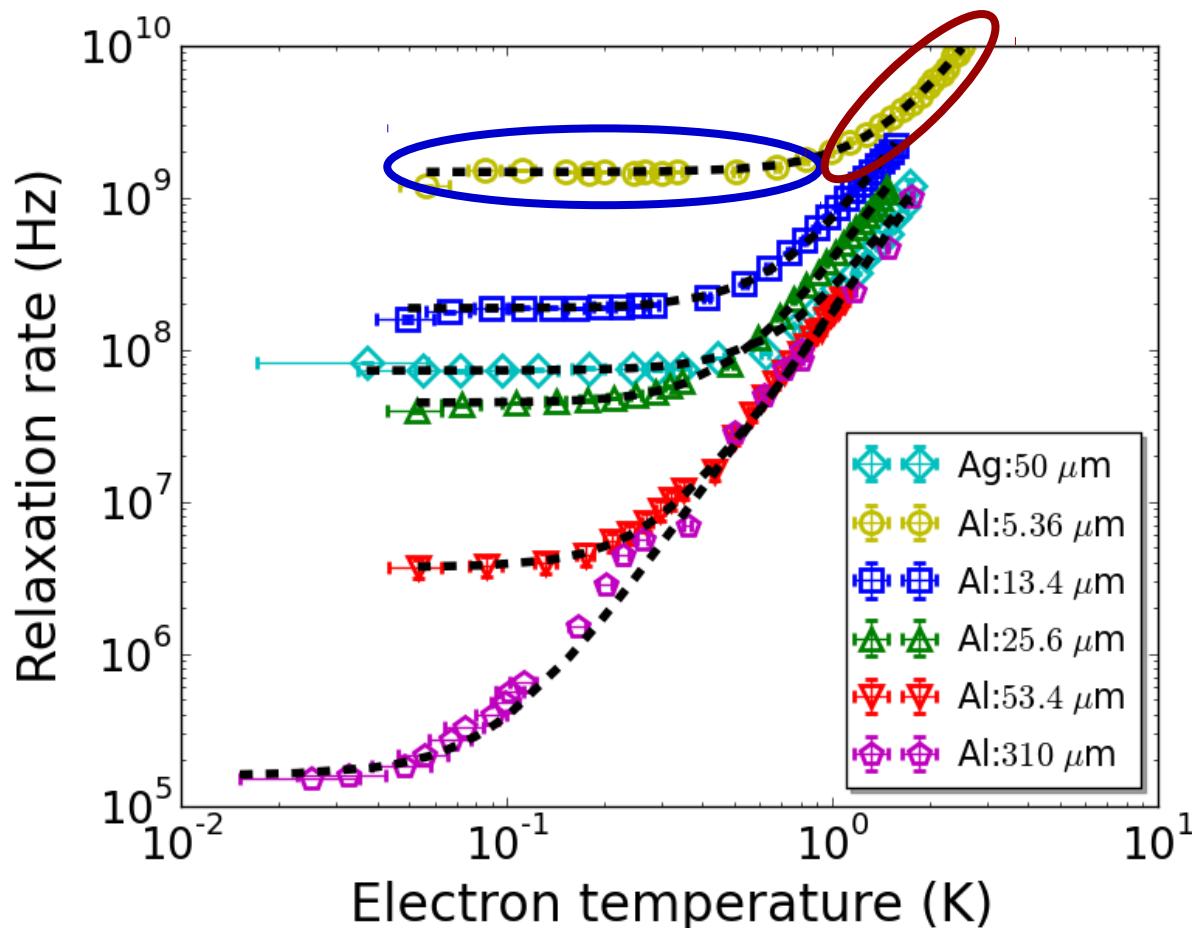
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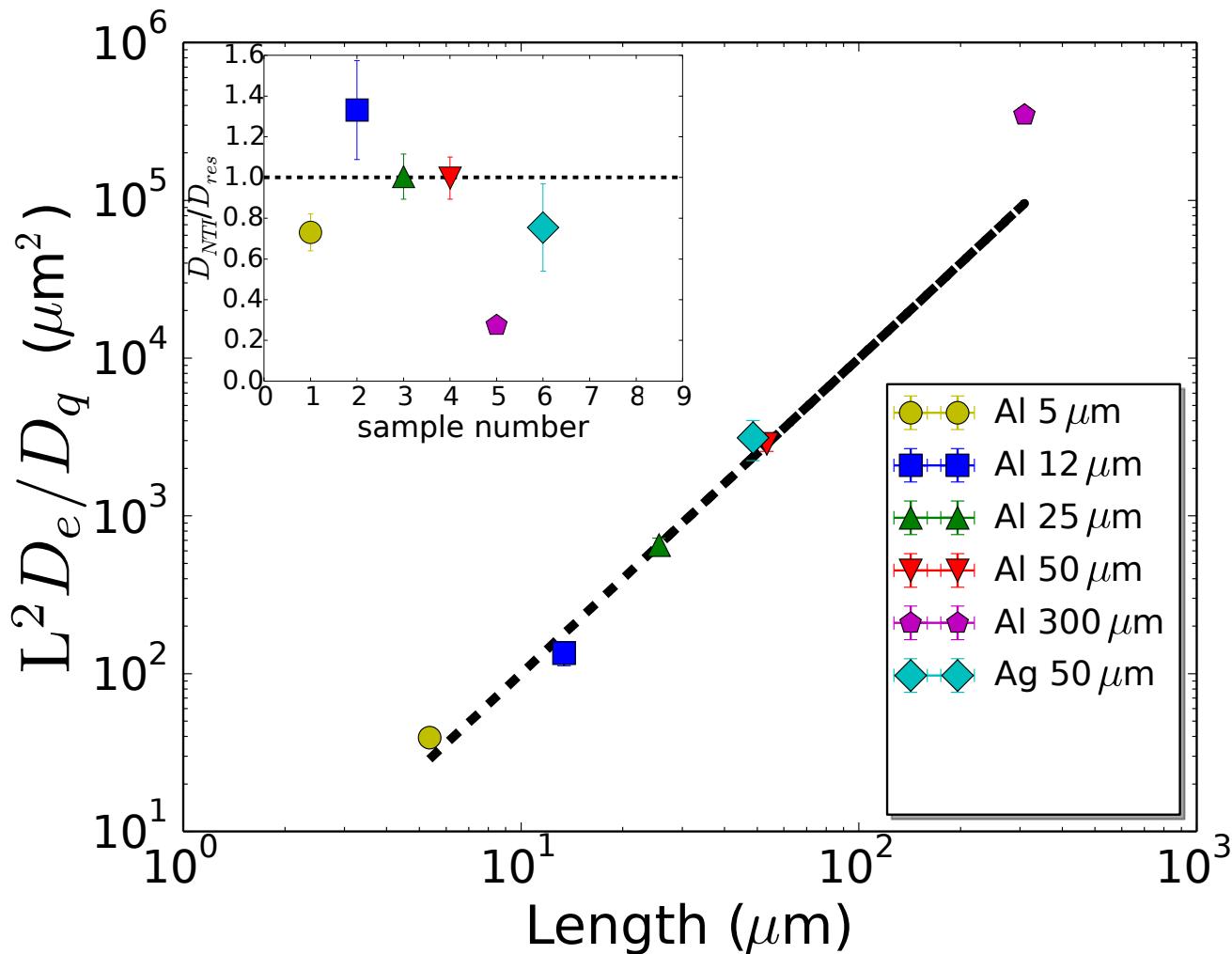
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Wiedemann-Franz and Diffusion law



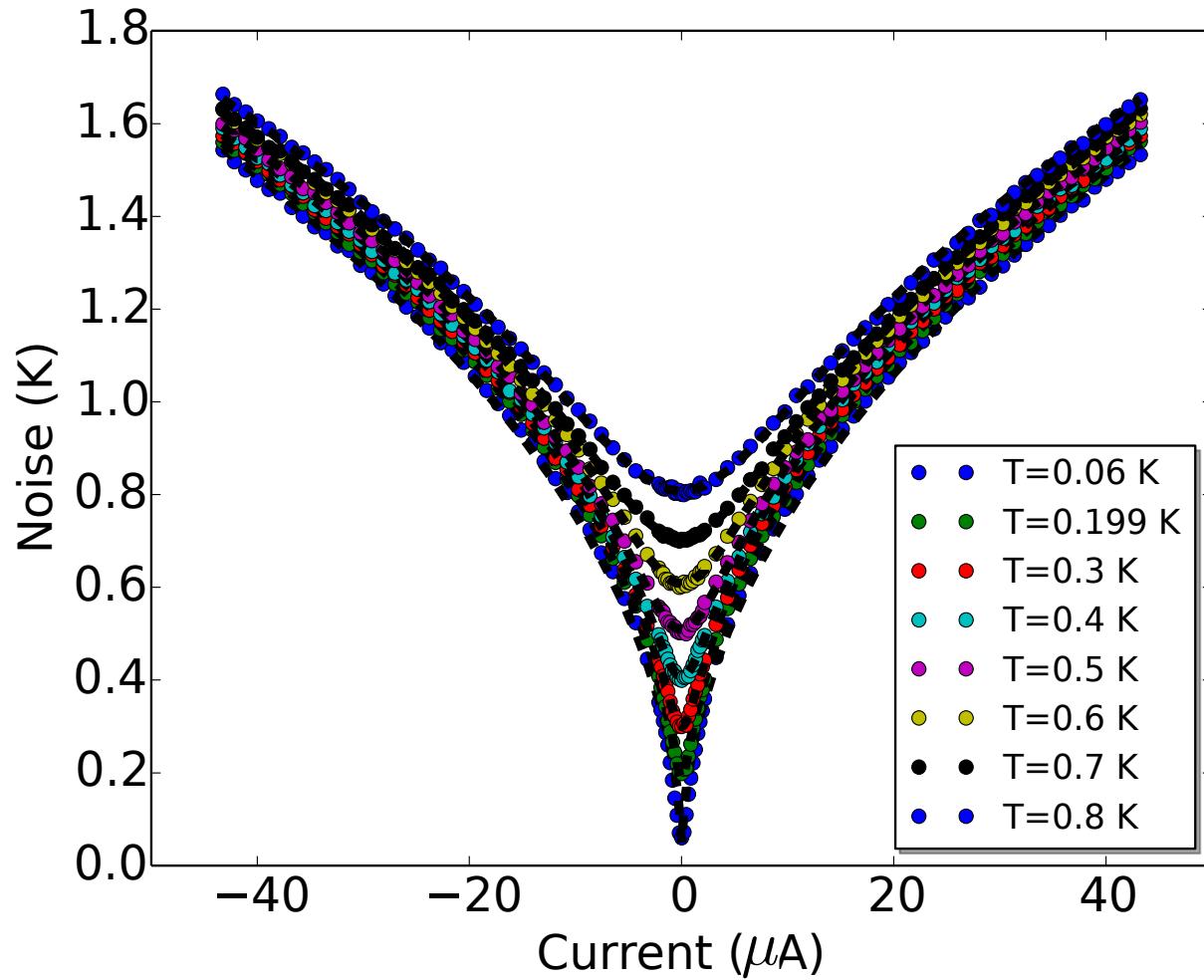
Diffusion Law:

$$L^2 = \tau_D D_q$$

Einstein Relation:

$$\sigma = n e^2 D_e$$

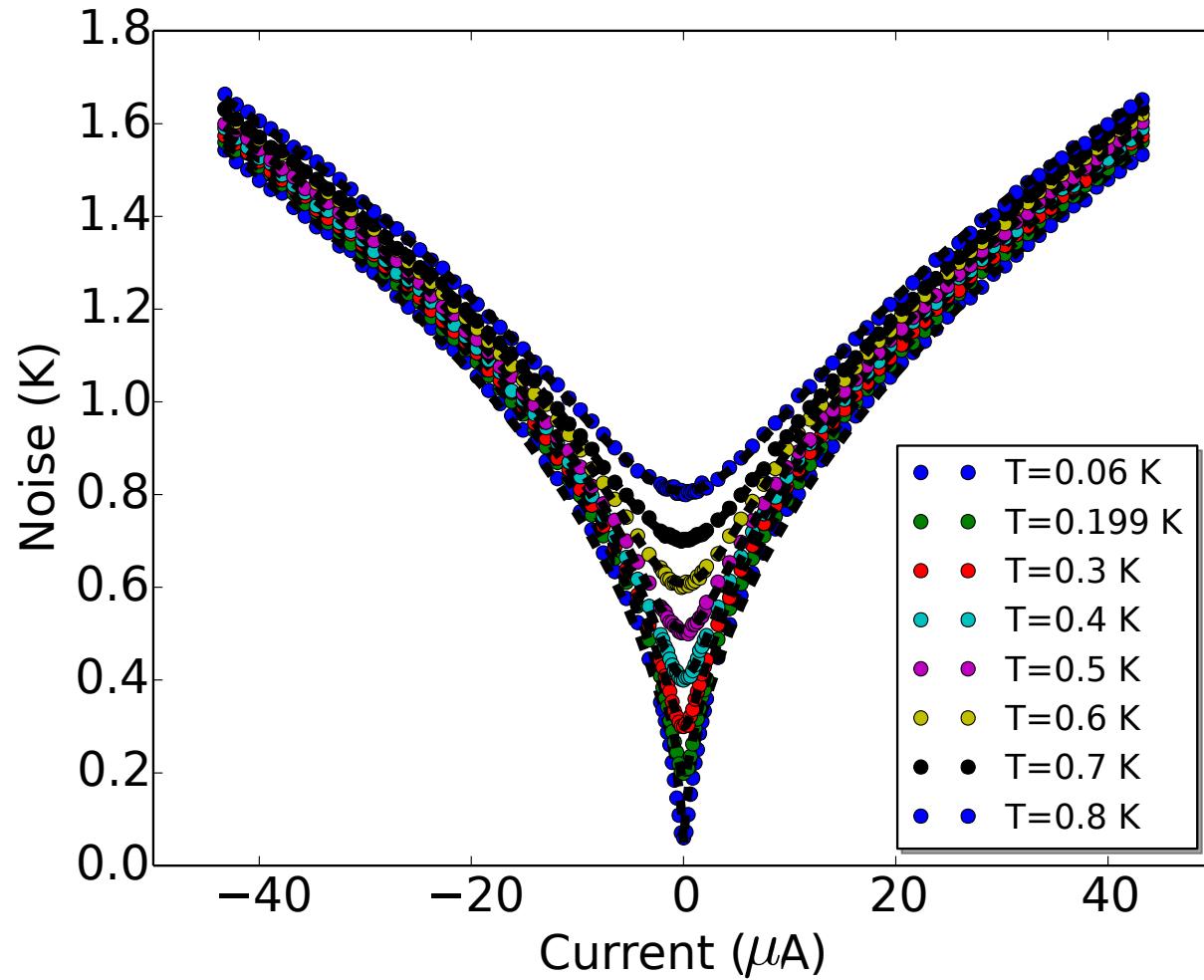
Noise at Zero Frequency



Thermal Noise in presence
of joule heating:

$$S = 2Rk_B T_e$$

Noise at Zero Frequency



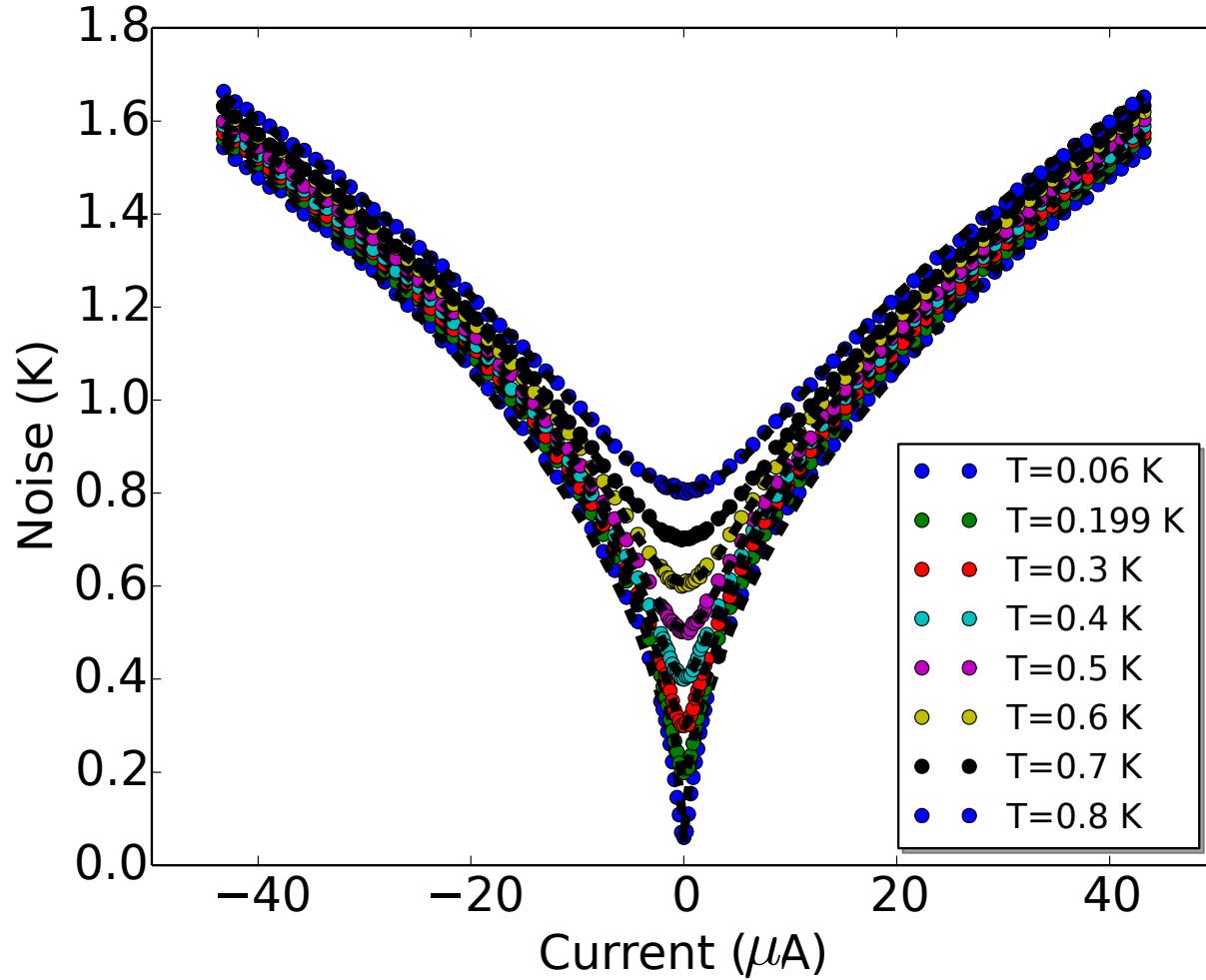
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$$T_e = \left(\frac{RI^2}{\Sigma V} + T_{e-ph}^5 \right)^{1/5}$$

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Thermal conductance:

$$G_{e-ph} = 5\Sigma V T_e^4$$

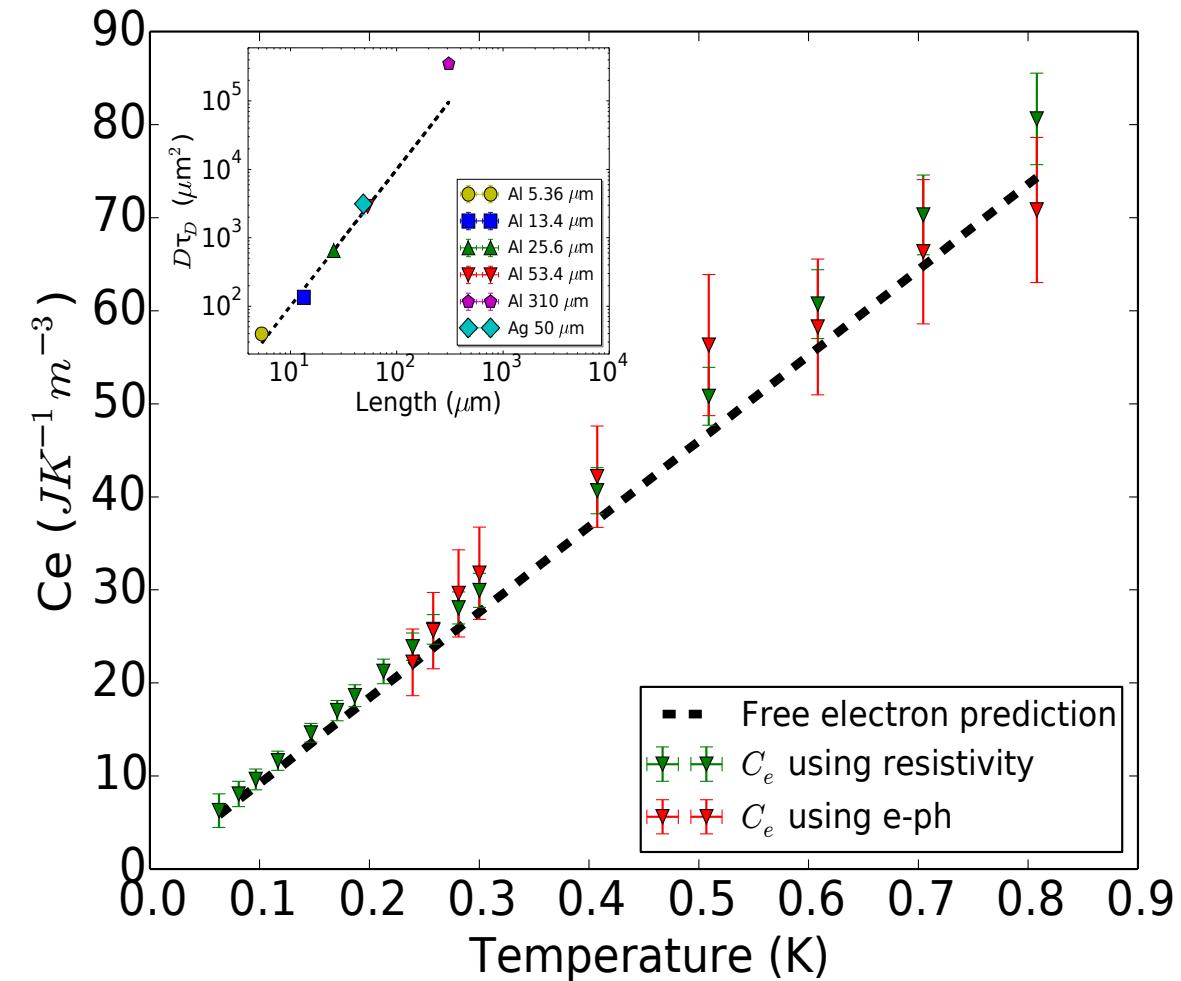
Specific Heat

We can compute the specific heat.

From electron-phonon regime:

$$C_e = G_{e-ph} \tau_{e-ph}$$

Specific Heat

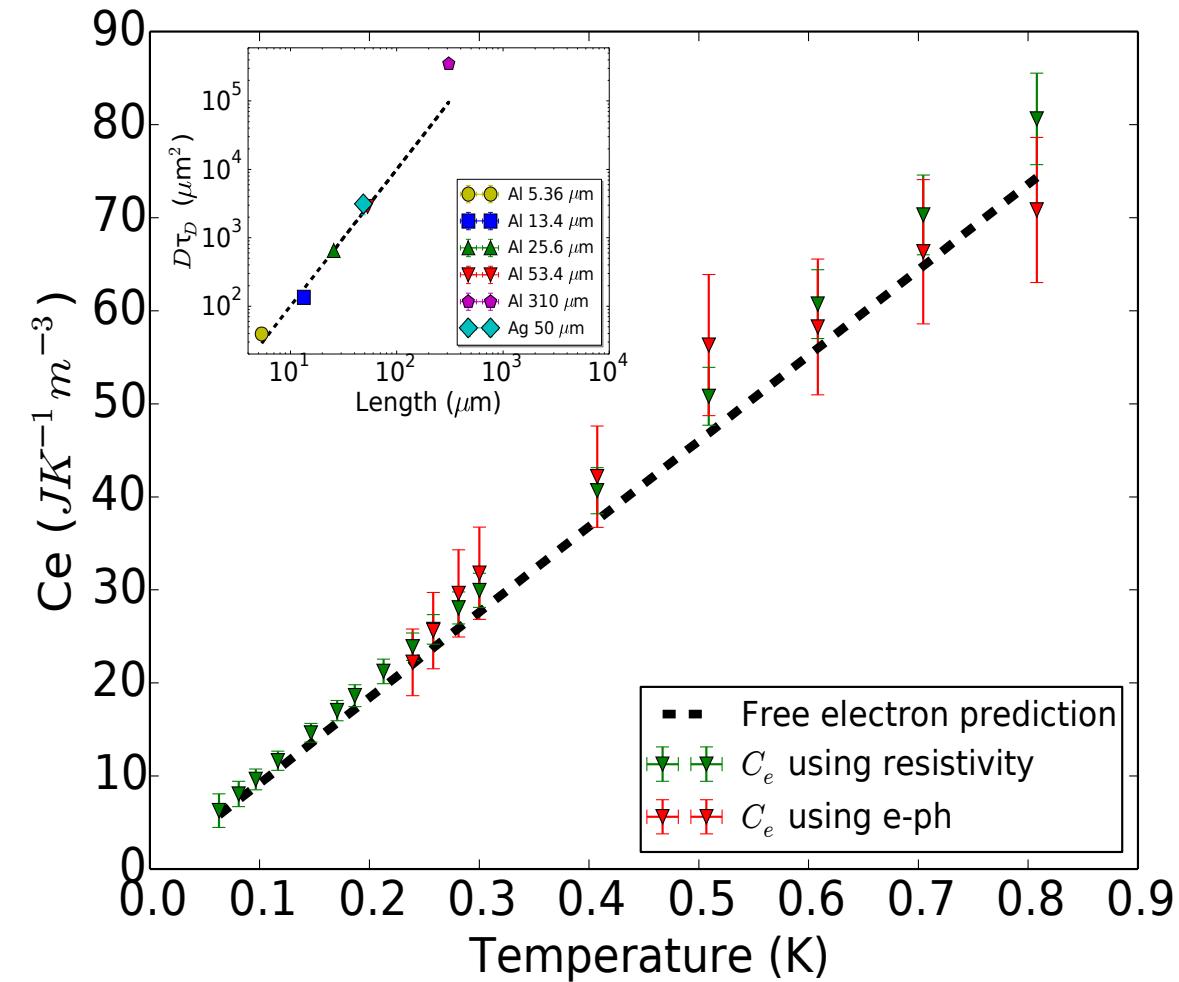


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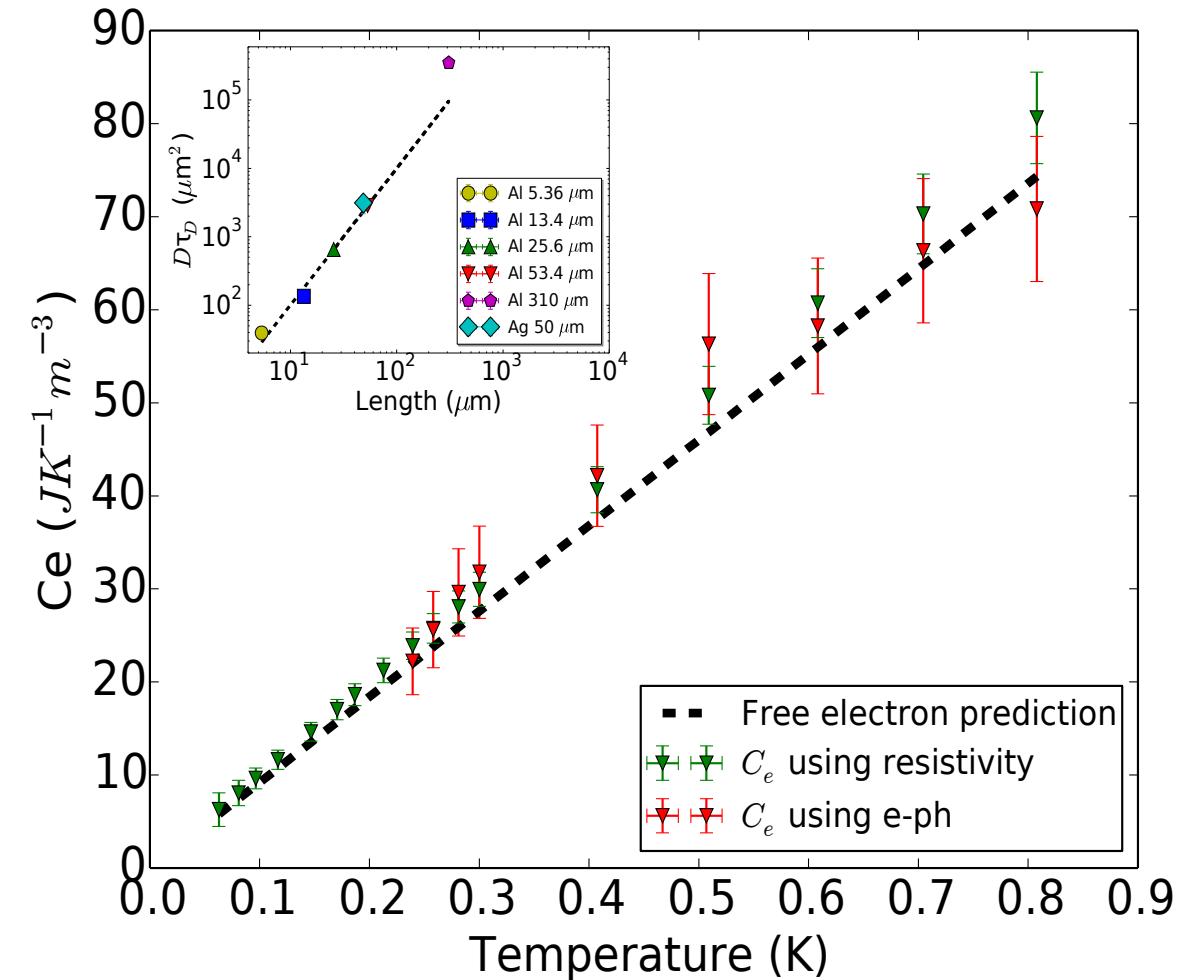
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$$C_e = \gamma T$$

with:

$$\gamma = L G \tau_D$$

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with:

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$$\gamma = 115 \text{ J.K}^{-2} \text{m}^{-3}$$

Effective Mass

Specific heat in a free electron gas:

$$\gamma = \frac{1}{3}\pi^2 k_B^2 n(E_F) = \frac{k_B^2}{3\hbar^2} V k_F m$$

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Metal	Theory			Experiment	
	λ_{ph}	λ_c	m_b/m	m^*/m	m^*/m
Al	0.49(0.38)	- 0.01	1.04 ^a	1.54 ^c	1.48 ^b
	0.53	- 0.01		1.58 ^d	

A. Tari, *The specific Heat of Matter at low Temperatures*, Imperial College Press

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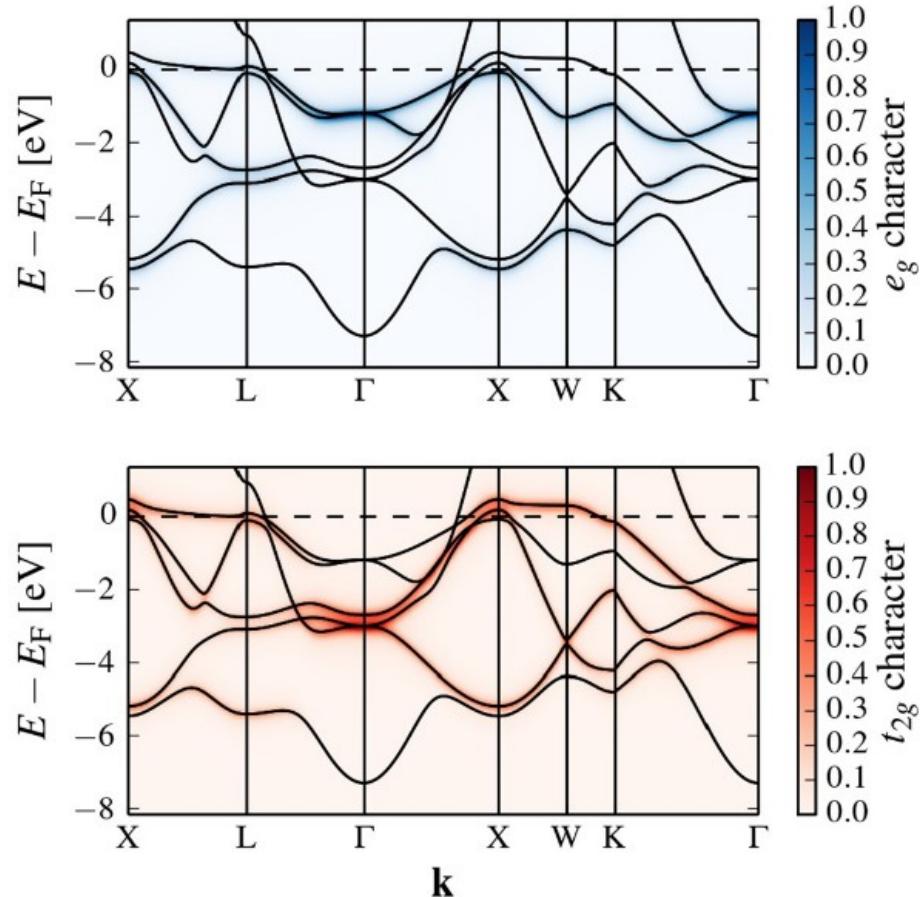
Effective mass and density of states:

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Metal	Theory			Experiment	
	λ_{ph}	λ_c	m_b/m	m^*/m	m^*/m
Al	0.49(0.38)	- 0.01	1.04 ^a	1.54 ^c	1.48 ^b
	0.53	- 0.01		1.58 ^d	

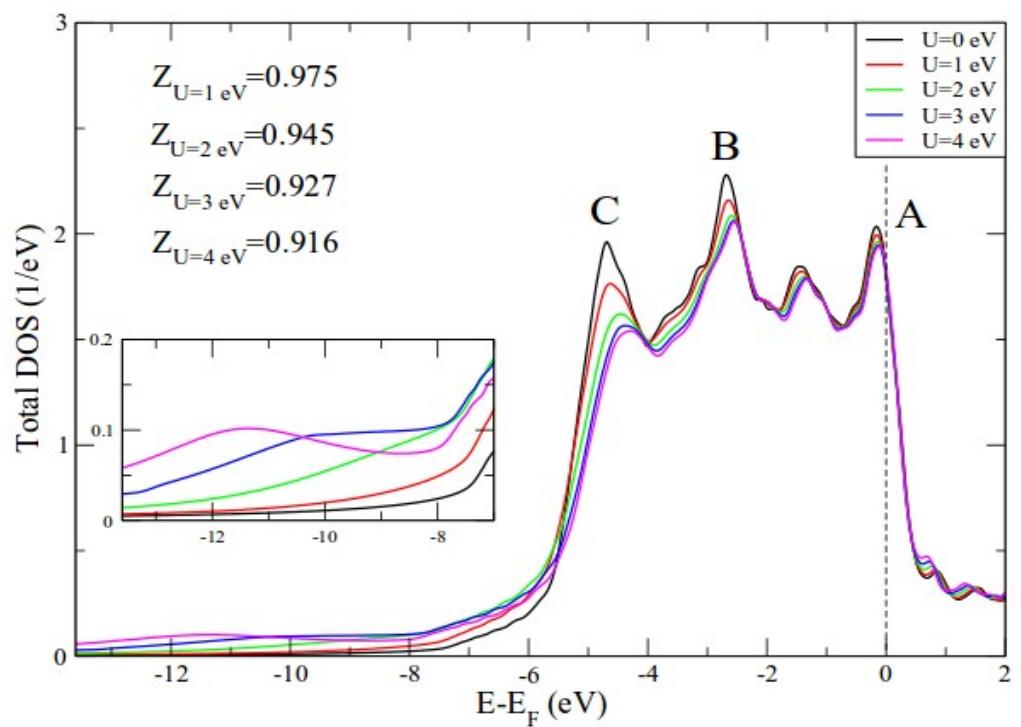
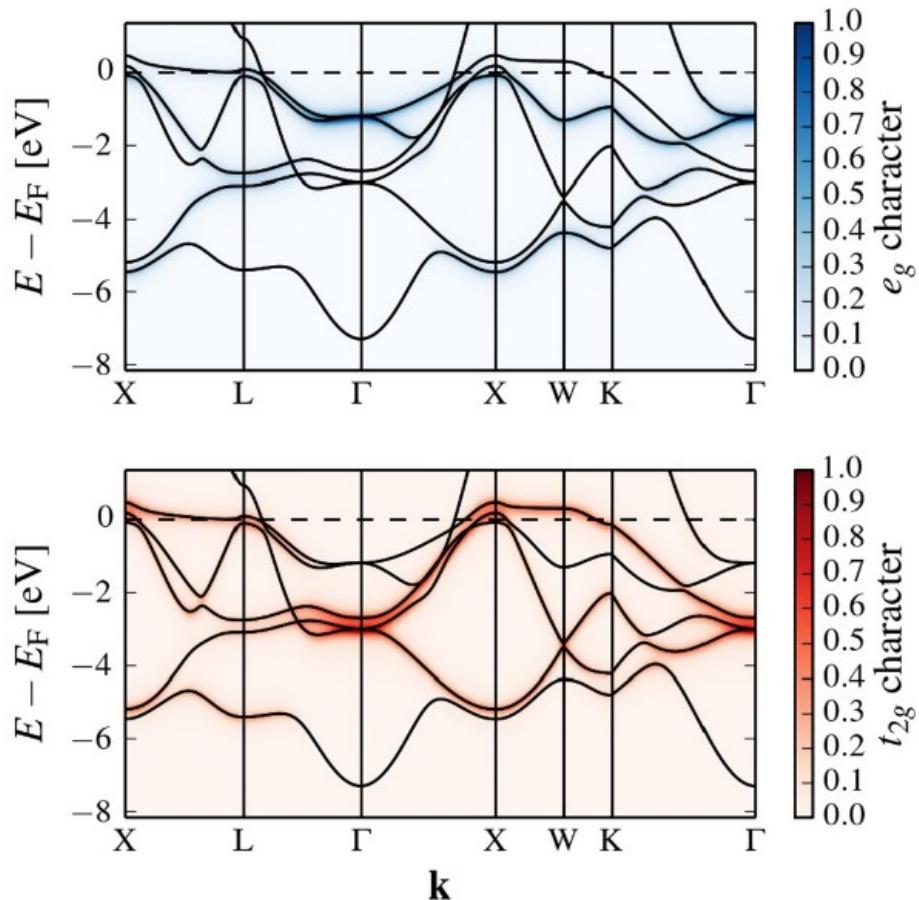
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Band structure of Palladium



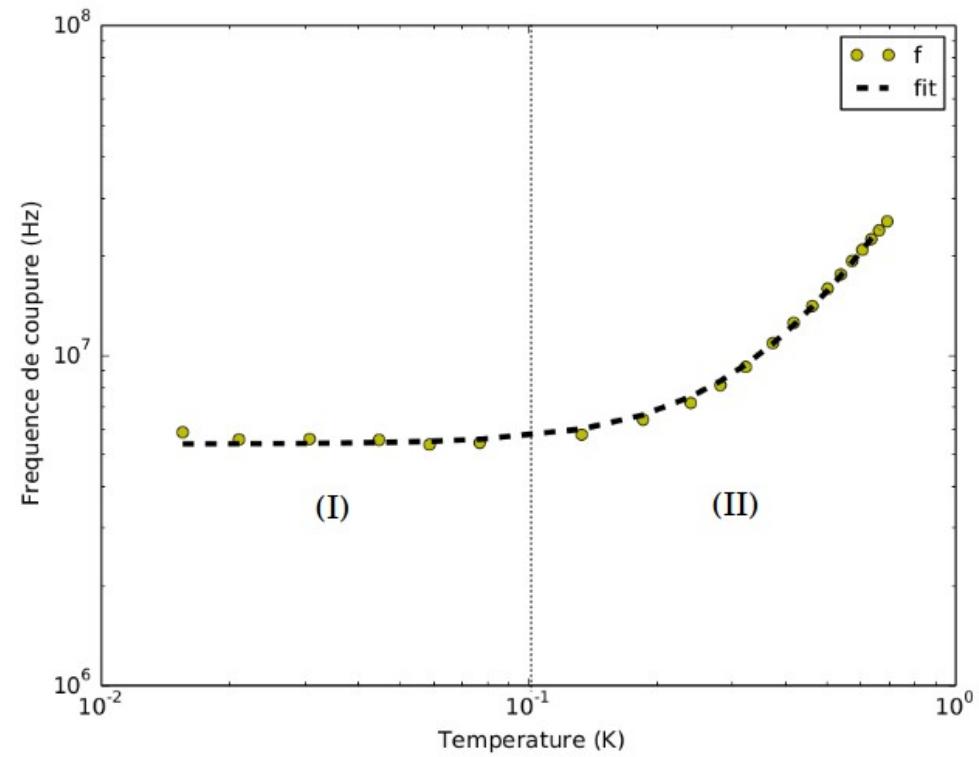
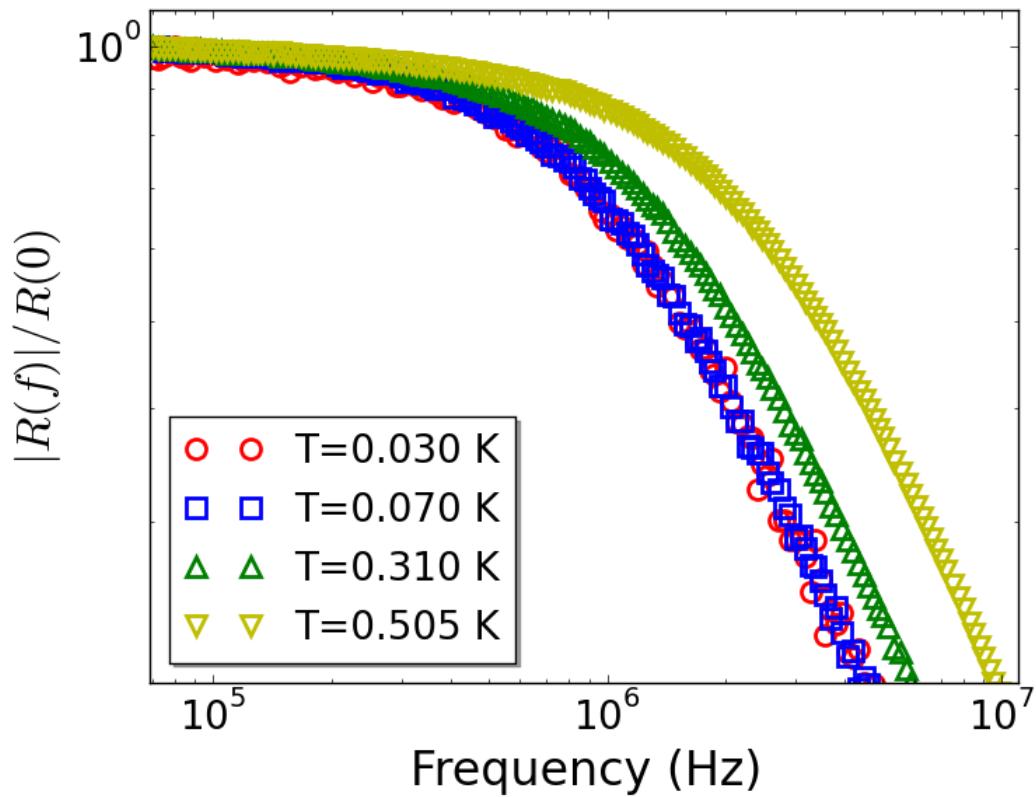
A. Ostlin et al. PRB 93, 155152

Band structure of Palladium



A. Ostlin et al. PRB 93, 155152

NTI on Palladium



Effective Mass

Specific heat measured in our Palladium sample:

$$\gamma = 830 \text{ J.K}^{-2}.\text{m}^{-3}$$

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Which gives an effective mass:

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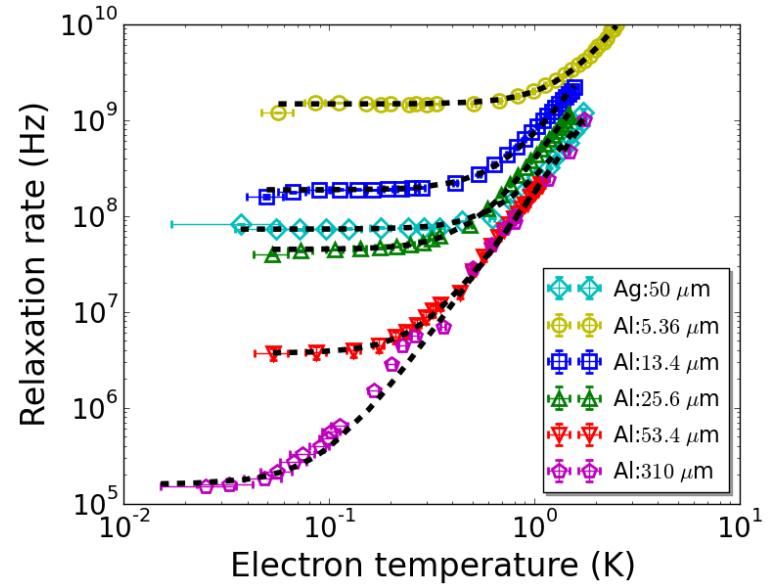
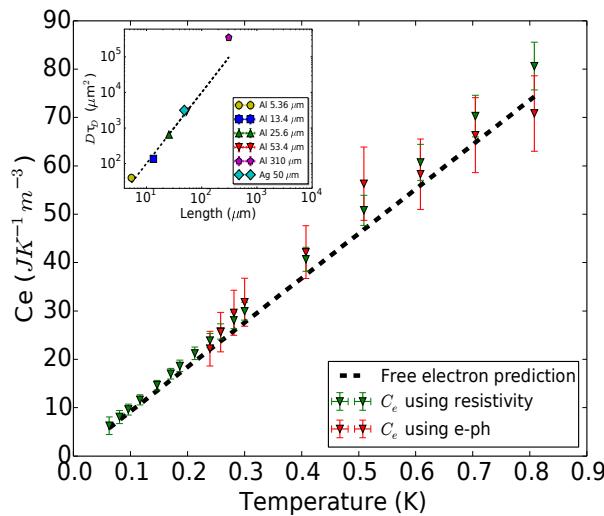
$$\frac{m^*}{m} = 1.33$$

Metal	Theory			Experiment	
	λ_{ph}	λ_c	m_b/m	m^*/m	m^*/m
Pd	0.20 ^e	0.37 ^f	1.66 ^g	1.71 ^h	

Conclusion

We have developed a technique to measure directly relaxation times such as:

- Electron-phonon time.
- Diffusion time.



- By combining with resistivity, we were able to check the validity of the Wiedemann-Franz law
- By combining with noise measurements, we were able to extract the specific heat of our samples.